

Physics of Electronics:

5. Semiconductors

July – December 2009

Contents overview

- Electrical conduction in semiconductors.
- Continuity equation.
- Semiconductor measurements.
- Diodes

Electrical conduction in sC.

- Current Flow = response to \mathcal{E} + diffusion

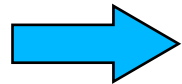
$$J_e = ne\mu_e \mathcal{E} + eD_e \nabla n$$

$$J_h = pe\mu_h \mathcal{E} - eD_h \nabla p$$

- Diffusion and mobility:

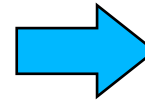
$$D = \frac{\tau_r kT}{M}$$

$$\mu = e\tau_r/m$$



$$D_e = (kT/e)\mu_e$$

$$D_h = (kT/e)\mu_h$$



$$D_e/\mu_e = D_h/\mu_h = kT/e$$

Continuity Equation for Minority Carriers

- I.CASE: No current flow (no gradient & no \mathcal{E})
 - Let's consider a **p-type** material:

Variation of minority carriers = Generation – Recombination

$$\frac{dn}{dt} = G(T) - rn(t)p(t)$$

- Consider first equilibrium:

$$dn/dt = 0 \Rightarrow G(T) = rn_0 p_0 = rn_i^2$$

- Now consider variations from the equilibrium:

$$n = n_0 + \delta n \quad \longleftrightarrow \quad p = p_0 + \delta p \quad \text{with} \quad \delta n = \delta p$$

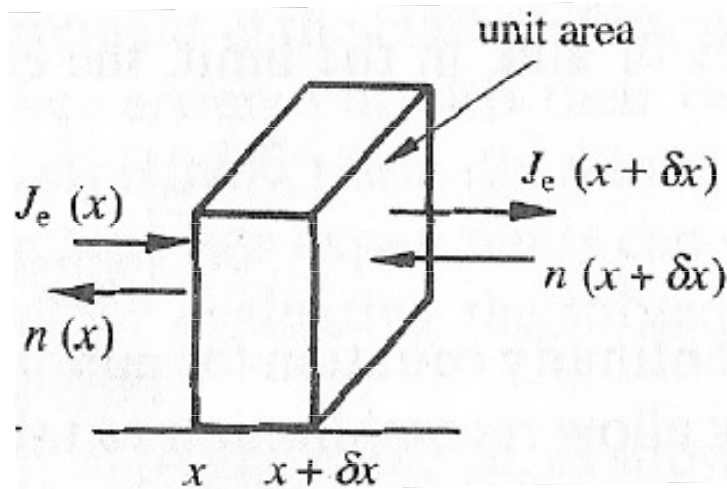
$$\Rightarrow [d(\delta n)/dt]|_{J=0} = -rp_0 \delta n = -\delta n/\tau_{Le}$$

- For a **n-type** material:

$$d(\delta p)/dt = -\delta p/\tau_{Lh}$$

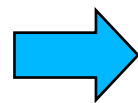
Continuity Equation for Minority Carriers

- II.CASE: With current flow



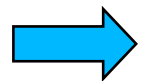
$$J_e = -nev_D$$

$$\left\{ \begin{array}{l} n_{x+\delta x} = J_e(x + \delta x)/ev_D \\ n_x = J_e(x)/ev_D \end{array} \right.$$



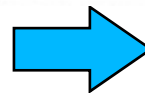
$$n_{x+\delta x} - n_x = \delta n = [J_e(x + \delta x) - J_e(x)]/ev_D$$

$$= \frac{1}{ev_D} \left(J_e(x) + \frac{\partial J_e}{\partial x} \delta x - J_e(x) \right)$$



$$\delta n = \frac{1}{ev_D} \frac{\partial J_e}{\partial x} \delta x$$

$$v_D = \partial x / \partial t$$



$$\left. \frac{\partial(\delta n)}{\partial t} \right|_{\text{current flow}} = \frac{1}{e} \frac{\partial J_e}{\partial x}$$

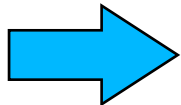
Continuity Equation for Minority Carriers

- II.CASE: Total continuity equation

$$\left. \frac{\partial(\delta n)}{\partial t} \right|_{\text{total}} = -\frac{\delta n}{\tau_{Le}} + \frac{1}{e} \frac{\partial J_e}{\partial x}$$

But $J_e = ne\mu_e \mathcal{E} + eD_e \nabla n$

$$n = n_0 + \delta n$$

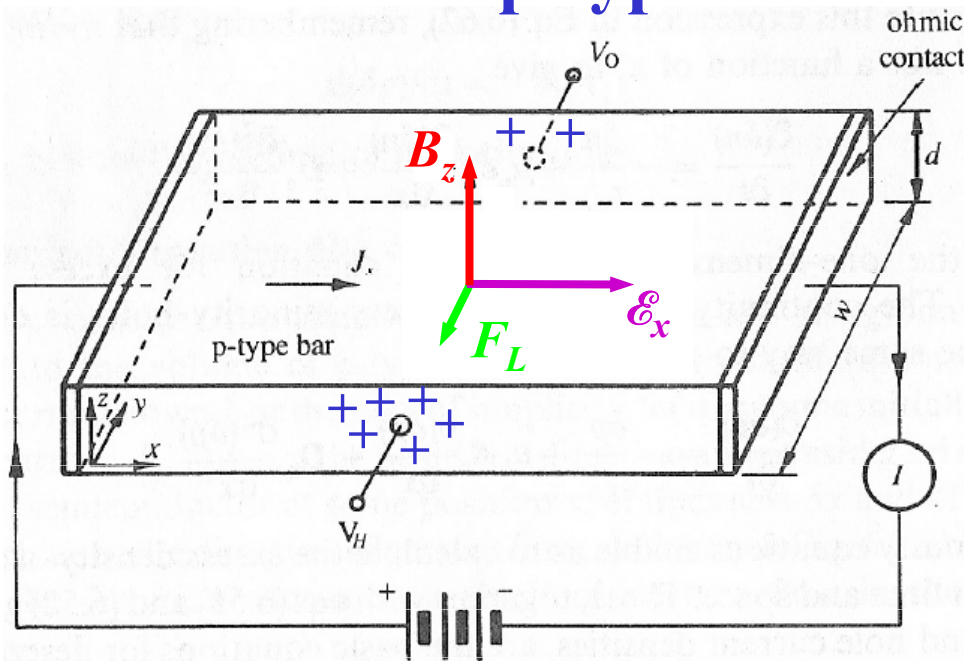


$$\frac{\partial(\delta n)}{\partial t} = -\frac{\delta n}{\tau_{Le}} + \mu_e \mathcal{E}_x \frac{\partial(\delta n)}{\partial x} + D_e \frac{\partial^2(\delta n)}{\partial x^2}$$

(idem for holes in an n-type material)

