

1 Aux 3/

$$U(x) = V \cos(\alpha x) - Fx, \quad V, F \text{ const.}$$

ptos. eq. y estabilidad

Lo L: es pensar que V y F son no nulos

1) $U(x)$ no es ident. nulo ($V \neq 0, F \neq 0$)

$$U(x) = V \cos(\alpha x) - Fx$$

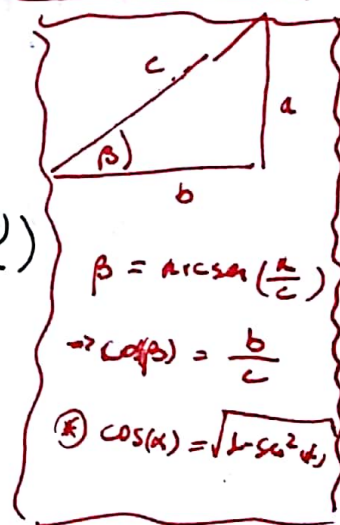
ptos $\Rightarrow U'(x) = -\alpha V \sin(\alpha x) - F \stackrel{!}{=} 0$

$$\Rightarrow \sin(\alpha x^*) = \frac{-F}{\alpha V} \therefore \arcsen\left(\frac{F}{\alpha V}\right) = x^*$$

estabilidad $\Rightarrow U''(x) = -\alpha^2 V \cos(\alpha x) + 0$

$$\Rightarrow U''(x^*) = -\alpha^2 V \cos\left(\alpha \arcsen\left(\frac{F}{\alpha V}\right)\right)$$

$$\Rightarrow U''(x^*) = -\alpha^2 V \sqrt{1 - \left(\frac{F}{\alpha V}\right)^2}$$



Estable > 0
Inestable < 0

$$\Rightarrow U''(x^*) > 0 \Leftrightarrow -\alpha^2 V \sqrt{1 - \left(\frac{F}{\alpha V}\right)^2} > 0$$

$$\Rightarrow V < 0 \text{ y } 1 - \left(\frac{F}{\alpha V}\right)^2 > 0 \quad \text{Estable}$$

Caso cont. a $V < 0$ ($V > 0$) es inestable pero

si $1 - \left(\frac{F}{\alpha V}\right)^2 < 0 \Rightarrow$ no existe sol. en $\mathbb{R}$

2) $V \neq 0, F = 0, \alpha \neq 0$

pts $U(x) = V \cos(\alpha x) \Rightarrow U'(x) = -V \alpha \sin(\alpha x)$

$\Rightarrow U'(x^*) \stackrel{!}{=} 0 \Leftrightarrow \sin(\alpha x^*) = 0 \Rightarrow \alpha x^* = k\pi$ $k \in \mathbb{Z}$

$\therefore x_k^* = \frac{k\pi}{\alpha}$

$\left\{ \begin{array}{l} \text{1 } x^* \text{ segun } k \\ \text{(hay infinitas)} \end{array} \right\}$

Est. $U''(x) = -V \alpha^2 \cos(\alpha x)$

$= \pm 1$ segun k
 $\Leftrightarrow$ par o impar

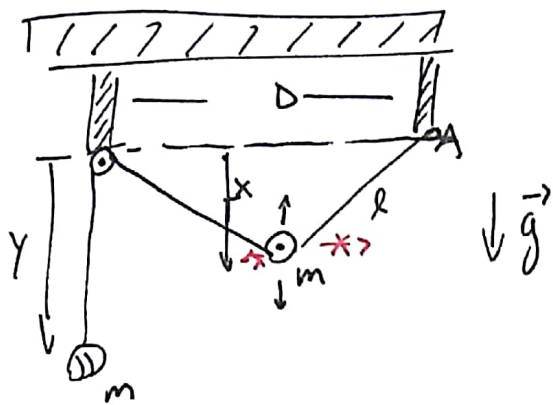
$\Rightarrow U''\left(\frac{k\pi}{\alpha}\right) = -V \alpha^2 \cos(k\pi)$

$\Rightarrow$ k par $U''(\cdot) = -V \alpha^2 \Rightarrow V > 0$ inst.

