

# Physics of Electronics:

## 2. The Electronic Structure of Atoms (cont.)

July – December 2008

# Contents overview

- Interpretation of wave function
- Uncertainty principle
- Beams of particles and potential barriers
- A particle in a 1D potential well
- The hydrogen atom
- The exclusion principle

# Interpretation of the Wave Function

- Born interpretation:
  - The probability of finding the particle in the space volume  $dV$ , at the time  $t$ , is given by:

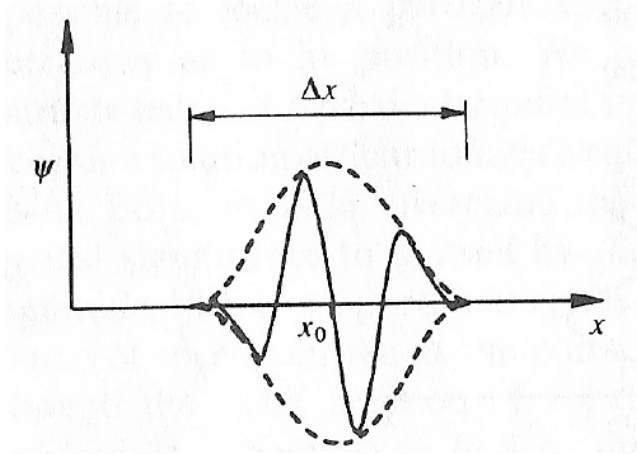
$$|\psi(x, y, z, t)|^2 dV \Rightarrow \int_{\text{whole space}} |\psi(x, y, z, t)|^2 dV = 1$$

(normalization)

- $\psi(x, t)$  also has to be:
  - Continuous and single valued on  $x$ .
  - Idem with its spatial first derivatives.

# Heisenberg Principle

- A rigorous demonstration using matrix mechanics



$$\Delta p \Delta x \geq \hbar/2$$

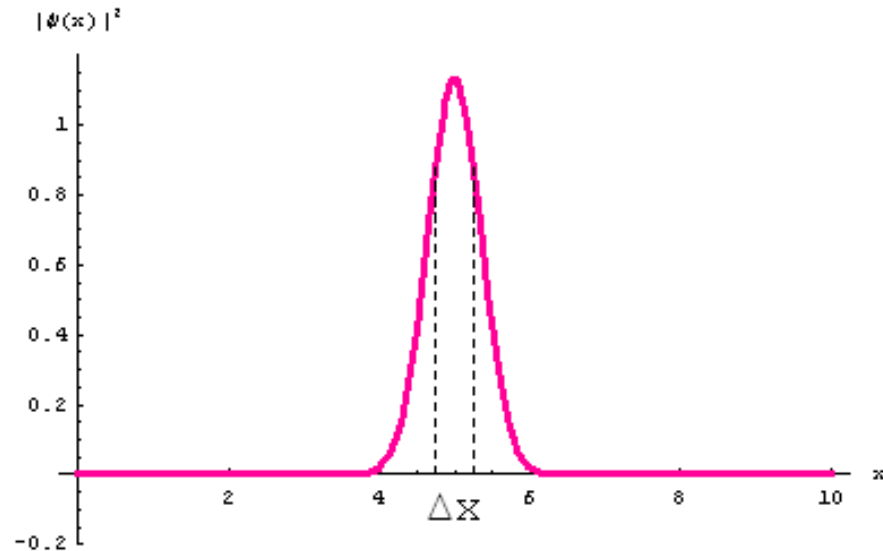
- In QM there are pair of physical quantities (called conjugates) for which this relation holds, e.g.:

$$\Delta E \Delta t \geq \hbar$$

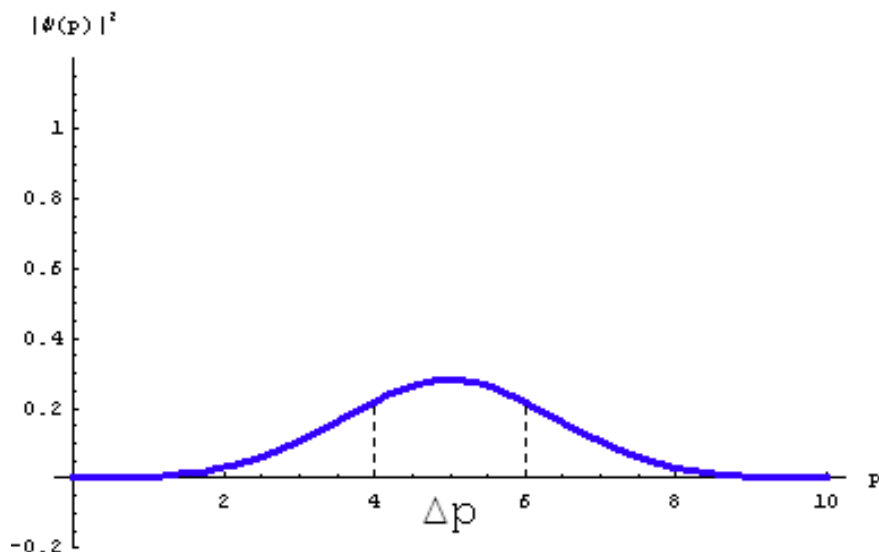
For an experimental demo: *Nature* **371**, 594 - 595 (13 October 2002)

# Heisenberg Principle

- What does exactly  $\Delta q$  mean?

