

Physics of Electronics:

7. Junction Diodes

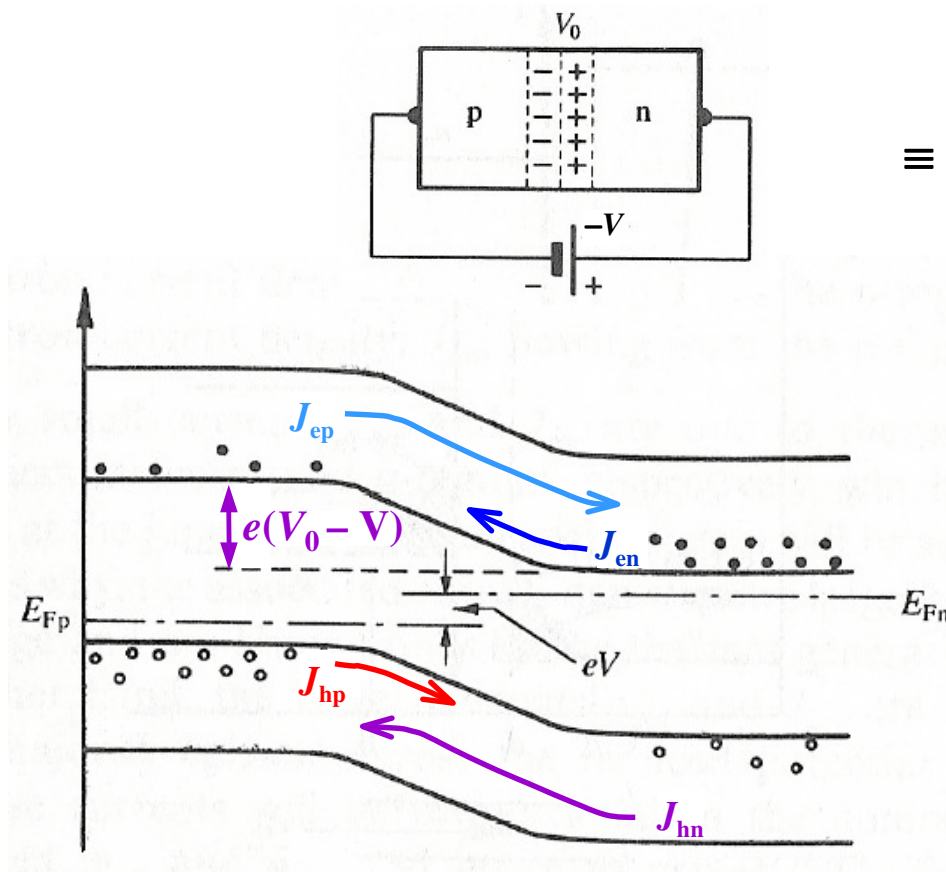
July – December 2008

Contents overview

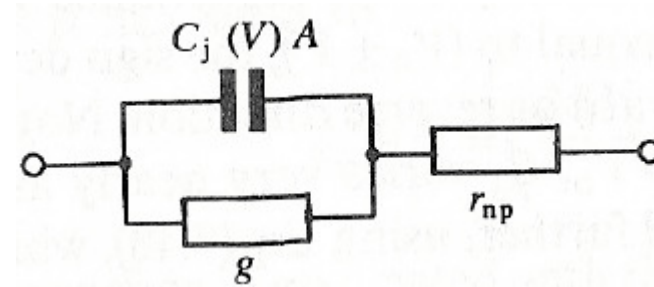
- Small-signal equivalent circuit.
- Switching characteristics.
- Metal-semiconductor junctions.
- The Zener diode.
- The tunnel diode

Small-AC-signal equivalent circuit

- Reversed biased junction



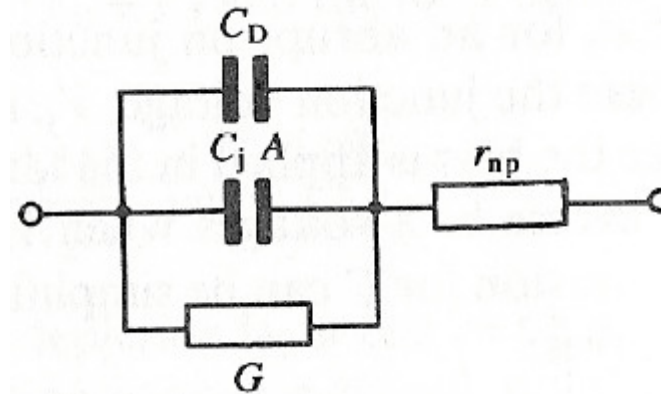
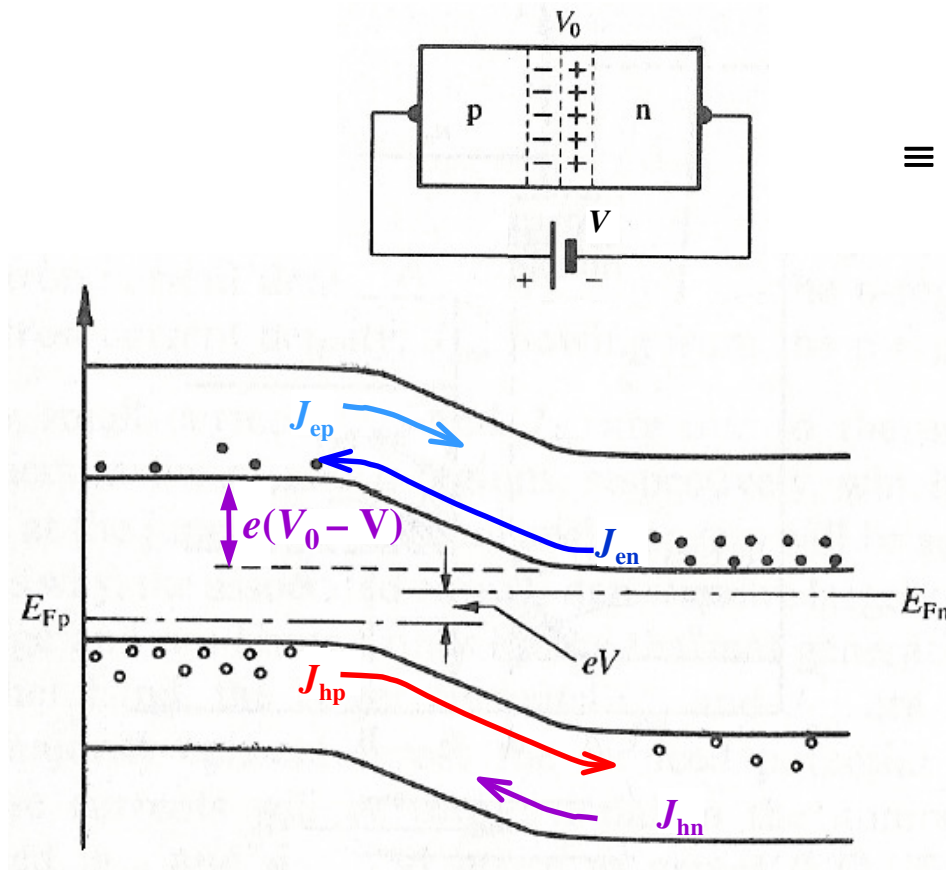
$\equiv$