Semiconductor Measurements

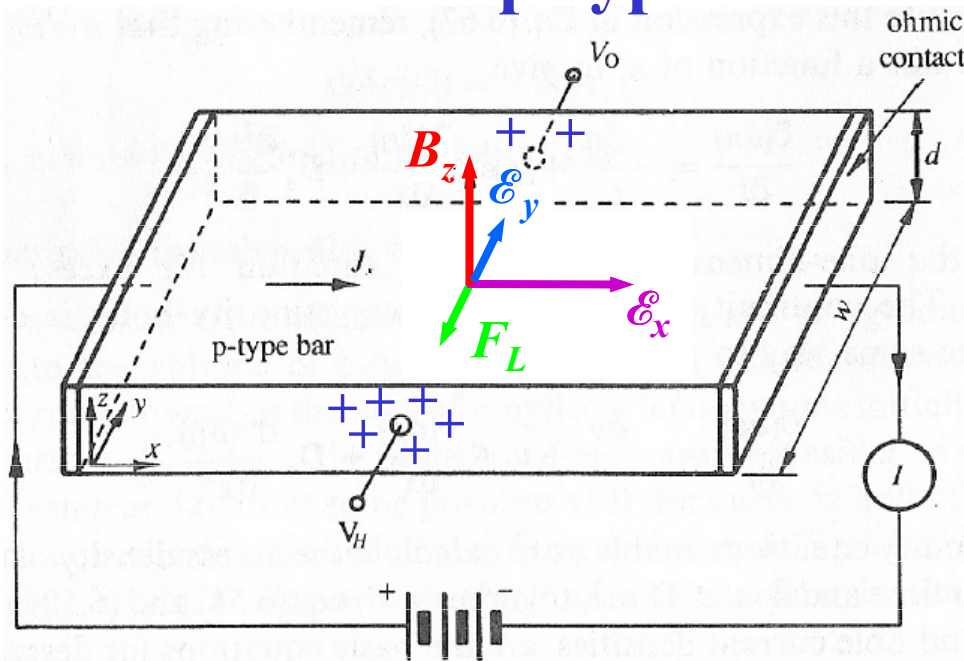
- Hall effect
 - Consider **p-type** material



$$\mathbf{F}_L = e \mathbf{v} \times \mathbf{B} \Rightarrow F_L = e v_{Dx} B_z$$

Semiconductor Measurements

- Hall effect
 - Consider **p-type** material



$$\mathbf{F}_L = e \mathbf{v} \times \mathbf{B} \Rightarrow F_L = e v_{Dx} B_z$$

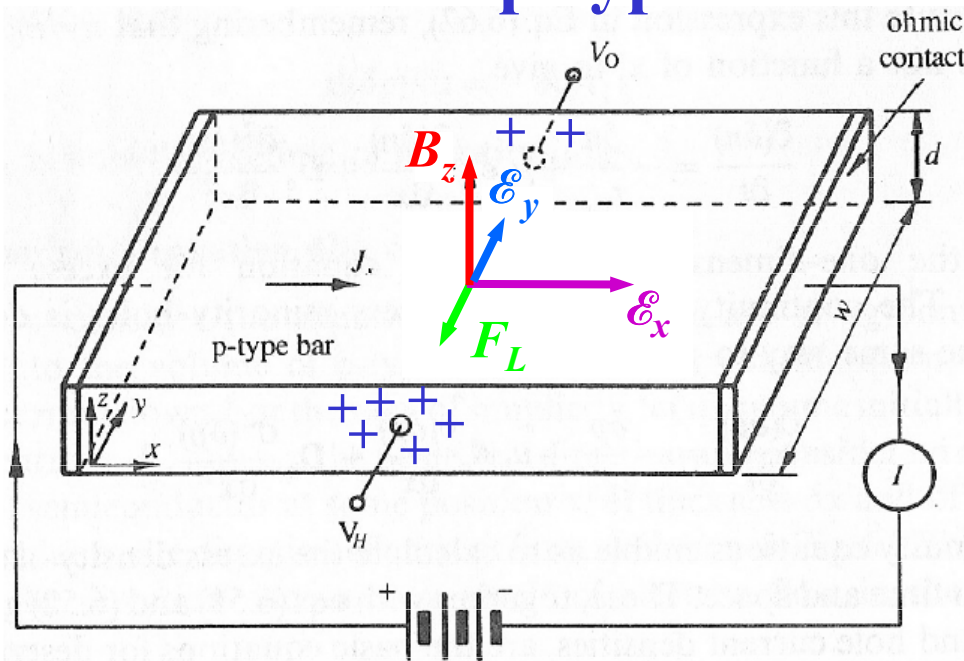
$$\left. \begin{aligned} e\mathcal{E}_y &= F_L = e v_{Dx} B_z \\ &\& J_x \simeq p e v_{Dx} \end{aligned} \right\} \mathcal{E}_y = J_x B_z / p e$$

The **Hall coefficient** is defined:

$$R_H = \mathcal{E}_y / J_x B_z = 1 / p e$$

Semiconductor Measurements

- Hall effect
 - Consider **p-type** material



$$\mathbf{F}_L = e \mathbf{v} \times \mathbf{B} \Rightarrow F_L = e v_{Dx} B_z$$

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$$R_H = \mathcal{E}_y / J_x B_z = 1 / p e$$

- If V_H and I are measured:

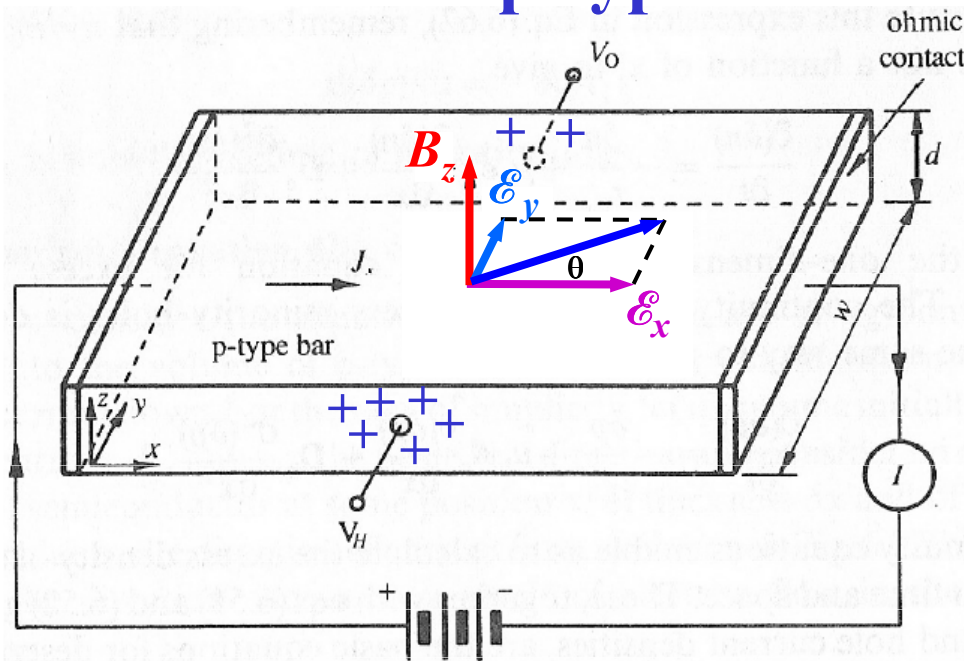
$$V_H = \mathcal{E}_y w$$

$$I = J_x w d$$

$$\left. \begin{aligned} V_H &= \mathcal{E}_y w \\ I &= J_x w d \end{aligned} \right\} R_H = \frac{V_H}{w I B_z} w d = \frac{V_H d}{I B_z} = \frac{1}{p e}$$

Semiconductor Measurements

- Hall effect
 - Consider **p-type** material



The **Hall angle** is defined:

$$\tan \theta = \mathcal{E}_y / \mathcal{E}_x$$

$$\left\{ \begin{aligned} \tan \theta &= \frac{J_x B_z}{pe} \frac{\sigma}{J_x} = R_H B_z \sigma \\ \tan \theta &= \frac{J_x B_z}{pe} \frac{\sigma}{J_x} = \mu_h B_z \end{aligned} \right.$$

$$\mu_h = R_H \sigma$$

- For a **n-type** material:

$$R_{He} = -1/ne$$

Semiconductor Measurements

- Hall effect
 - Consider **both type of carriers** to be present:

$$\begin{array}{l} v_{Dh} = \mu_H \mathcal{E}_x \\ v_{De} = -\mu_H \mathcal{E}_x \end{array} \quad \Rightarrow \quad \begin{array}{l} F_h = -e(v_{Dh} \times B) = -ev_{Dh}B_z \\ F_e = e(v_{De} \times B) = -ev_{De}B_z \end{array}$$

- A net current is created in the y direction

$$\sigma \mathcal{E}_y = e(pv_{yh} - nv_{ye})$$

- Now, using the expressions for the Hall angle

$$v_{yh} = \mu_h \mathcal{E}_y = \mu_h (\mathcal{E}_x \tan \theta) = \mu_h (\mathcal{E}_x \mu_h B_z) \quad \& \quad v_{ye} = \mu_e^2 \mathcal{E}_x B_z$$