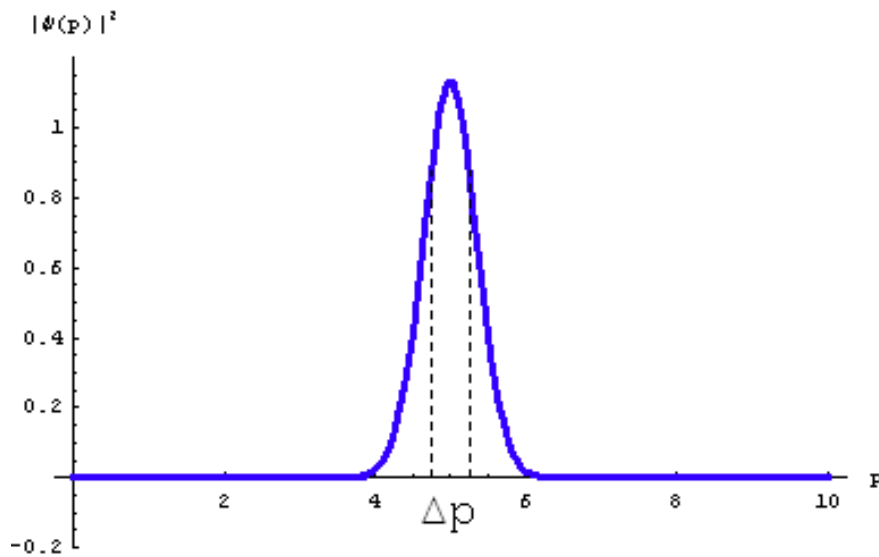
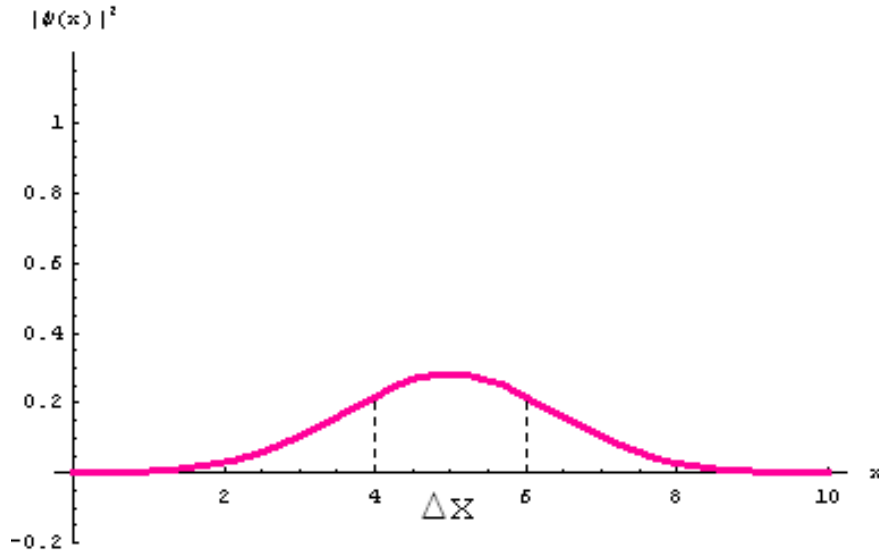


$\Delta q$  is the standard deviation of the probability density



# Heisenberg Principle

- What does exactly  $\Delta q$  mean?

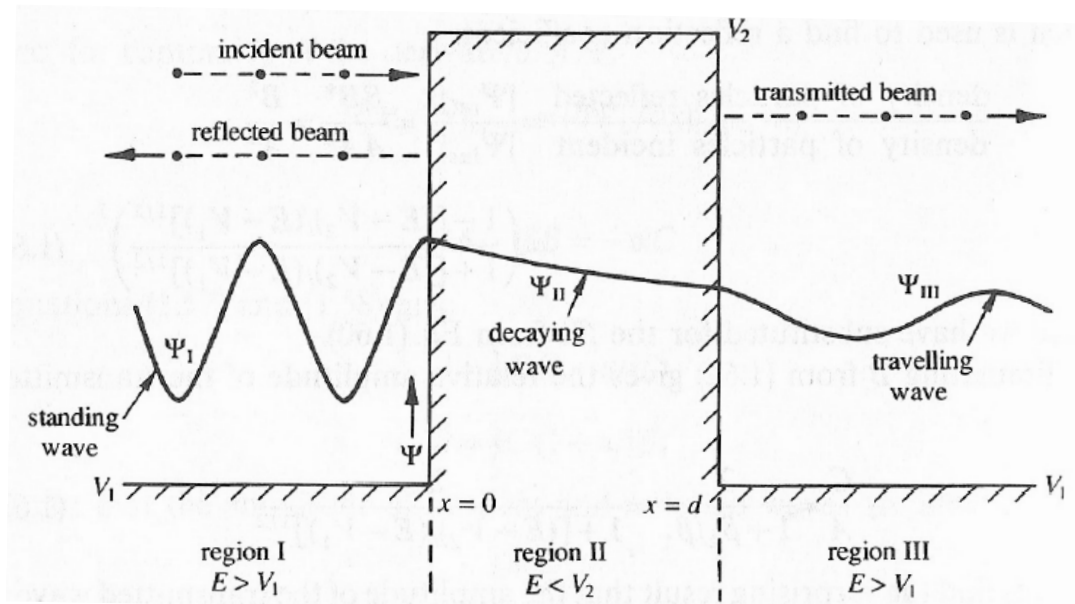


$\Delta q$  is the standard deviation of the probability density

The Heisenberg principle applies to quantities whose standard deviations are not independent. It happens when those quantities are FT related.