r_{np} : Resistance of the junction @ $V=0$
 $C_j(V)$: Depletion layer capacitance
 g : Conductance of the reverse flow

Small-AC-signal equivalent circuit

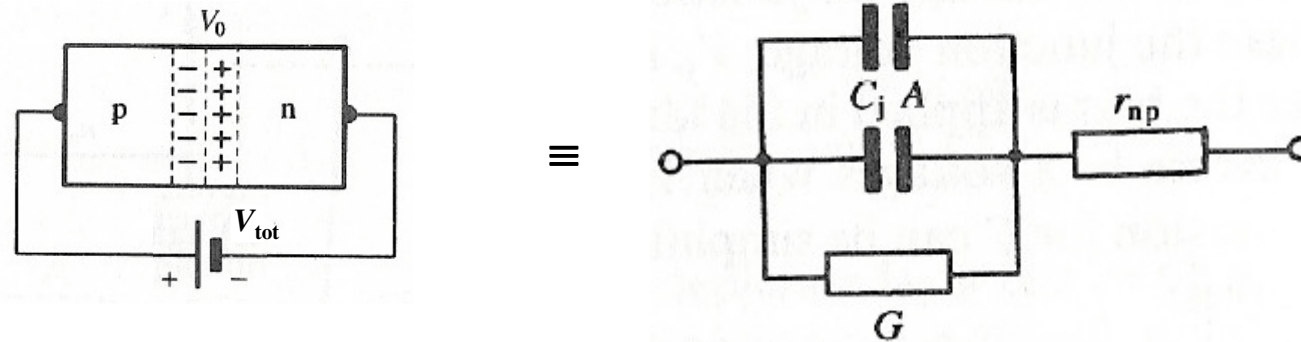
- Forward biased junction



- r_{np} : Resistance of the junction @ $V=0$
- $C_j(V)$: Depletion layer capacitance
- G : Conductance of the forward flow
- C_D : Diffusion capacitance

Small-AC-signal equivalent circuit

- Forward biased junction



$$\left. \begin{aligned} V_{tot} &= V + V_1 \exp(i\omega t) \\ p_{n0} &= p_n \exp(eV/kT) \end{aligned} \right\} \Rightarrow p_{n0} = p_n \exp\left(\frac{e}{kT}[V + V_1 \exp(j\omega t)]\right)$$

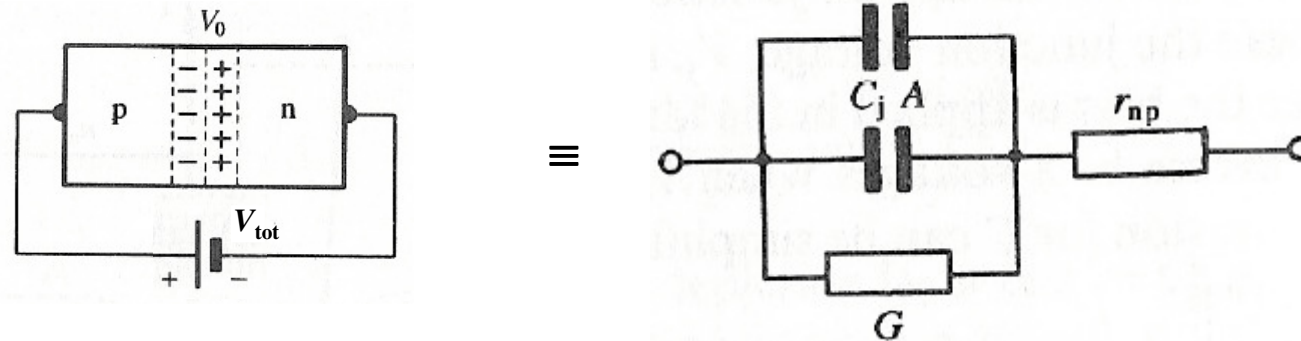
$$V_1 \ll V_0 \Rightarrow p_{n0} = \underbrace{p_n \exp\left(\frac{eV}{kT}\right)}_{\text{DC}} \underbrace{\left(1 + \frac{eV_1}{kT} \exp(j\omega t)\right)}_{\text{AC}}$$

- Therefore we can assume:

$$\delta p = p(x, t) - p_n = \underbrace{p_0(x)}_{\text{DC}} + \underbrace{p_1(x) \exp(j\omega t)}_{\text{AC}} - p_n$$

Small-AC-signal equivalent circuit

- Forward biased junction



- Applying continuity equation:

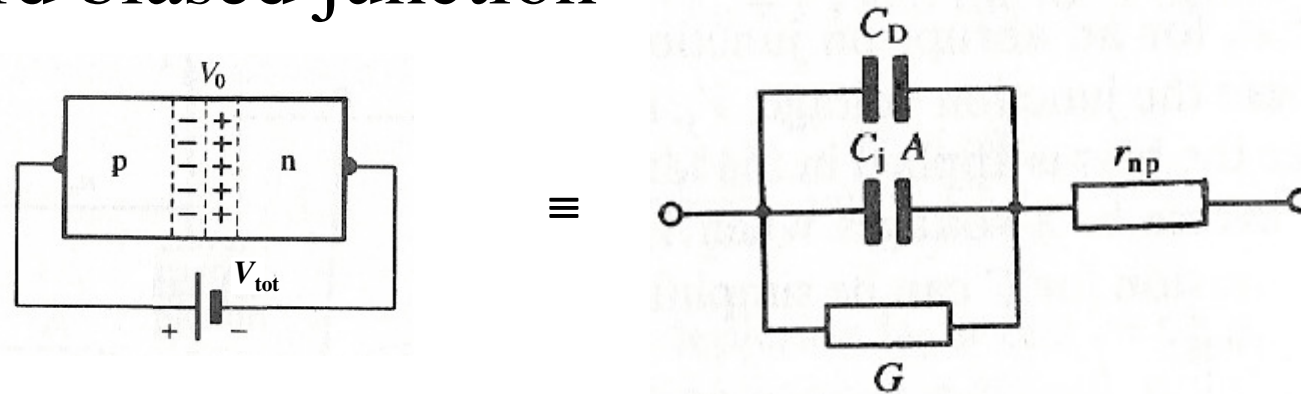
$$\frac{\partial(\delta p)}{\partial t} = -\frac{\delta p}{\tau_{Lh}} + \mu_h \cancel{E_x}^0 \frac{\partial(\delta p)}{\partial x} + D_h \frac{\partial^2(\delta p)}{\partial x^2} \quad \Rightarrow \quad j\omega p_1 = \frac{-p_1}{\tau_{Lh}} + D_h \frac{\partial^2 p_1}{\partial x^2}$$