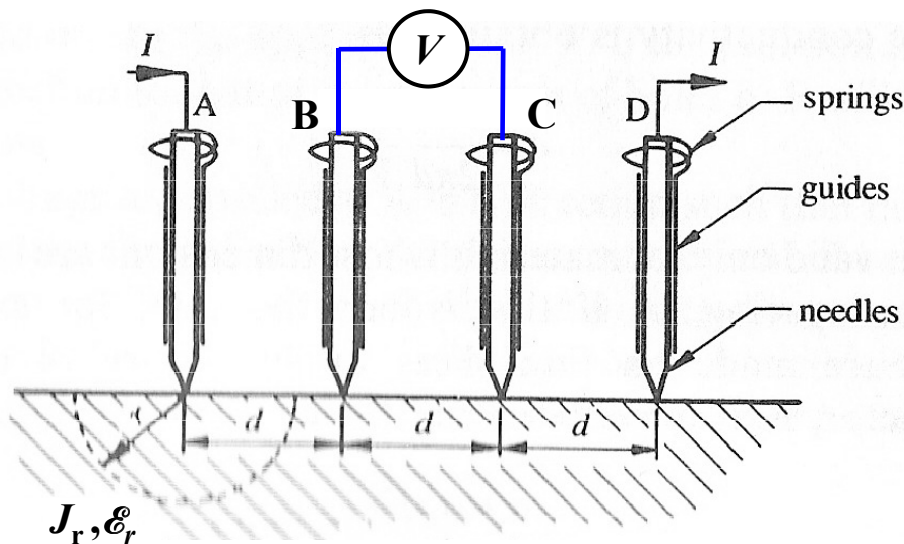
$$\Rightarrow \sigma \mathcal{E}_y = e(p\mu_h^2 - n\mu_e^2) \mathcal{E}_x B_z$$

- Hall coefficient

$$R_H = \frac{\mathcal{E}_y}{J_x B_z} = \frac{e(p\mu_h^2 - n\mu_e^2)}{\sigma^2} \quad \xrightarrow{\sigma = e(n\mu_e + p\mu_h)} \quad R_H = \frac{p\mu_h^2 - n\mu_e^2}{e(p\mu_h + n\mu_e)^2}$$

Semiconductor Measurements

- Four-point probe method for conductivity meas.
 - Consider $d \ll L$



For the current leaving A :

$$J_r = I / (2\pi r^2) \Rightarrow \mathcal{E}_r = J / \sigma = I / (2\pi \sigma r^2)$$

Potential at a distance a from A :

$$V_a = \int_{-\infty}^a \mathcal{E}_r dr = -\frac{I}{2\pi\sigma} \int_{-\infty}^a \frac{1}{r^2} dr = \frac{I}{2\pi\sigma a}$$

$$\Rightarrow V_{BC} = \frac{I}{2\pi\sigma d} - \frac{I}{2\pi\sigma(2d)} = \frac{I}{4\pi\sigma d}$$

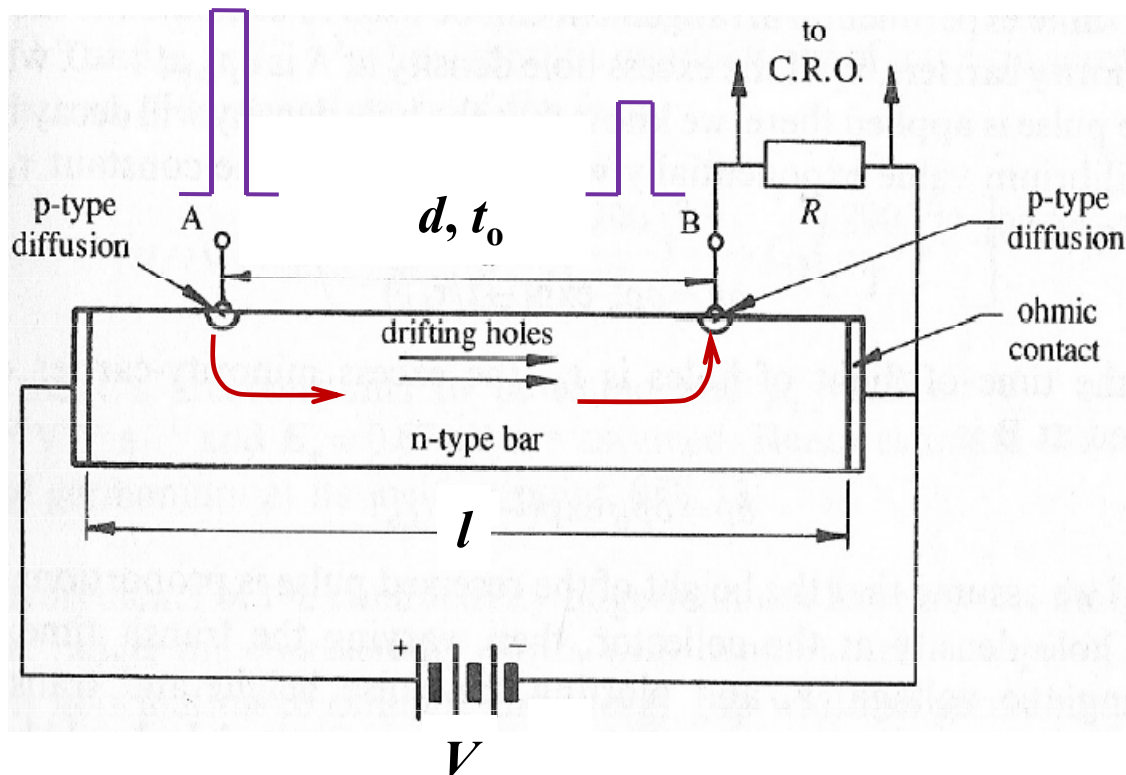
Identically for the current entering D , then:

$$V = 2V_{BC} = \frac{I}{2\pi\sigma d}$$

$$\Rightarrow \sigma = \frac{1}{2\pi d} \frac{I}{V}$$

Semiconductor Measurements

- Minority carrier life-time and mobility.



Mobility:

$$v_D = d/t_0 \Rightarrow \mu E = \mu V/l = d/t_0$$

$$\Rightarrow \boxed{\mu = ld/t_0 V}$$

Life-time:

$$d(\delta p)/dt = -\delta p/\tau_{Lh}$$

$$\Rightarrow \delta p = \delta p_0 \exp(-t/\tau_{Lh})$$

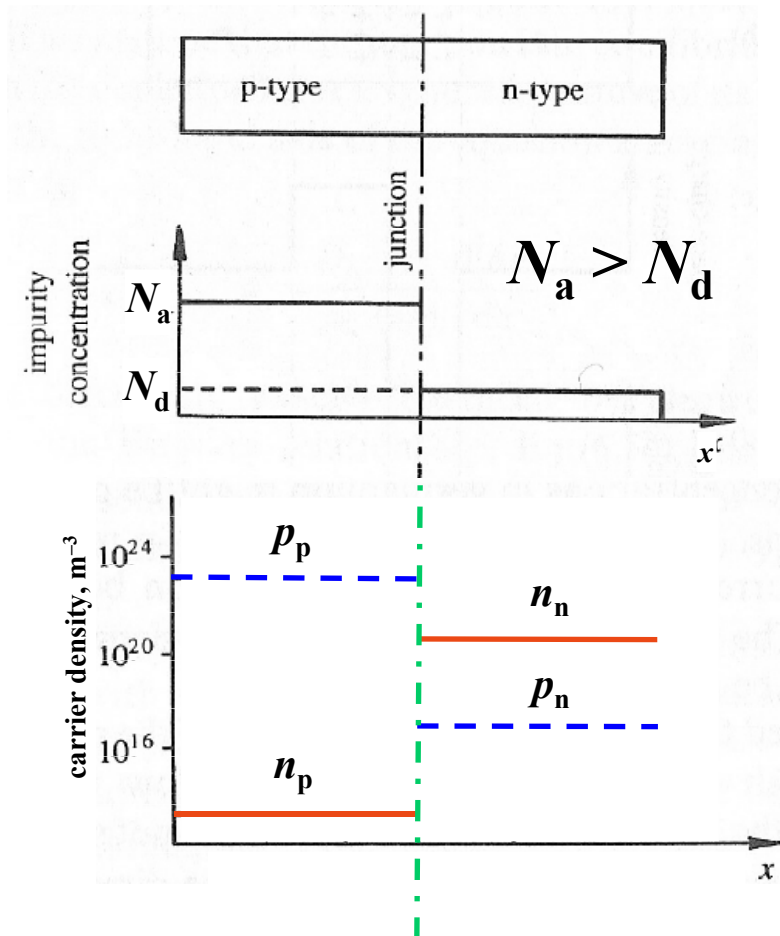
Physics of Electronics:

6. Junction Diodes

July – December 2009

Pn junction in equilibrium

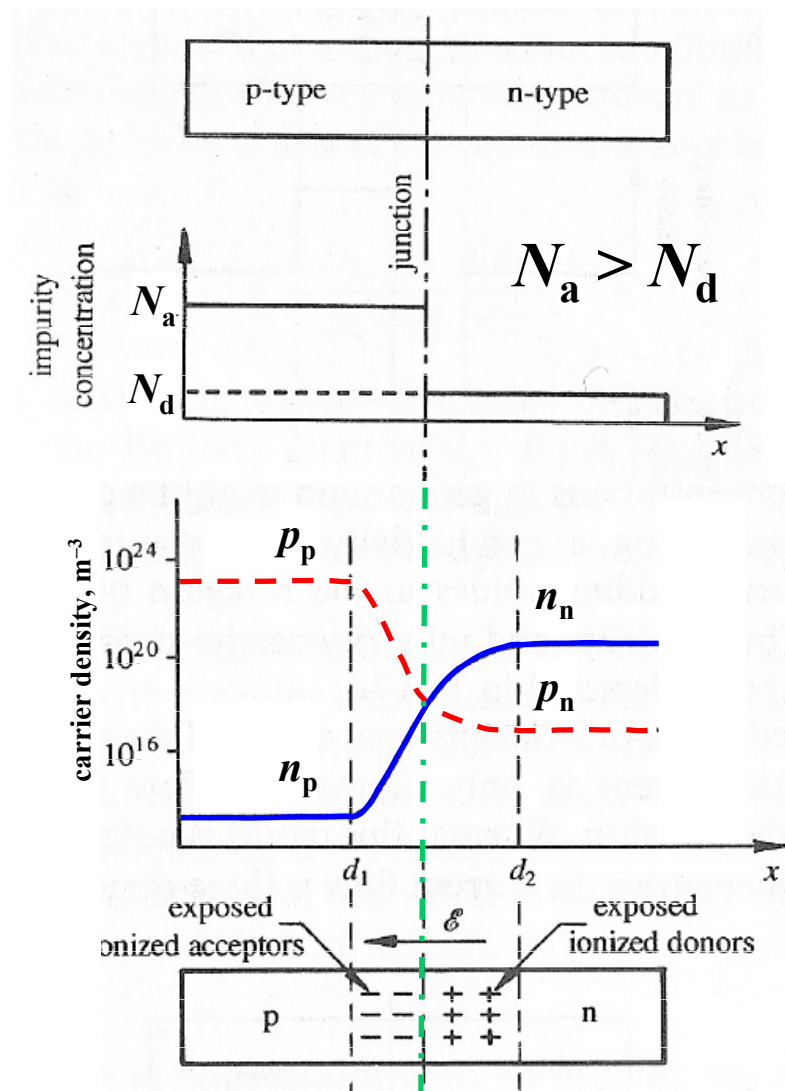
- Consider a pn junction with abrupt transition



Initially, the density of carriers change abruptly at the junction.

Pn junction in equilibrium

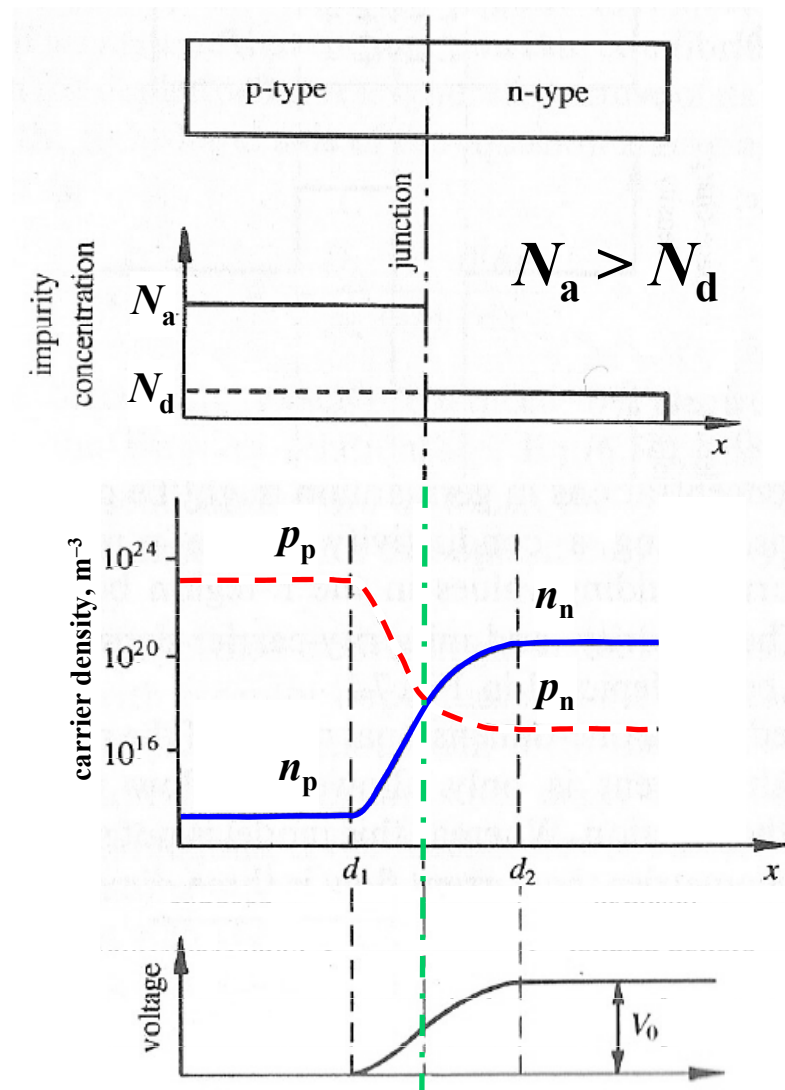
- Consider a pn junction with abrupt transition



This change of concentration causes diffusion which is stopped by exposition of ionized impurities at the transition (called depletion layer $\sim 1\mu\text{m}$).

Pn junction in equilibrium

- Consider a pn junction with abrupt transition



This change of concentration causes diffusion which is stopped by exposition of ionized impurities at the transition (called depletion layer $\sim 1\mu\text{m}$).

In equilibrium, a voltage difference is created across the depletion layer.