# Potential Barriers

- Narrow potential barrier ( $V_1 < E < V_2$ ):



SE is written for every region, from which:

$$\Psi_I = A \exp(j\beta x) + B \exp(-j\beta x)$$

$$\Psi_{II} = C \exp(-\alpha x) + D \exp(\alpha x)$$

$$\Psi_{III} = F \exp(j\beta x)$$

where  $\alpha^2 = \frac{2m}{\hbar^2}(V_2 - E)$

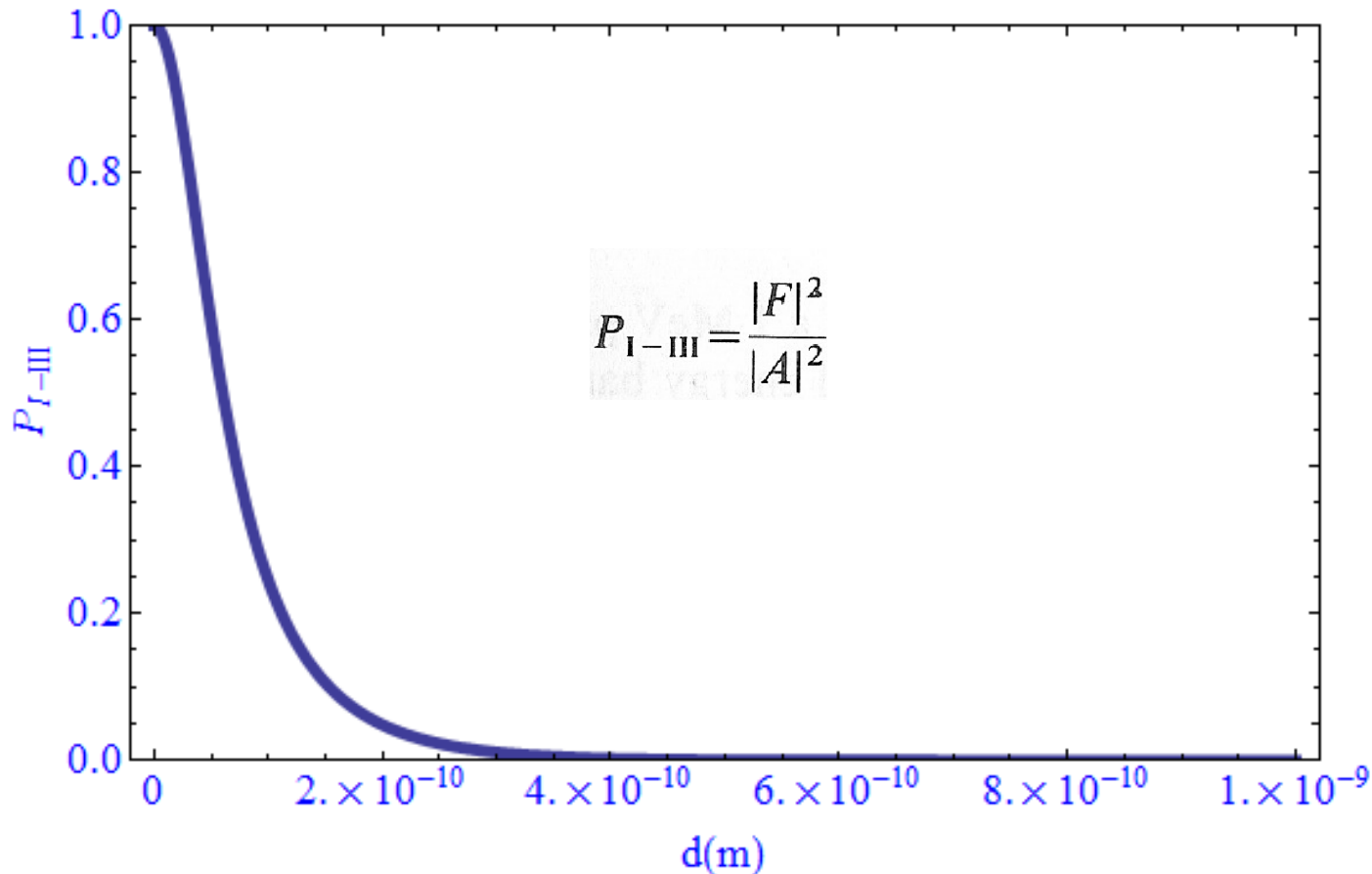
$$\beta^2 = \frac{2m}{\hbar^2}(E - V_1)$$

Constants  $A, B, C, D, F$  are obtained from continuity and normalization. In particular:

$$F = A \exp(-j\beta d) [\cosh(\alpha d) + \frac{1}{2}(\alpha/\beta - \beta/\alpha) \sinh(\alpha d)]^{-1}$$