$$\Rightarrow \frac{\partial^2 p_1}{\partial x^2} = \frac{1 + j\omega\tau_{Lh}}{L_h^2} p_1 \quad \Rightarrow \quad p_1 = C \exp\left(-\frac{(1 + j\omega\tau_{Lh})^{1/2} x}{L_h}\right)$$

- But $p_1(x=0) = p_{n0} \Rightarrow p_1 = p_n \left(\frac{eV_1}{kT}\right) \exp\left(\frac{eV}{kT}\right) \exp\left(-\frac{(1 + j\omega\tau_{Lh})^{1/2} x}{L_h}\right) \exp(j\omega t)$

Small-AC-signal equivalent circuit

- Forward biased junction



- Alternating currents of holes injected into n-region:

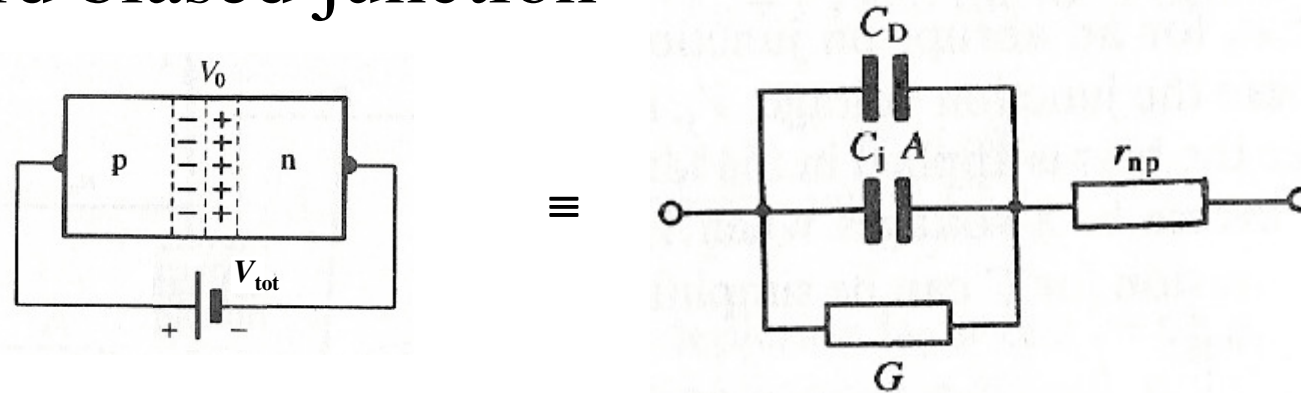
$$J_{h1} = -eD_h \left(\frac{dp(x)}{dx} \right) \Big|_{x=0} = \frac{(1 + i \omega \tau_{Lh})^{1/2}}{L_h} \frac{p_n D_h e^2 V_1}{kT} \exp\left(\frac{eV}{kT}\right) \exp(j\omega t)$$

- Alternating currents of electrons injected into p-region:

$$J_{e1} = -eD_e \left(\frac{dn(x)}{dx} \right) \Big|_{x=0} = \frac{(1 + i \omega \tau_{Le})^{1/2}}{L_e} \frac{n_p D_e e^2 V_1}{kT} \exp\left(\frac{eV}{kT}\right) \exp(j\omega t)$$

Small-AC-signal equivalent circuit

- Forward biased junction



- Total alternating currents of minority injected carriers:

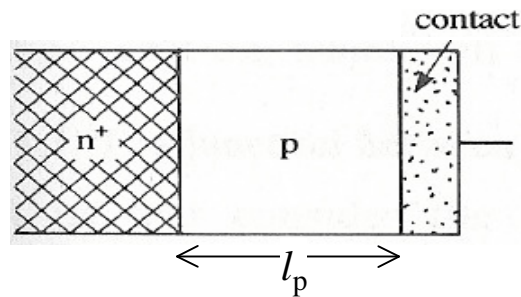
$$J_1 = J_{h1} + J_{e1}$$

- Admittance ($Y_1 = J_1/V_1$) for $\omega\tau \ll 1$:

$$Y_1 \cong \underbrace{\frac{e^2}{kT} \exp\left(\frac{eV}{kT}\right) \left(\frac{D_h p_n}{L_h} + \frac{D_e n_p}{L_e}\right)}_G + \underbrace{\frac{j\omega e^2}{2kT} \left(\frac{D_h p_n \tau_{Lh}}{L_h} + \frac{D_e n_p \tau_{Le}}{L_e}\right)}_{C_D}$$