$V < 0$ estable

$\Rightarrow$ k impar $U''(\cdot) = V \alpha^2 \Rightarrow V > 0$ estable

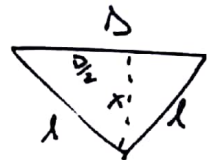
$V < 0$ inestable



Largo : L

a) $y \sim x$ y eq.

$\Rightarrow$ Sabemos que $L = 2l + y$, donde



$$\Rightarrow L = 2\sqrt{x^2 + \left(\frac{D}{2}\right)^2} + y$$

$$\Rightarrow l^2 - \left(\frac{D}{2}\right)^2 = x^2$$

$$l = \sqrt{x^2 + \left(\frac{D}{2}\right)^2}$$

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$$y(x) = L - \sqrt{4x^2 + D^2}$$

Como solo hay energia pot. grav.

$$\Rightarrow U = -m(gx + gy)$$

$$\Rightarrow U(x) = -mg(x + L - \sqrt{4x^2 + D^2})$$

$$\rightarrow \frac{dU}{dx}(x_{eq}) \stackrel{!}{=} 0 \quad \neq$$

$$= -mg \left(1 + 0 - \frac{4x_{eq}}{\sqrt{D^2 + 4x_{eq}^2}} \right)$$

$$\rightarrow \boxed{x_{eq} = \pm \frac{D}{2\sqrt{3}}}$$

$$* \quad \frac{d^2U}{dx^2}(x_{eq}) = \frac{4mg D^2}{(D^2 + 4x_{eq}^2)^{3/2}} \Rightarrow \frac{d^2U}{dx^2}(x_{eq}) = \frac{3\sqrt{3} mg}{2D} > 0$$

pos. de eq.
estable

b)

$$K = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} m \dot{y}^2$$

$$\hookrightarrow \dot{y} = \frac{dy}{dx} \cdot \dot{x} = - \frac{4x\dot{x}}{\sqrt{D^2 + 4x^2}}$$

$$\Rightarrow K = \frac{1}{2} m \underbrace{\left(\frac{D^2 + 20x^2}{D^2 + 4x^2} \right)}_{L(x)} \dot{x}^2$$

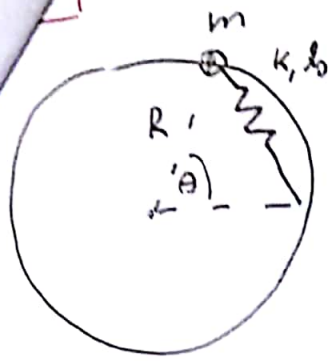
Cuando $x = x(x)$, podemos sup. que cerca de x_{eq}

$$L = L(x_{eq}) = \text{cte}$$

$$\Rightarrow K = \frac{1}{2} (2m) \dot{x}^2$$

$\Rightarrow$

$$\frac{d^2 U}{dx^2}(x_{eq}) \Rightarrow \boxed{\omega_{p.o.}^2 = \frac{3\sqrt{3}g}{4D}}$$



*) pto eq. , est. y free. pto. (pto. estable)

Teo. coseno $L^2 = R^2 + R^2 - 2R^2 \cos(\theta)$

$$\Rightarrow L = R \sqrt{2(1 - \cos(\theta))}$$

*) Energia pot. $\rightarrow U = \frac{k}{2} [R \sqrt{2(1 - \cos(\theta))} - l_0]^2$

$$\Rightarrow \frac{\partial U}{\partial \theta} = \frac{2k}{2} (R \sqrt{2(1 - \cos(\theta))} - l_0) \left(\frac{R \sin(\theta)}{\sqrt{2(1 - \cos(\theta))}} \right) \stackrel{!}{=} 0$$

$$\Rightarrow \sin(\theta_{eq}) = 0$$

$$\hookrightarrow \boxed{\theta_{eq} = 0, \pi}$$

*) Cuando vemos el denominador, este tambien se va a cero $\rightarrow \theta_{eq} = 0$ no puede ser un eq.