# Potential Barriers

- Narrow potential barrier ( $V_1 < E < V_2$ ):
  - Transmission from I to III



As comparison, Bohr radius = 0.05 nm

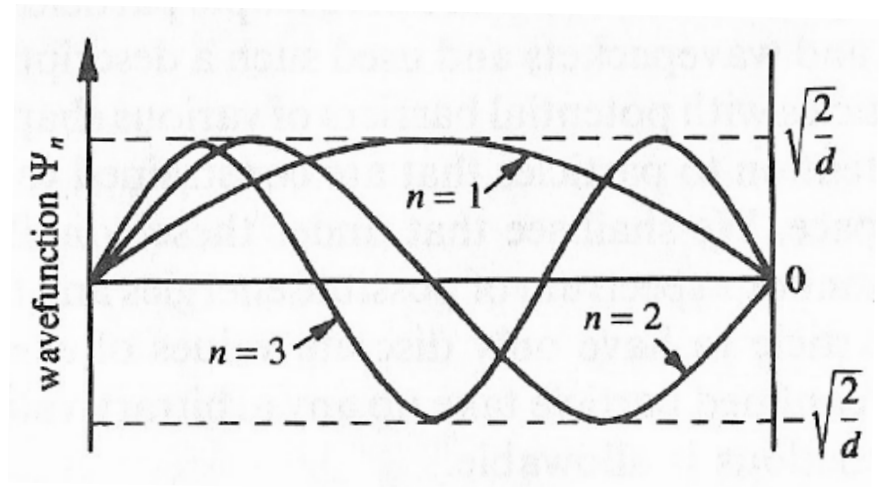
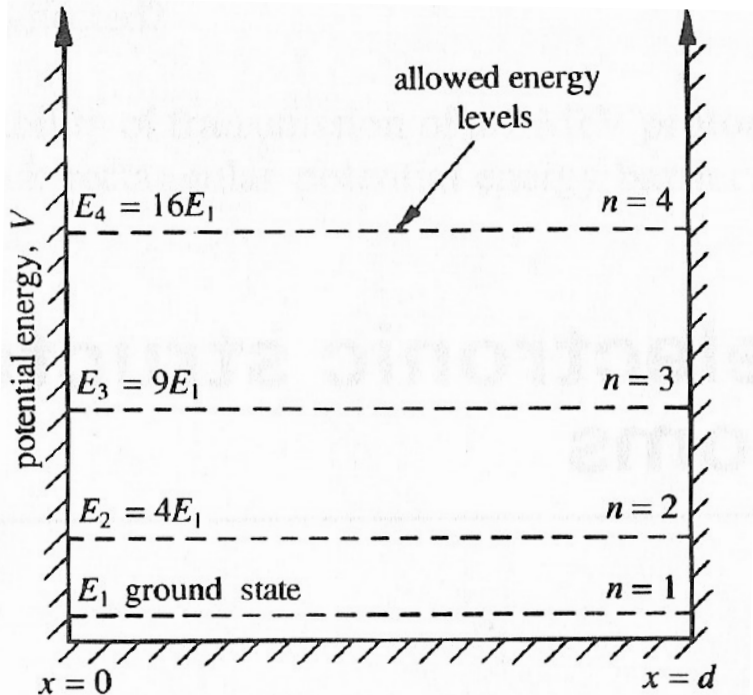


# A Particle in a 1D Potential Well

- Infinite well:

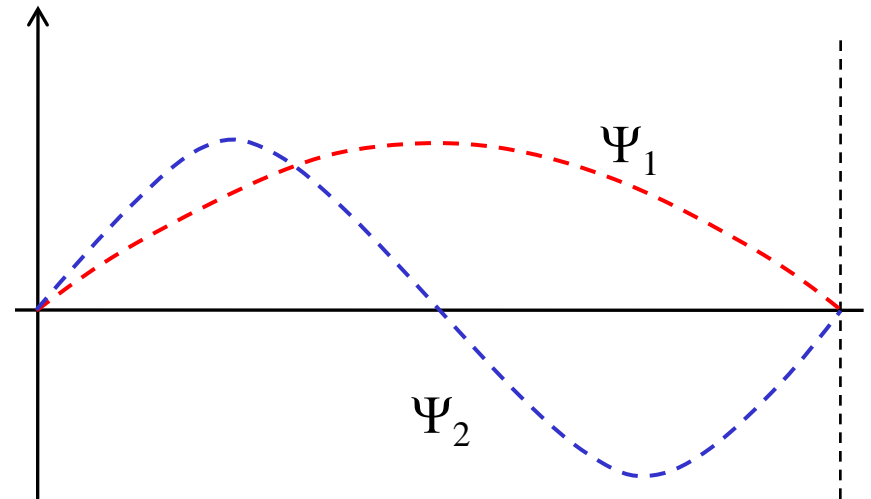
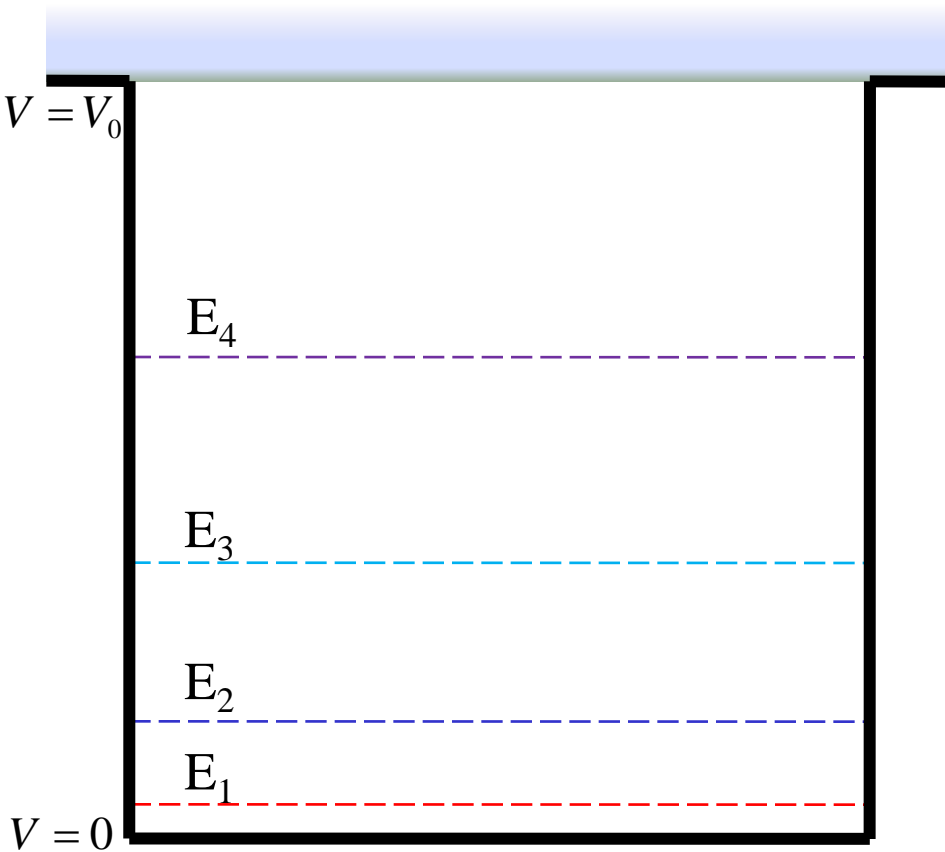
$$V = \begin{cases} 0 & \text{for } 0 < x < d \\ \infty & \text{elsewhere} \end{cases} \xrightarrow{\text{SE}} \frac{d^2\Psi}{dx^2} + \frac{2m}{\hbar^2} E\Psi = 0 \xrightarrow{\beta^2 = (2m/\hbar^2)E} \Psi = Ae^{j\beta x} + Be^{-j\beta x}$$

$$\xrightarrow{\quad} \Psi = (2/d)^{1/2} \sin(n\pi x/d) \quad E = \frac{\hbar^2 n^2 \pi^2}{2md^2} = \frac{n^2 \hbar^2}{8md^2} \quad n = 1, 2, 3, \dots$$