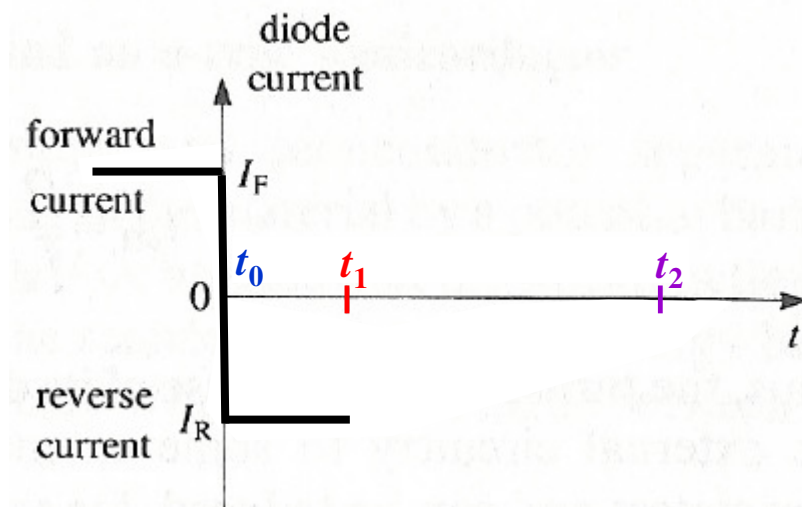
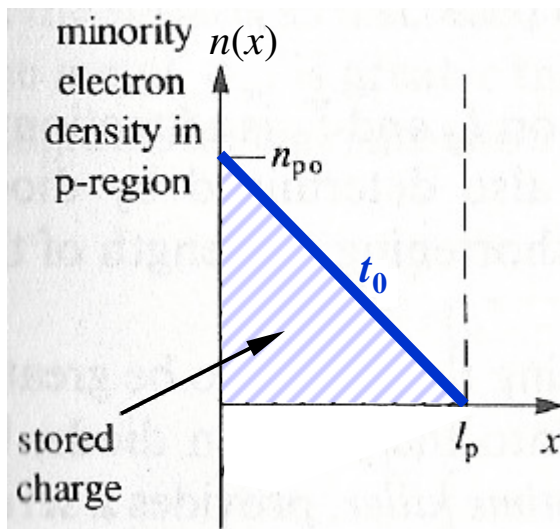
Switching Characteristics

- Consider a forward biased np junction



$$n(x) = \left(\frac{n_p \exp(-l_p/L_e) + n_{p0} - n_p}{1 - \exp(-2l_n/L_h)} \right) \exp(-x/L_e) - \left(\frac{n_p \exp(-l_p/L_e) + n_{p0} - n_p}{1 - \exp(-2l_p/L_e)} - (n_{p0} - n_p) \right) \exp(+x/L_e) + n_p$$

- With $l_p \leq L_e$, what if the junction is suddenly reversed?



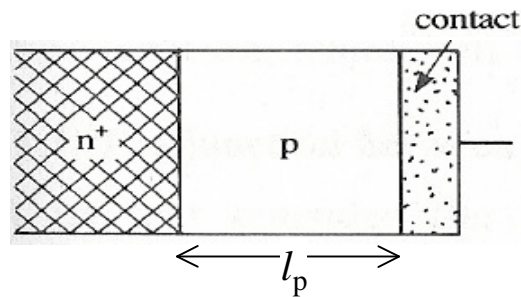
$t < t_0$: junction forward biased, $n(x)$ is linear.

$t = t_0$: junction reversed biased.

$t < t_1$: $n(x)$ starts decreasing but $I_R \cong I_F$.

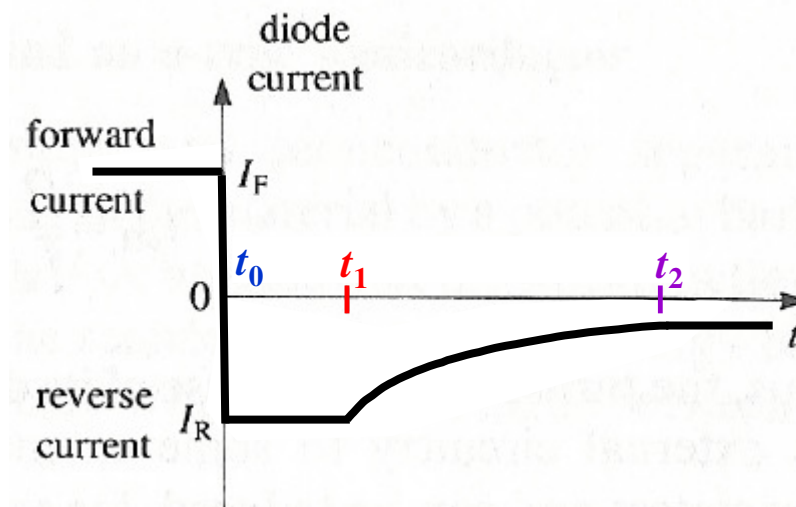
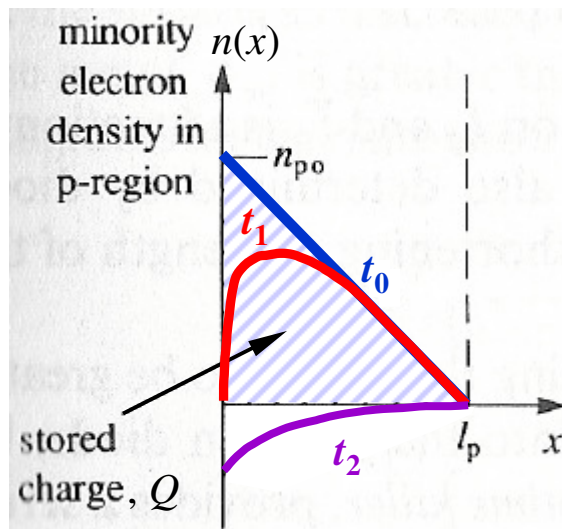
Switching Characteristics

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$$n(x) = \left(\frac{n_p \exp(-l_p/L_e) + n_{p0} - n_p}{1 - \exp(-2l_n/L_h)} \right) \exp(-x/L_e) - \left(\frac{n_p \exp(-l_p/L_e) + n_{p0} - n_p}{1 - \exp(-2l_p/L_e)} - (n_{p0} - n_p) \right) \exp(+x/L_e) + n_p$$

- With $l_p \leq L_e$, what if the junction is suddenly reversed?



$t = t_1$: $n(x)$ keeps decreasing but I_R cannot be maintained.

$t < t_2$: $n(x)$ and current decrease further.

$t = t_2$: $n(x)$ and current reach the steady reversed values.