Pn junction in equilibrium

- Voltage across the depletion layer
 - Start with continuity eq. for holes

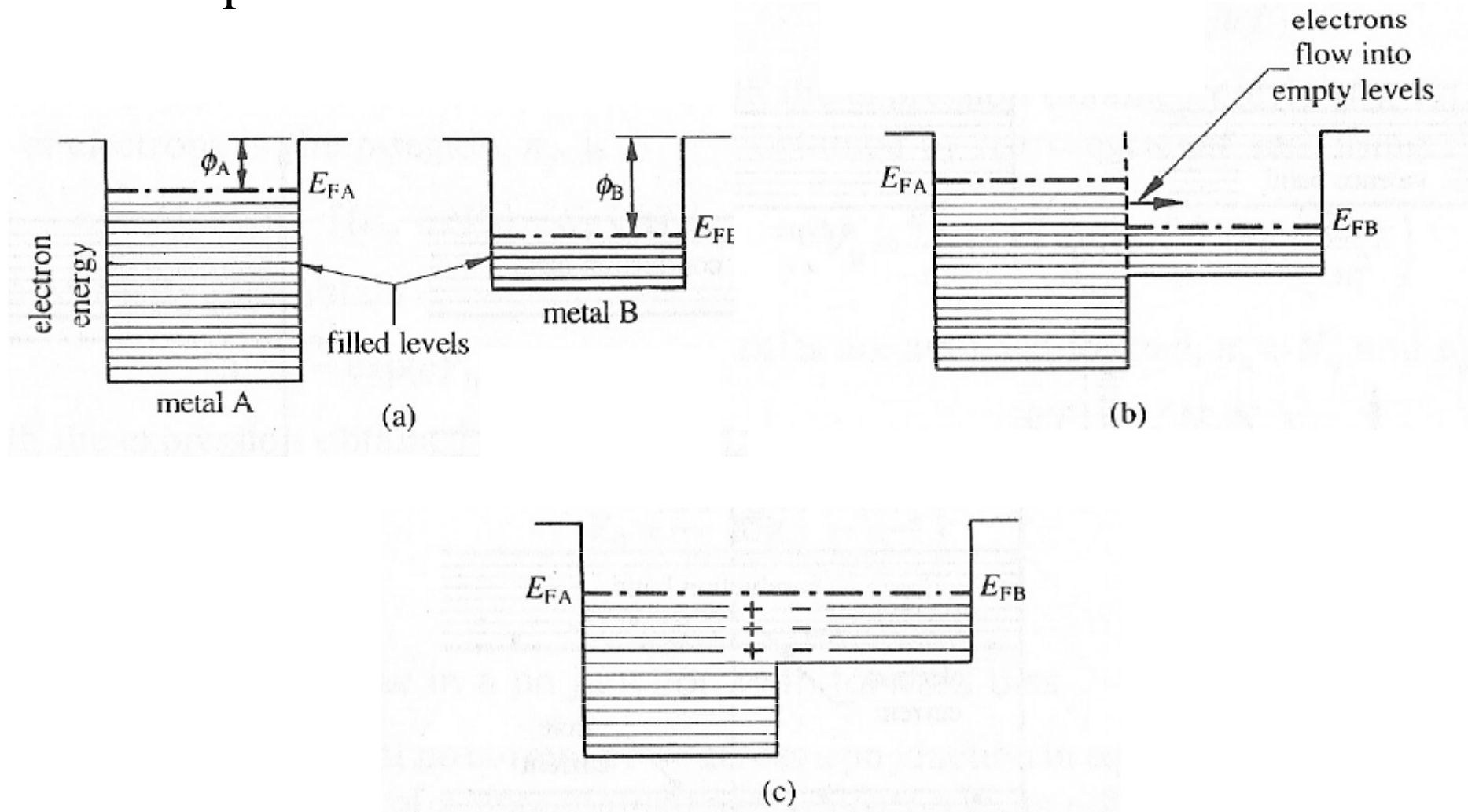
$$\begin{aligned}
 \frac{\partial(\delta p)}{\partial t} &= -\frac{\delta p}{\tau_{Lh}} + \mu_h \mathcal{E}_x \frac{\partial(\delta p)}{\partial x} + D_h \frac{\partial^2(\delta p)}{\partial x^2} \quad \Rightarrow \quad \mathcal{E}_x = \frac{D_h}{\mu_h} \frac{1}{(\delta p)} \frac{d(\delta p)}{dx} \\
 &\quad D_h/\mu_h = kT/e \\
 &\quad \Rightarrow \quad \mathcal{E}_x = \frac{kT}{e} \frac{1}{(\delta p)} \frac{d(\delta p)}{dx} \quad \Rightarrow \quad e \int_{d_1}^{d_2} \mathcal{E}_x \, dx = -kT \int_{p_0}^{p_n} \frac{d(\delta p)}{\delta p} \\
 &\quad \Rightarrow \quad \boxed{p_p = p_n \exp(eV_0/kT)}
 \end{aligned}$$

- Idem for electrons

$$\boxed{n_n = n_p \exp(eV_0/kT)}$$

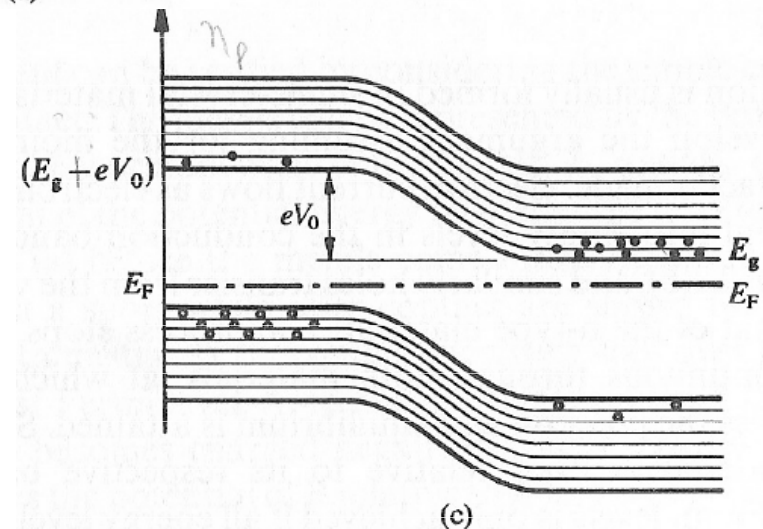
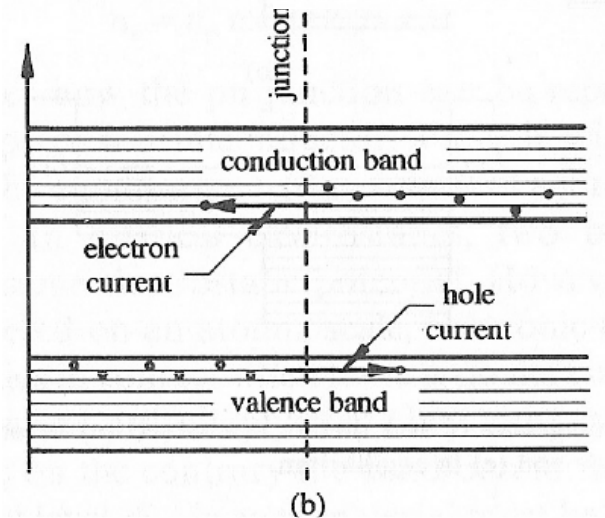
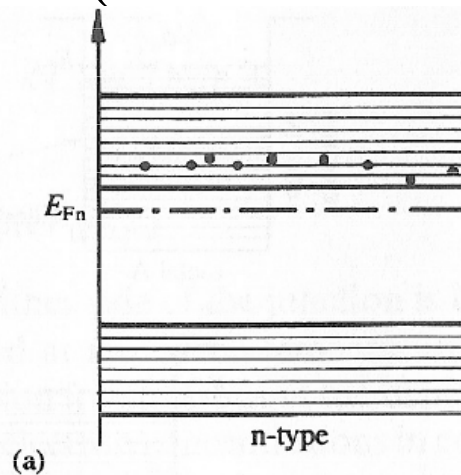
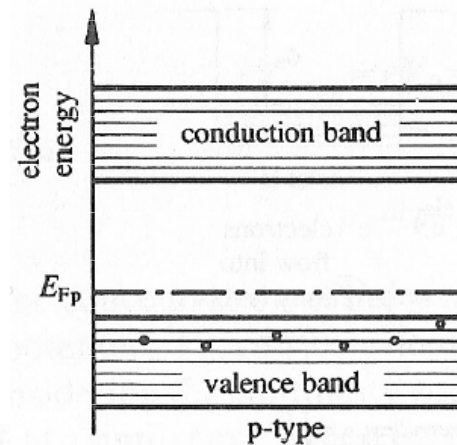
Pn junction in equilibrium