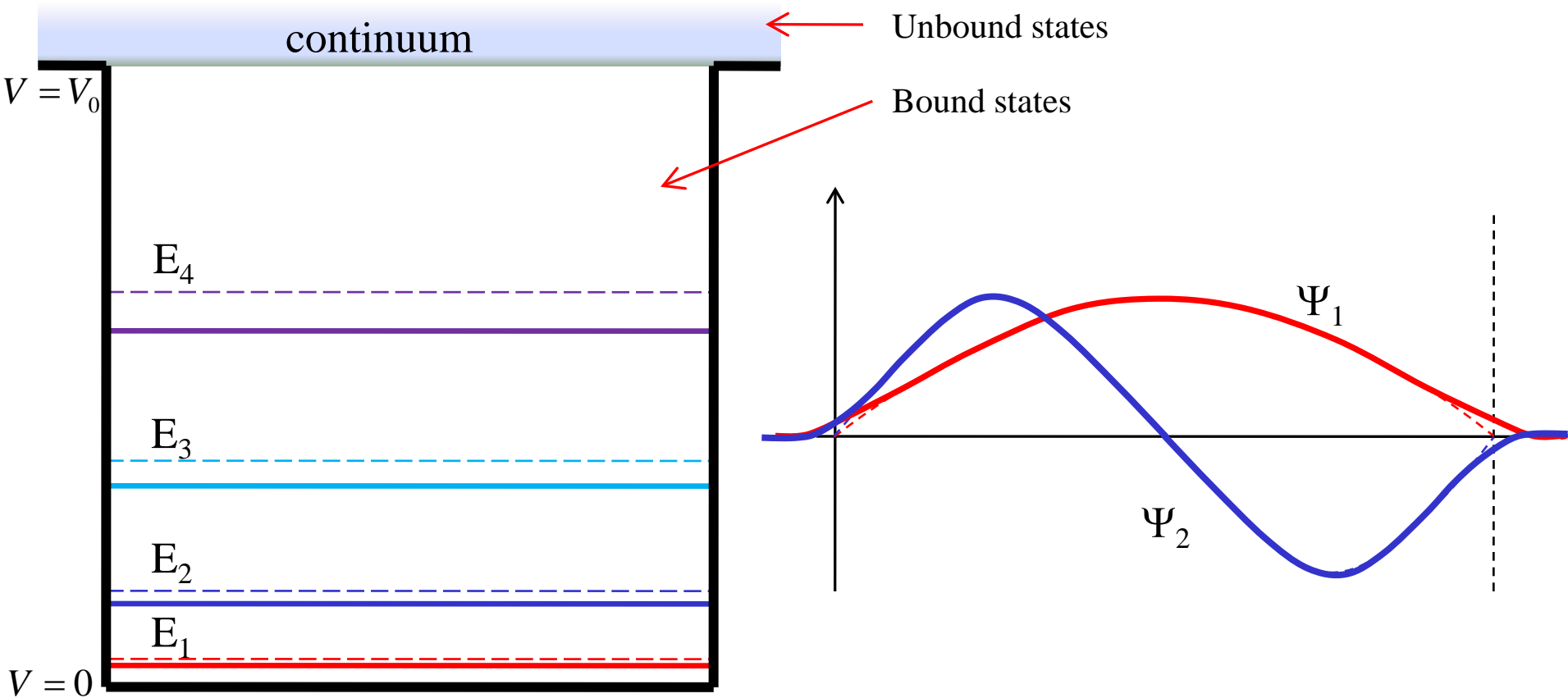
# A Particle in a 1D Potential Well

- Finite well:



# A Particle in a 1D Potential Well

- Finite well (left as exercise):



# The Hydrogen Atom

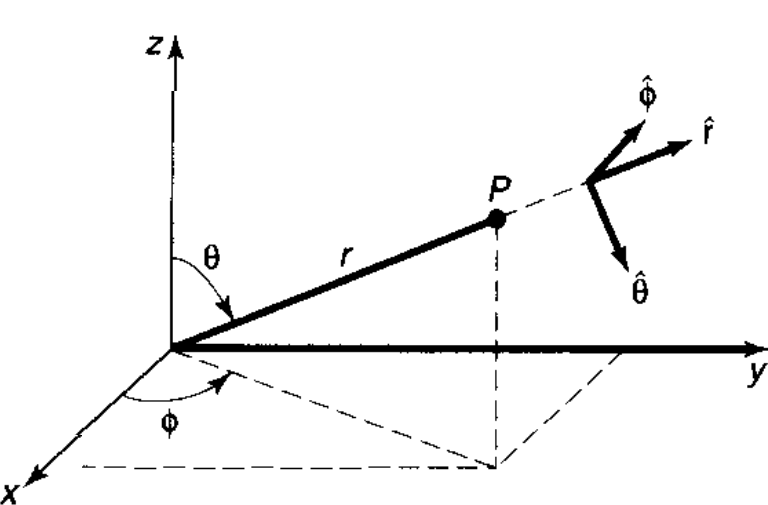
- Electron moving in the electric field of the nucleus:

$$V = -e^2/(4\pi\epsilon_0 r)$$

- We have to solve the time-independent SE:

$$\nabla^2\Psi + \frac{2m}{\hbar^2}(E - V)\Psi = 0$$

- Given the symmetry, we use it in polar coordinates:



$$\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial \Psi}{\partial r} \right) + \frac{1}{r^2} \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \Psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \Psi}{\partial \phi^2} + \frac{2m}{\hbar^2} (E - V) \Psi = 0$$

# The Hydrogen Atom

- Electron moving in the electric field of the nucleus:

$$V = -e^2/(4\pi\epsilon_0 r)$$

- Full solution by separation of variables assuming:

$$\Psi(r, \theta, \phi) = R(r) \underbrace{\Theta(\theta)\Phi(\phi)}_{Y(\theta, \phi)}$$

- Some algebra (and convenient election of constants) gives:

$$\frac{1}{R} \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right) - \frac{2mr^2}{\hbar^2} [V(r) - E] = l(l+1);$$

$$\left. \begin{aligned} \frac{1}{\Theta} \left[ \sin \theta \frac{d}{d\theta} \left( \sin \theta \frac{d\Theta}{d\theta} \right) \right] + l(l+1) \sin^2 \theta &= m^2; \\ \frac{1}{\Phi} \frac{d^2 \Phi}{d\phi^2} &= -m^2. \end{aligned} \right\} \text{Independent of } V$$

# The Hydrogen Atom

- Electron moving in the electric field of the nucleus:

$$V = -e^2/(4\pi\epsilon_0 r)$$

– Azimuth angle part:

$$\frac{d^2\Phi}{d\phi^2} = -m^2\Phi \Rightarrow \Phi(\phi) = e^{im\phi}, \text{ but } \Phi(\phi + 2\pi) = \Phi(\phi). \quad \Rightarrow \quad m = 0, \pm 1, \pm 2, \dots$$

– Polar angle part:

$$\frac{1}{\Theta} \left[ \sin\theta \frac{d}{d\theta} \left( \sin\theta \frac{d\Theta}{d\theta} \right) \right] + l(l+1) \sin^2\theta = m^2; \quad \Rightarrow \quad \Theta(\theta) = A P_l^m(\cos\theta),$$

where:

$$P_l^m(x) \equiv (1-x^2)^{|m|/2} \left( \frac{d}{dx} \right)^{|m|} P_l(x), \quad \Rightarrow \quad |m| \leq l \quad \Rightarrow \quad m = -l, -l+1, \dots, -1, 0, 1, \dots, l-1, l.$$

$$P_l(x) \equiv \frac{1}{2^l l!} \left( \frac{d}{dx} \right)^l (x^2-1)^l. \quad \Rightarrow \quad l = 0, 1, 2, \dots$$

# The Hydrogen Atom

- Electron moving in the electric field of the nucleus:

$$V = -e^2/(4\pi\epsilon_0 r)$$

- (Total) angular part depends on two integer numbers:

$$Y(\theta, \phi) = \Theta(\theta)\Phi(\phi) = Y_l^m(\theta, \phi)$$

- After normalization (R and Y are normalized independently):

$$Y_l^m(\theta, \phi) = \epsilon \sqrt{\frac{(2l+1)}{4\pi} \frac{(l-|m|)!}{(l+|m|)!}} e^{im\phi} P_l^m(\cos\theta),$$

where  $\epsilon = (-1)^m$  for  $m \geq 0$  and  $\epsilon = 1$  for  $m \leq 0$ .