Switching Characteristics

- Estimation of the time t_{off} that takes to deplete the stored charge Q :

- Q is given by:

$$Q = A \int_0^{l_p} \rho(x) dx \quad \text{where} \quad \rho(x) = n(x)e = en_{p0}(1 - x/l_p)$$

- $n(x)$ can be expressed in terms of the currents:

$$\left. \begin{array}{l} \text{– Definition of density current} \quad J_{Fe} = ev(x)n(x) \\ \text{– Assuming only diffusion} \quad J_{Fe} = eD_e dn/dx \\ \text{– From the measured } I_F \quad J_{Fe} = I_F/A \end{array} \right\} \Rightarrow \rho(x) = I_F(l_p - x)/AD_e$$

$$\therefore Q = I_F l_p^2 / 2D_e$$

- In average, the current I_R is then:

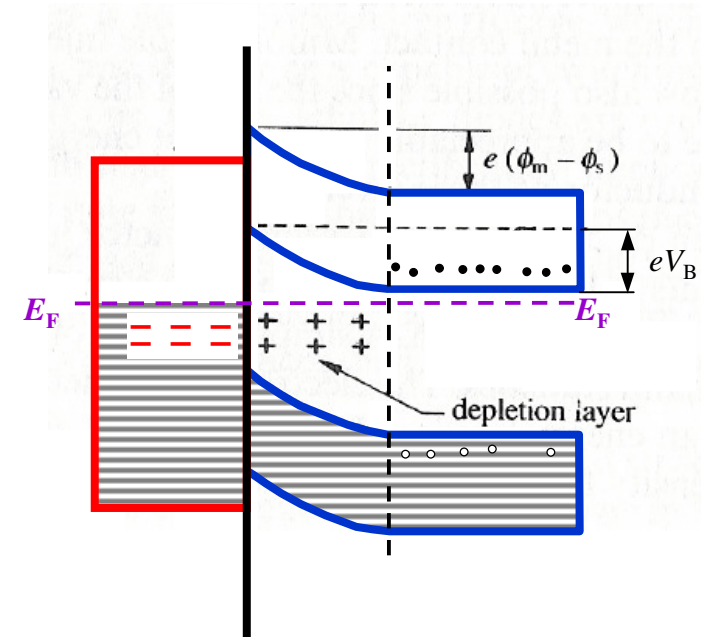
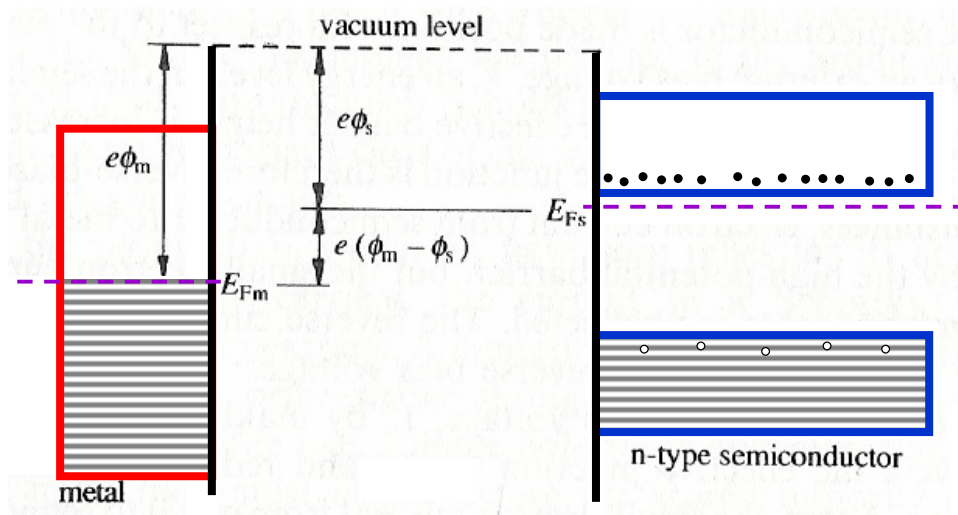
$$\bar{I}_R = Q/t_{\text{off}} \Rightarrow$$

$$t_{\text{off}} = \frac{I_F}{\bar{I}_R} \frac{l_p^2}{2D_e}$$

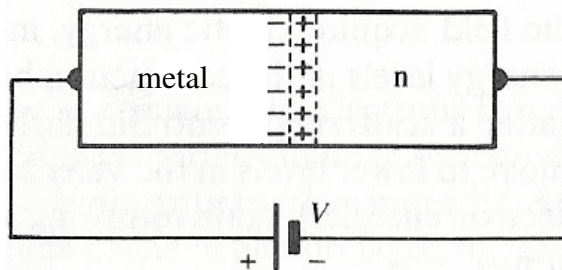
It can be tailored!!!

Metal-sC junctions

- Metal – n-type sC or Schottky junction ($\phi_m > \phi_s$):
 - Unbiased junction:



- Biased junction:

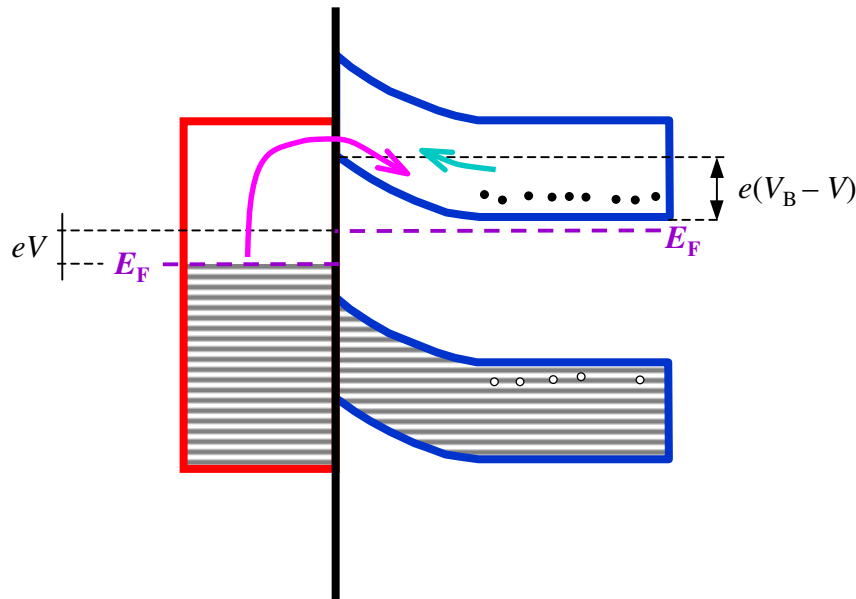


- Forward biased: $V_B \rightarrow V_B - V$ (sC-to-m currents fostered)
- Reversed biased: $V_B \rightarrow V_B + V$ (sC-to-m currents suppressed)

Contact is, therefore, not ohmic!!!!