$$*) R \sqrt{2(1 - \cos(\theta))} - l_0 = 0$$

$$\Rightarrow \cos(\theta) = 1 - \frac{1}{2} \left(\frac{l_0}{R} \right)^2$$

$$\Rightarrow l_0 < 2R$$

no tendríamos solución cuando $\cos(\theta) < -1$

$$\begin{aligned}
 \frac{\partial^2 U}{\partial \theta^2} \bigg|_{\theta=\pi} &= \frac{KR}{\sqrt{2}} \frac{(R\sqrt{2(1-\cos(\theta))} - l_0) \cos(\theta)}{\sqrt{2(1-\cos(\theta))}} \bigg|_{\theta=\pi} = \frac{KR}{\sqrt{2}} \cdot \frac{(2R - l_0)(-1)}{2} \\
 &= \frac{KR}{2\sqrt{2}} (l_0 - 2R)
 \end{aligned}$$

$$\Rightarrow \frac{\partial^2 U}{\partial \theta^2} \bigg|_{\theta=\pi} = \frac{KR}{2\sqrt{2}} (l_0 - 2R) \quad \rightarrow \quad \begin{aligned} & l_0 > 2R \quad \text{Est.} \\ & l_0 < 2R \quad \text{Inst.} \end{aligned}$$

$$\begin{aligned}
 \frac{\partial^2 U}{\partial \theta^2} \bigg|_{\cos(\theta) = 1 - \frac{1}{2}(\frac{l_0}{R})^2} &= \frac{KR^2 \sin(\theta)}{\sqrt{2(1-\cos(\theta))}} \cdot \frac{R \sin(\theta)}{\sqrt{2(1-\cos(\theta))}} \\
 &= \frac{KR^2 \overset{(1-\cos^2(\theta))}{\sin^2(\theta)}}{2(1-\cos(\theta))} \bigg|_{\cos(\theta)=\dots} = \frac{KR^2}{2} (1 + \cos(\theta)) \bigg|_{\cos(\theta)=\dots} \\
 &= \frac{KR^2}{2} \left(1 + 1 - \frac{1}{2}(\frac{l_0}{R})^2 \right) = KR^2 \left(1 - \frac{1}{4}(\frac{l_0}{R})^2 \right)
 \end{aligned}$$

*) minimum $\cos(\theta_0) = 1 - \frac{1}{2}(\frac{l_0}{R})^2$ see eq.

$\rightarrow$ es stable $\checkmark$

P.O.

Si podemos escribir la energía como

$$E = \frac{1}{2} \dot{x}^2 + U(x) \quad \rightarrow \quad \ddot{x} = -\omega_{p.o.}^2 x$$

$$\Rightarrow \omega_{p.o.}^2 = \frac{U''(x_{eq})}{\alpha}$$

Para este caso tenemos que $\alpha = mR^2 \rightarrow$ inercia

$$\Rightarrow \omega_{p.o.}^2 = \frac{1}{mR^2} \left. \frac{\partial^2 U}{\partial \theta^2} \right|_{\theta_{eq.}}$$

$$\Rightarrow \theta_{eq} = \pi \rightarrow \omega_{p.o.}^2 = \frac{KR}{2\sqrt{2}} (l - 2R), \quad l_0 > 2R$$

$$\Rightarrow \cos(\theta_{eq}) = 1 - \frac{1}{2} \left(\frac{l_0}{R} \right)^2 \rightarrow \omega_{p.o.}^2(\dots) = KR^2 \left(1 - \frac{1}{4} \left(\frac{l_0}{R} \right)^2 \right), \quad l_0 < 2R$$