# The Hydrogen Atom

- Electron moving in the electric field of the nucleus:

$$V = -e^2/(4\pi\epsilon_0 r)$$

– Radial part

$$\frac{1}{R} \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right) - \frac{2mr^2}{\hbar^2} [V(r) - E] = l(l+1);$$

whose solution is (see any book on Q.M.):

$$R_{nl}(r) = \frac{1}{r} \rho^{l+1} e^{-\rho} v(\rho), \quad \text{with} \quad \rho = \frac{r}{an}. \quad (a \text{ is the Bohr radius})$$

where:

$$\left. \begin{aligned} v(\rho) &= L_{n-l-1}^{2l+1}(2\rho), \\ L_{q-p}^p(x) &\equiv (-1)^p \left( \frac{d}{dx} \right)^p L_q(x) \\ L_q(x) &\equiv e^x \left( \frac{d}{dx} \right)^q (e^{-x} x^q) \end{aligned} \right\} \Rightarrow l = 0, 1, 2, \dots, n-1.$$



# The Hydrogen Atom

- Electron moving in the electric field of the nucleus:

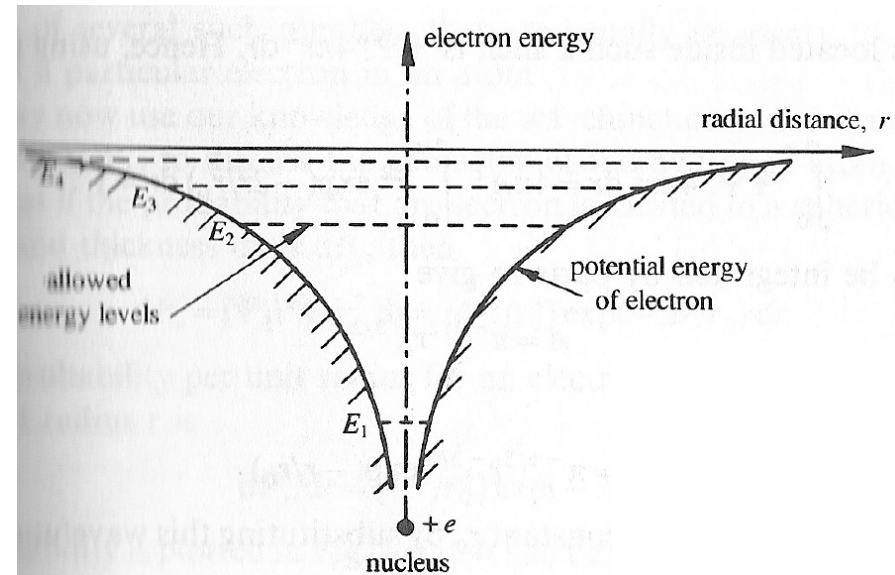
$$V = -e^2/(4\pi\epsilon_0 r)$$

– Radial part

$$\frac{1}{R} \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right) - \frac{2mr^2}{\hbar^2} [V(r) - E] = l(l+1);$$

whose solution is (see any book on Q.M.):

$$E_n = - \left[ \frac{m}{2\hbar^2} \left( \frac{e^2}{4\pi\epsilon_0} \right)^2 \right] \frac{1}{n^2} = \frac{E_1}{n^2}, \quad n = 1, 2, 3, \dots$$



# The Hydrogen Atom

- Electron moving in the electric field of the nucleus:

$$V = -e^2/(4\pi\epsilon_0 r)$$

- Total solution:

$$\psi_{nlm}(r, \theta, \phi) = R_{nl}(r) Y_l^m(\theta, \phi),$$

- After normalization:

$$\psi_{nlm} = \sqrt{\left(\frac{2}{na}\right)^3 \frac{(n-l-1)!}{2n[(n+l)!]^3}} e^{-r/na} \left(\frac{2r}{na}\right)^l L_{n-l-1}^{2l+1}\left(\frac{2r}{na}\right) Y_l^m(\theta, \phi).$$

- Notice that the w.f. depends on 3 integers (or quantum numbers) but the energy only in one of them. I.e. several w.f.'s can have the same value of energy. This is called DEGENERACY.