Metal-sC junctions

- Metal – n-type sC or Schottky junction ($\phi_m > \phi_s$):
 - Currents in a forward biased junction:



$$J_{ms} \propto \exp(-eV_B/kT)$$

$$J_{sm} \propto \exp[-e(V_B - V)/kT]$$

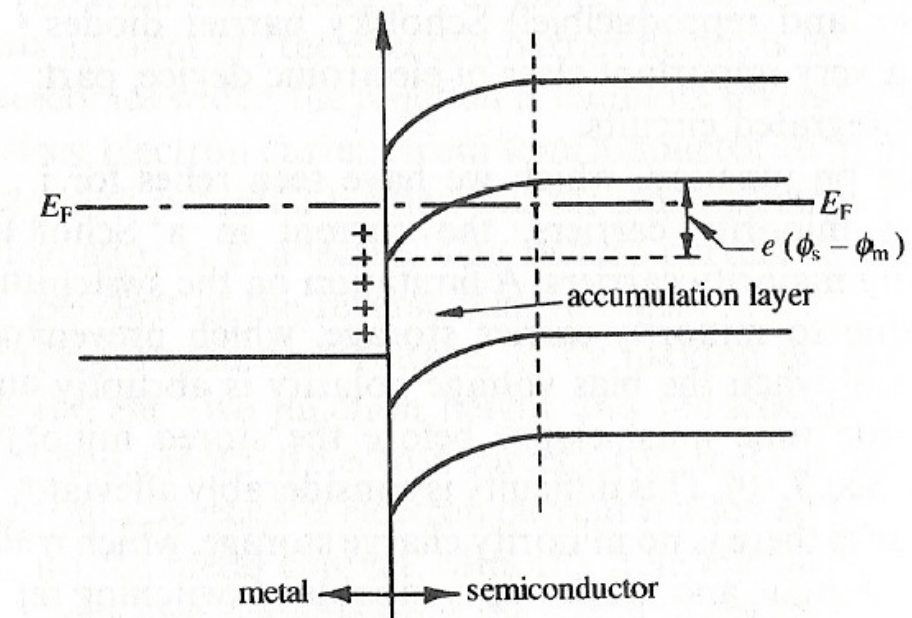
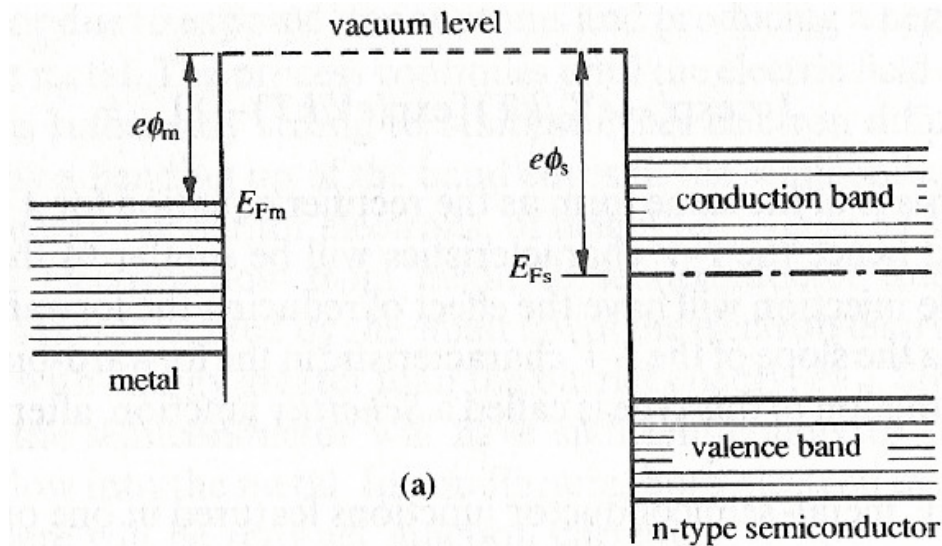
Electron majority current is then:

$$J \propto \exp(-eV_B/kT) [\exp(eV/kT) - 1]$$

- Similar to pn junction but much faster when switching (no minority charge storage)

Metal-sC junctions

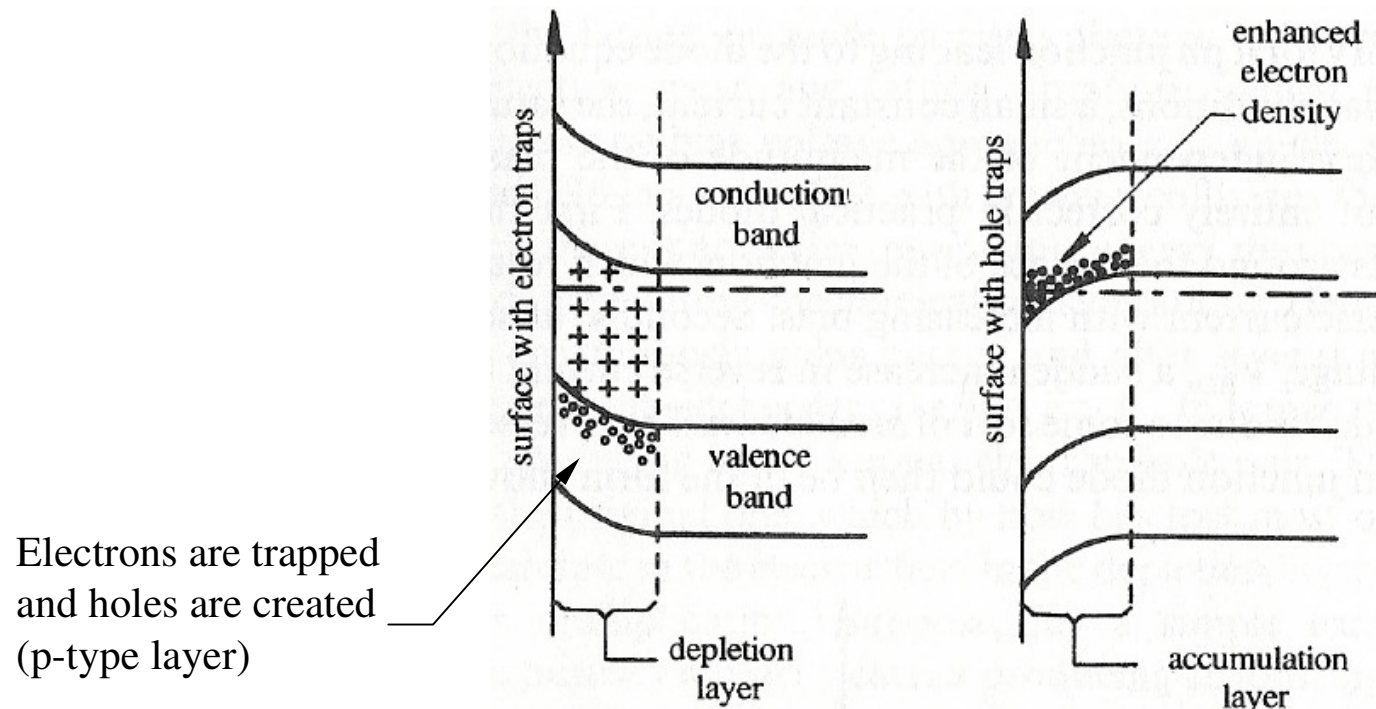
- Metal – n-type sC junction ($\phi_s > \phi_m$):
 - Unbiased junction:



- No potential barrier is formed.
 - Ohmic contact (i.e. no rectification).
- Homework: analyze a metal – p-type sC junction.