- Junctions and band structure
 - Equilibrium at atomic scale (in metals):



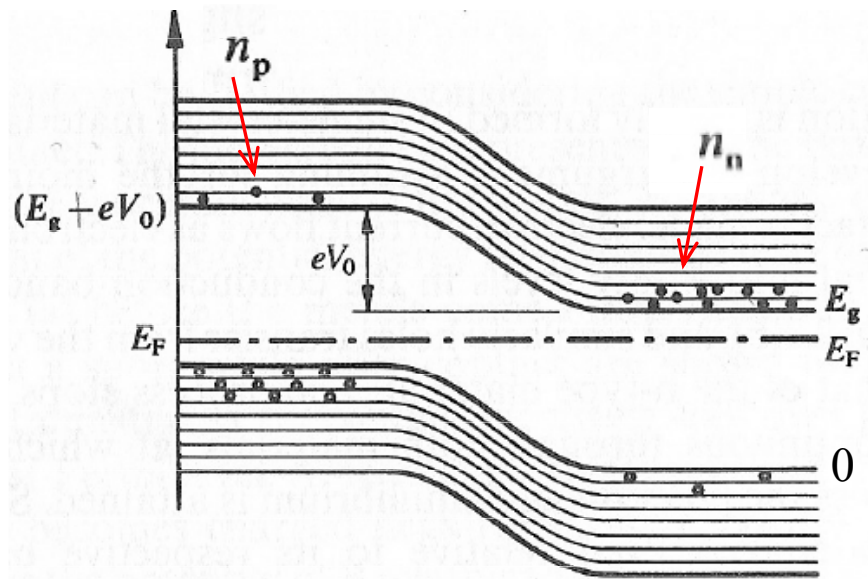
Pn junction in equilibrium

- Junctions and band structure
 - Equilibrium at atomic scale (in semiconductors)



Pn junction in equilibrium

- Number of electrons in the conduction band



$$n_n = N_c \exp\left[-(E_g - E_F)/kT\right]$$

$$n_p = N_c \exp\left[-(\{E_g + eV_0\} - E_F)/kT\right]$$

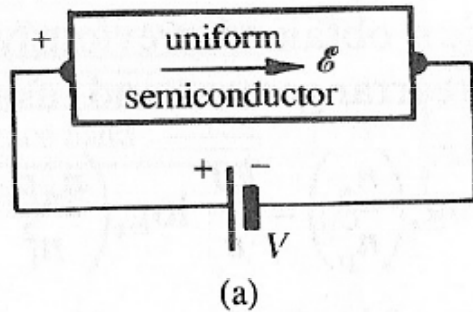
→ $n_n/n_p = \exp(eV_0/kT)$

- If all the impurities are ionized ($n_n \approx N_d$ and $p_p \approx N_a$):

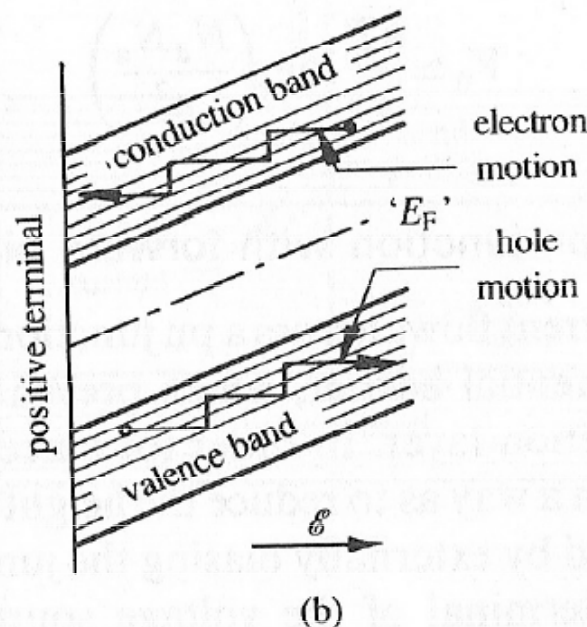
$$V_0 \simeq \frac{kT}{e} \log_e \left(\frac{N_d N_a}{n_i^2} \right)$$

Pn junction with forward bias

- Consider a uniform sC biased with a voltage V :



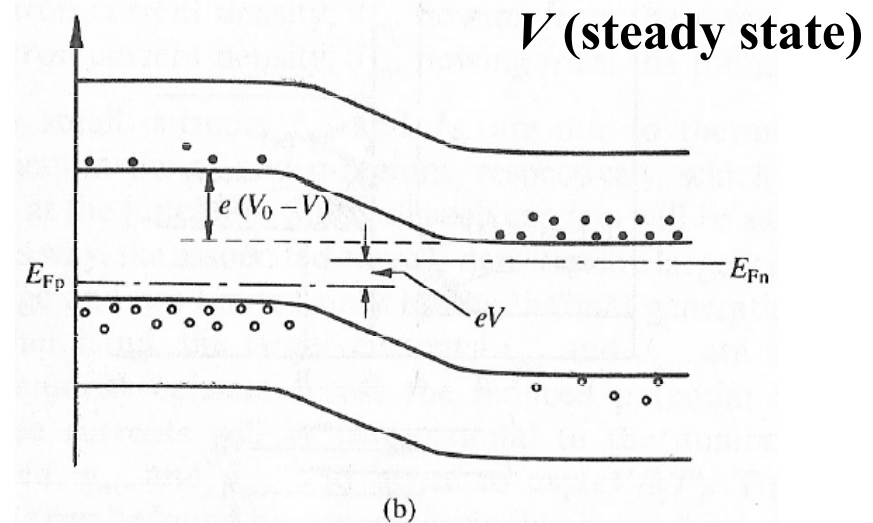
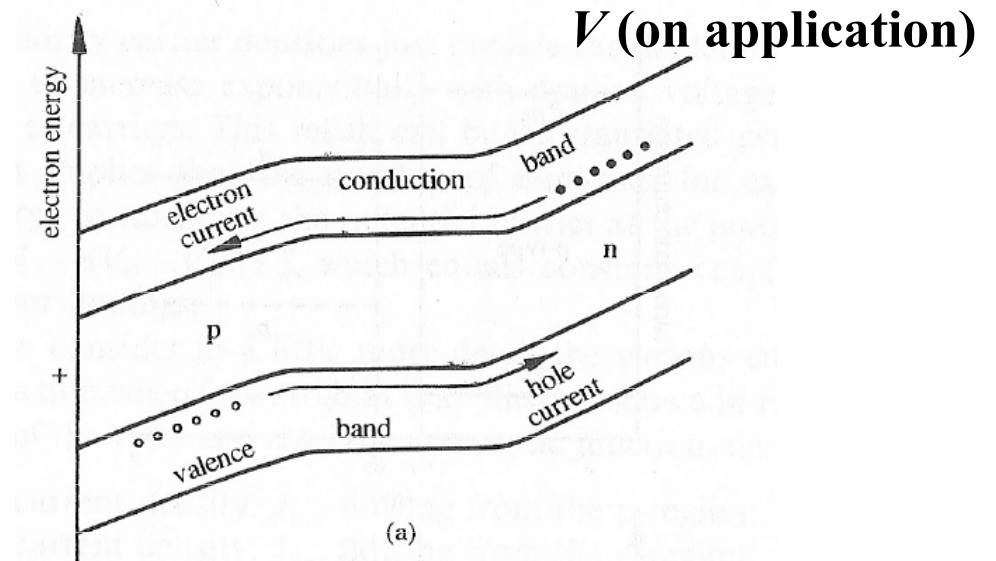
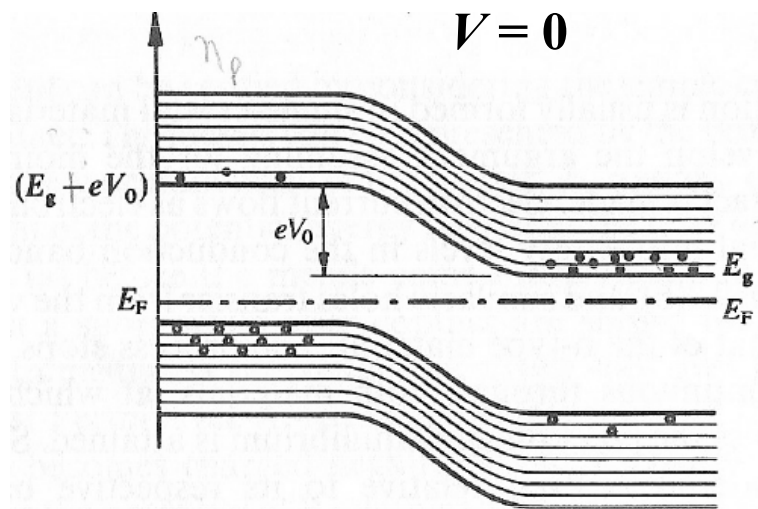
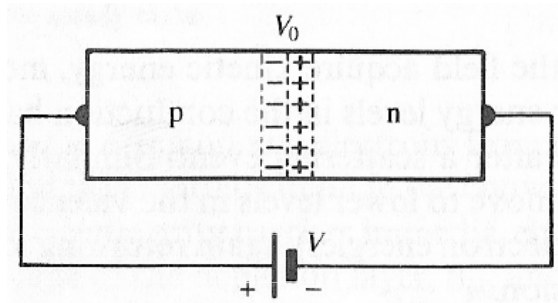
The whole band structure tilts (one end of the sC has more energy than the other) .