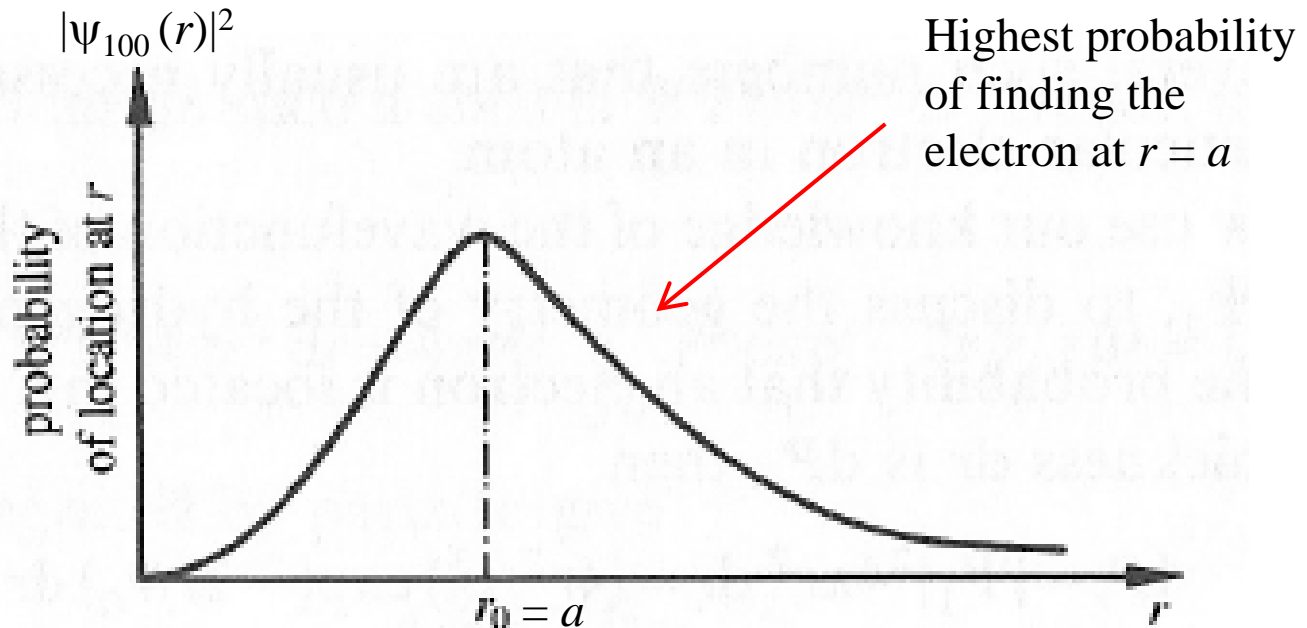
# The Hydrogen Atom

- Electron moving in the electric field of the nucleus:

$$V = -e^2/(4\pi\epsilon_0 r)$$

- The ground state (level with lowest energy) is:

$$\psi_{100}(r, \theta, \phi) = \frac{1}{\sqrt{\pi a^3}} e^{-r/a}.$$



# Spin & Exclusion Principle

- It exist another quantum number describing the electron in an atom: SPIN number. It represents the rotation of the electron over its own axis. It can be obtained by solving the relativistic SE.
- Therefore, an electron in an atom is completely specified by 4 quantum numbers:

$$(n, l, m, s).$$

- The exclusion principle says that no 2 electrons can have the same quantum numbers.

# Exclusion Principle

- This applies for any atom, which permits to construct the following table:

Table 2.1 Electronic structure of the lighter elements.

Element	Principal quantum number, $n$	Azimuthal quantum number, $l=0, 1, \dots, n-1$	Magnetic quantum number, $m=-l, \dots, +l$	Spectroscopic designation
H	1	0	0	1s
He	1	0	0	1s <sup>2</sup>
Li	2	0	0	1s <sup>2</sup> 2s
Be	2	0	0	1s <sup>2</sup> 2s <sup>2</sup>
B	2	1	-1	1s <sup>2</sup> 2s <sup>2</sup> 2p
C	2	1	-1	1s <sup>2</sup> 2s <sup>2</sup> 2p <sup>2</sup>
N	2	1	0	1s <sup>2</sup> 2s <sup>2</sup> 2p <sup>3</sup>
O	2	1	0	1s <sup>2</sup> 2s <sup>2</sup> 2p <sup>4</sup>
F	2	1	1	1s <sup>2</sup> 2s <sup>2</sup> 2p <sup>5</sup>
Ne	2	1	1	1s <sup>2</sup> 2s <sup>2</sup> 2p <sup>6</sup>
Na	3	0	0	1s <sup>2</sup> 2s <sup>2</sup> 2p <sup>6</sup> 3s
Mg	3	0	0	1s <sup>2</sup> 2s <sup>2</sup> 2p <sup>6</sup> 3s <sup>2</sup>
Al	3	1	-1	1s <sup>2</sup> 2s <sup>2</sup> 2p <sup>6</sup> 3s <sup>2</sup> 3p
Si	3	1	-1	1s <sup>2</sup> 2s <sup>2</sup> 2p <sup>6</sup> 3s <sup>2</sup> 3p <sup>2</sup>
P	3	1	0	1s <sup>2</sup> 2s <sup>2</sup> 2p <sup>6</sup> 3s <sup>2</sup> 3p <sup>3</sup>
etc.				

where:

$l$	State or subshell
0	s
1	p
2	d
3	f
etc.	

# Conclusions

- We have made a model of the atom:
  - One approximated (well)
  - One more realistic (hydrogen atom).
- The solution of the H atom gave rise to 3 quantum numbers.
- The complete description of an electronic state requires an extra number: spin.
- Exclusion principle. The set of 4 quantum numbers has to be different for every electron in an atom.