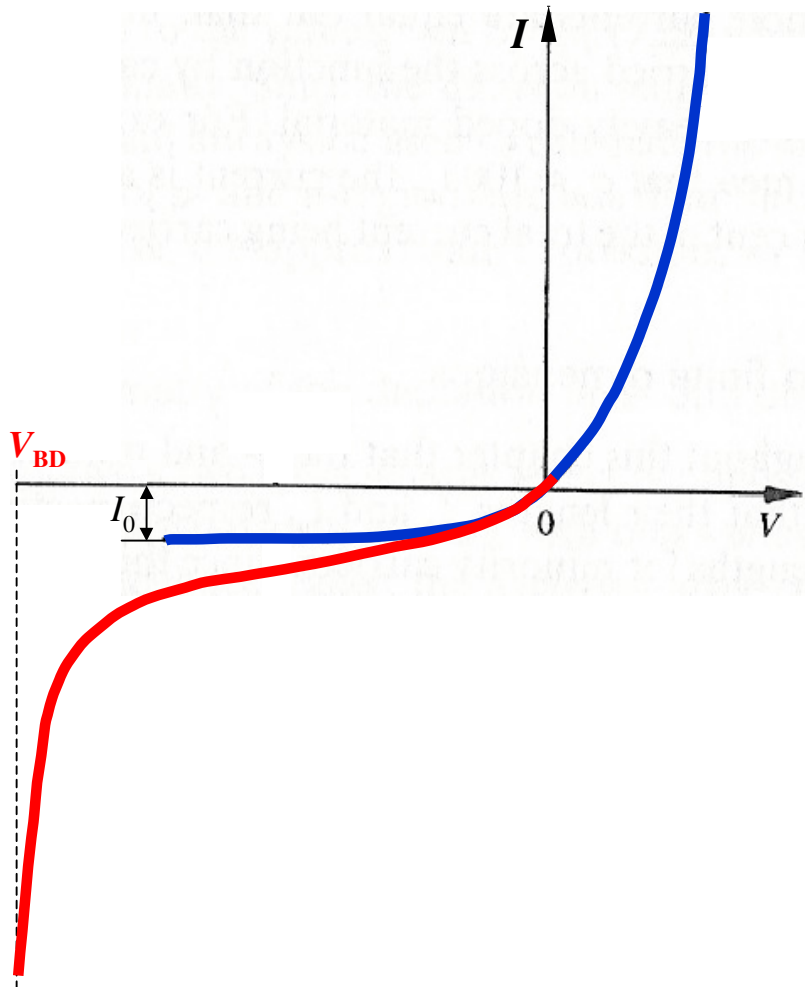
Metal-sC junctions

- Surface states:
 - The surface can be affected by oxidation/absorption of impurity atoms.
 - These states can act as electron/hole traps.
 - Consider an n-type sC:



The Zener Diode

- Breakdown voltage

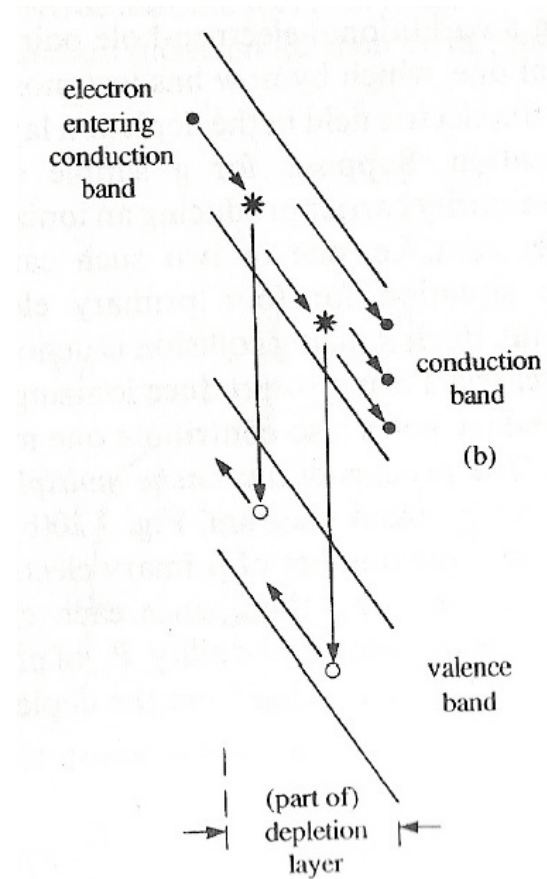
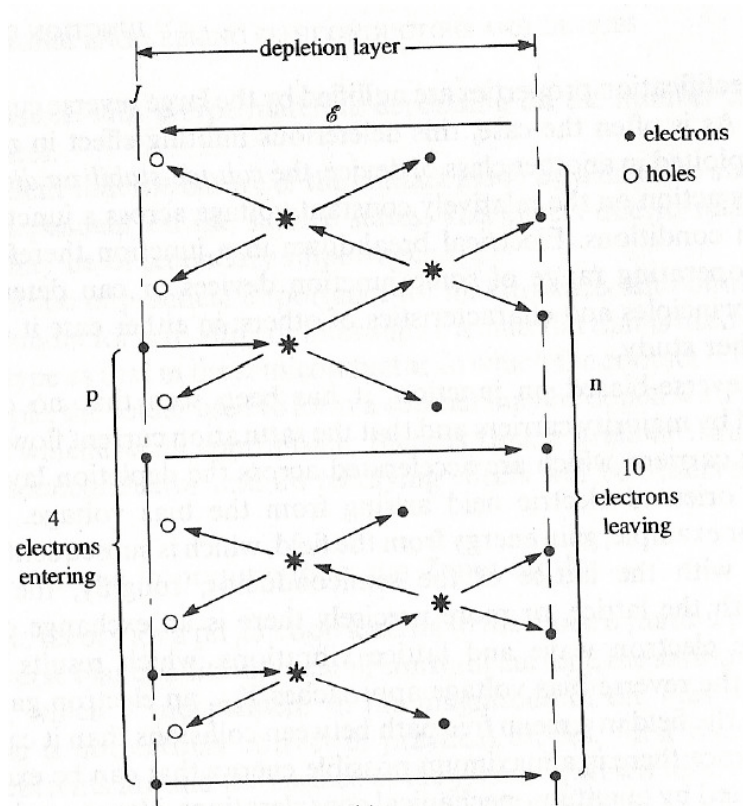