Electrons and holes move in the field (in opposite directions).

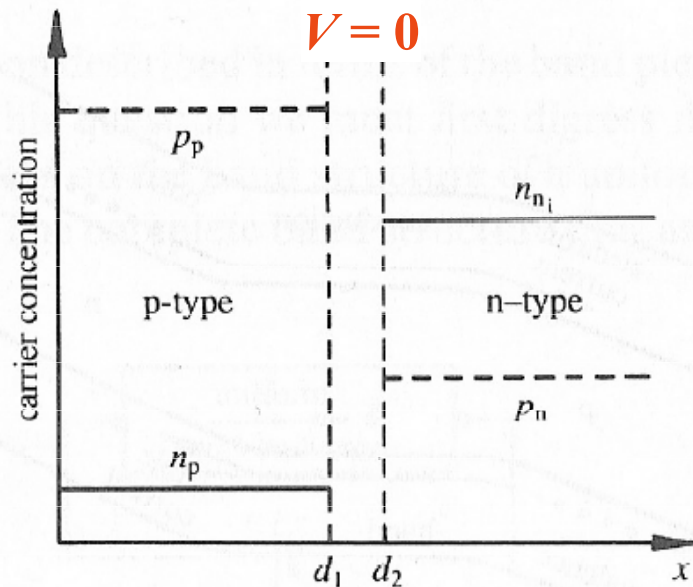
Pn junction with forward bias

- Consider a junction biased with a voltage V :

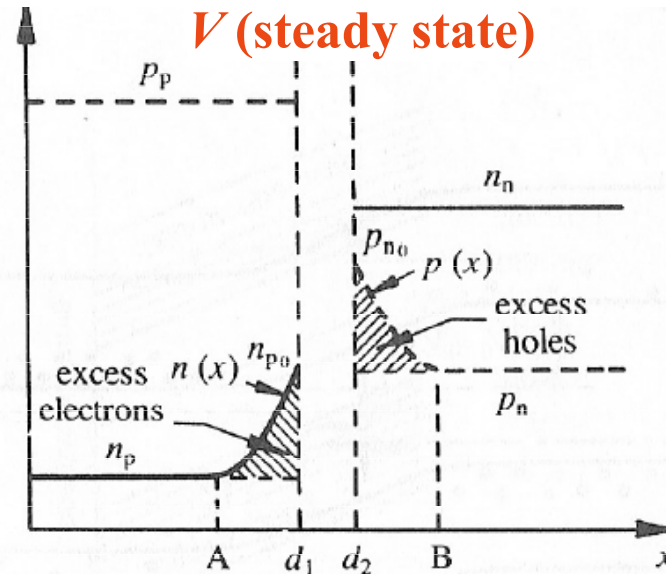


Pn junction with forward bias

- Excess carriers in a biased junctions:



$$p_p = p_n \exp(eV_0/kT)$$



$$p_{n0} = p_p \exp[e(V - V_0)/kT]$$

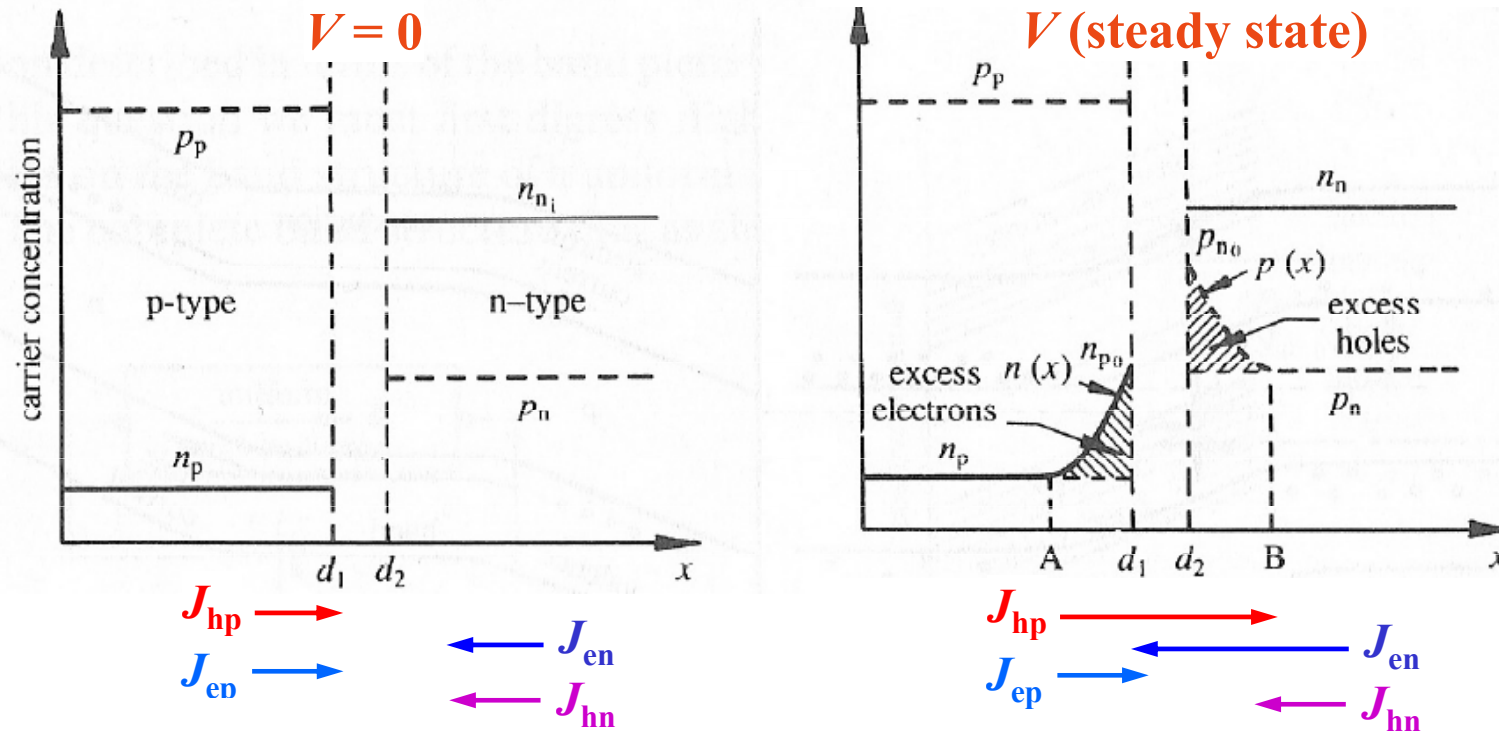


$$p_{n0} = p_n \exp(eV/kT)$$

$$n_{p0} = n_p \exp(eV/kT)$$

Pn junction with forward bias

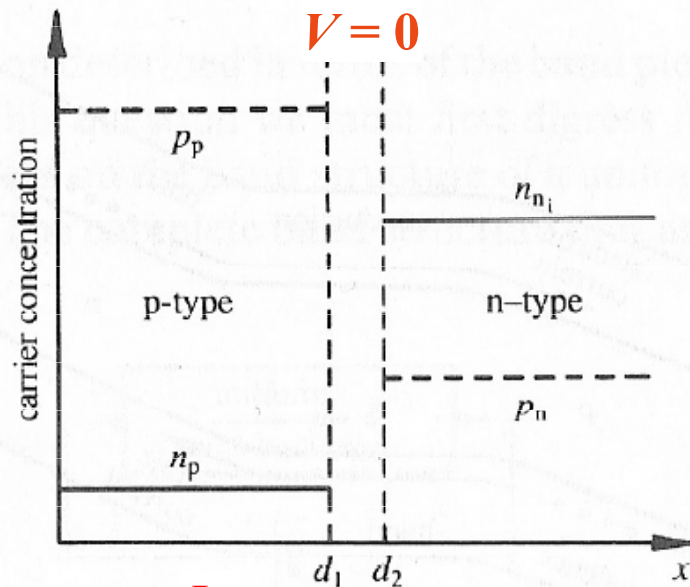
- Excess carriers in a biased junctions:



- (a) a hole current density, J_{hp} , flowing from the p-region;
 (b) a hole current density, J_{hn} , flowing from the n-region;
 (c) an electron current density, J_{en} , flowing from the n-region;
 (d) an electron current density, J_{ep} , flowing from the p-region.

Pn junction with forward bias

- Excess carriers in a biased junctions:

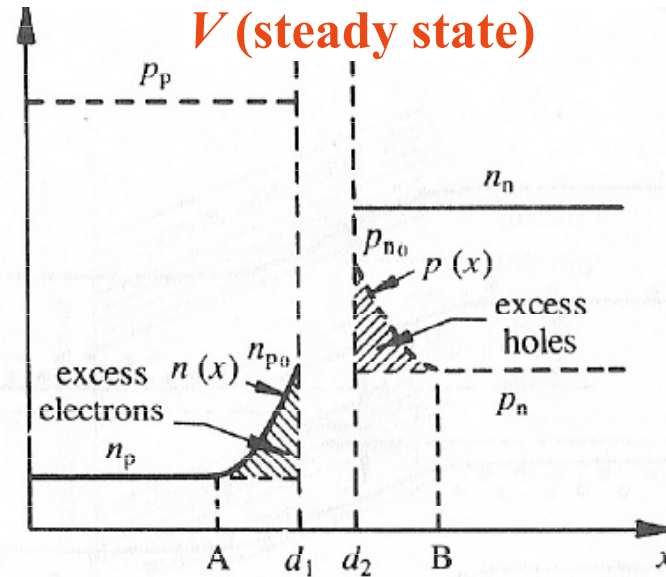


$$J_{hp} \rightarrow \quad \leftarrow J_{en}$$

$$J_{en} \rightarrow \quad \leftarrow J_{hn}$$

$$J_{hp} = J_{hn}$$

$$J_{en} = J_{ep}$$



$$J_{hp} \rightarrow \quad \leftarrow J_{en}$$

$$J_{ep} \rightarrow \quad \leftarrow J_{hn}$$

$$J_{hp} \propto p_{n0} = p_n \exp(eV/kT)$$



$$J_{hp} = J_{hn} \exp(eV/kT)$$

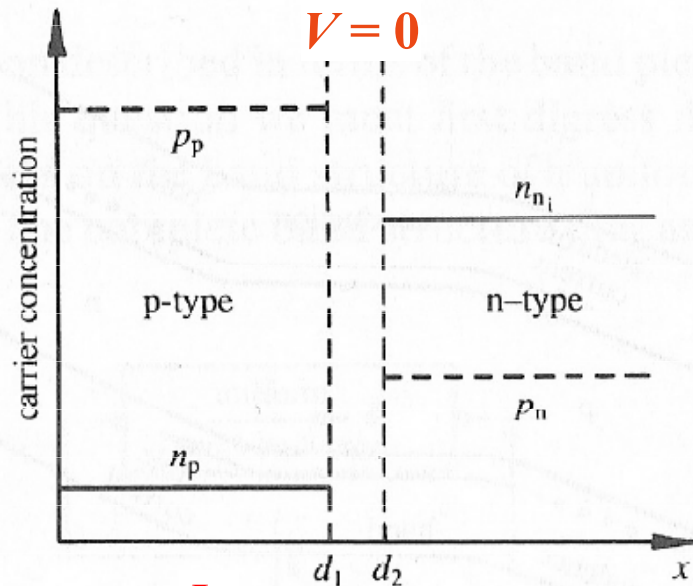
$$J_{en} \propto n_{p0} = n_p \exp(eV/kT)$$



$$J_{en} = J_{ep} \exp(eV/kT)$$

Pn junction with forward bias

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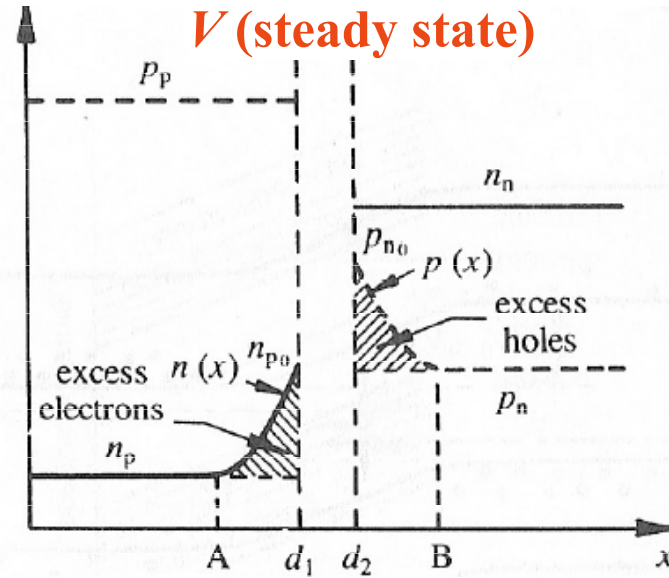


$$J_{hp} \rightarrow \quad \leftarrow J_{en}$$

$$J_{en} \rightarrow \quad \leftarrow J_{hn}$$

$$J_{hp} = J_{hn}$$

$$J_{en} = J_{ep}$$



$$J_{hp} \rightarrow \quad \leftarrow J_{en}$$

$$J_{ep} \rightarrow \quad \leftarrow J_{hn}$$

Total currents:

$$J_h = J_{hp} - J_{hn} = J_{hn} [\exp(eV/kT) - 1]$$

$$J_e = J_{en} - J_{ep} = J_{ep} [\exp(eV/kT) - 1]$$

$$\Rightarrow J = J_0 [\exp(eV/kT) - 1] \quad \text{Diode equation}$$

$$J_0 = J_{hn} + J_{ep}$$

Conclusions

- Extrinsic sC are created by adding impurities:
 - Donors (extra electrons) and acceptors (extra holes).
- Electron processes in sC were qualitatively studied.
- Components of the total current were studied.
- We have deduced the continuity equation for MINORITY carriers.
- Several measurements in sC were studied.
- The junction diode was studied.