Ideal diode: current limited if reversed biased

Real diode: voltage limited if reversed biased

The Zener Diode

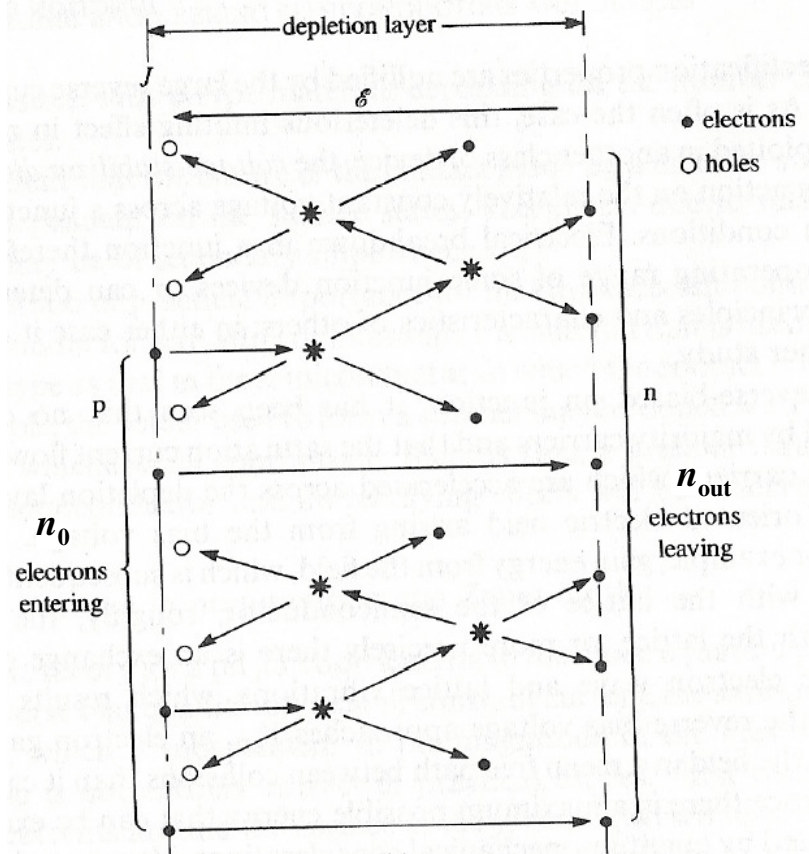
- Avalanche breakdown



- Energy acquired by electrons is limited by collisions.
- Dominant effect in wider junctions and moderate doping.
- V_{BD} is T dependent (because mean free path is T dependent).

The Zener Diode

- Avalanche breakdown



$$n_{out} = n_0 + P_i n_0 + P_i(P_i n_0) + \dots = n_0 \sum_{m=0}^{\infty} P_i^m$$

$\swarrow$ 1st collision $\swarrow$ 2nd collision

- Multiplication factor:

$$M = n_{out}/n_0 = \sum_{m=0}^{\infty} P_i^m = 1/(1 - P_i)$$

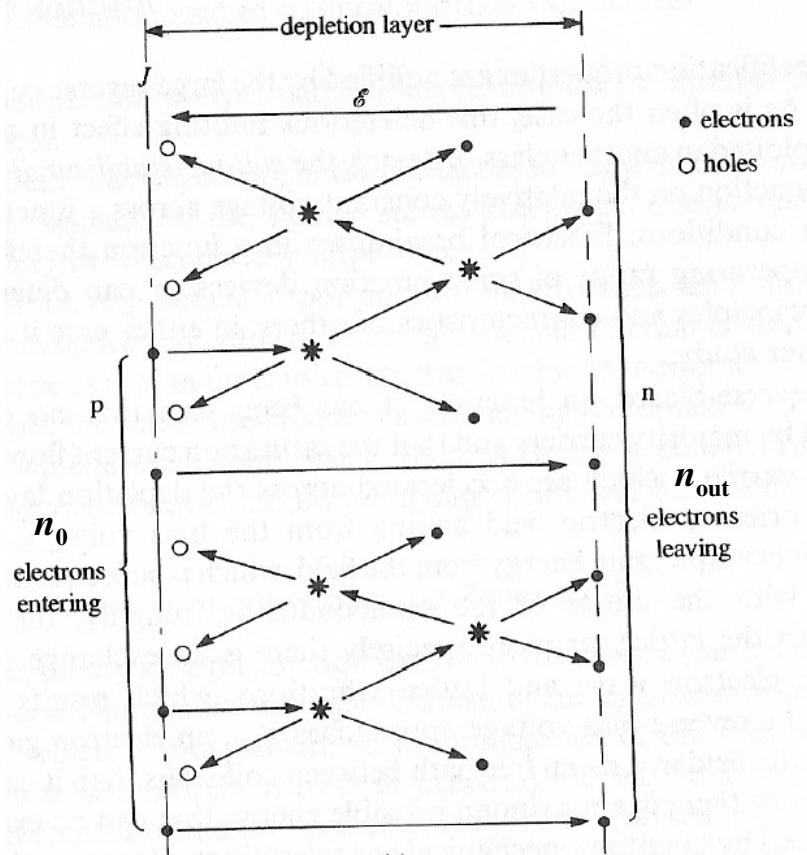
where $P_i = 0$ when $V = 0$
 $P_i \rightarrow 1$ as $V \rightarrow V_{BD}$

- Experimentally it is found:

$$P_i = (V/V_{BD})^n \quad \longrightarrow \quad M = [1 - (V/V_{BD})^n]^{-1}$$

The Zener Diode

- Avalanche breakdown



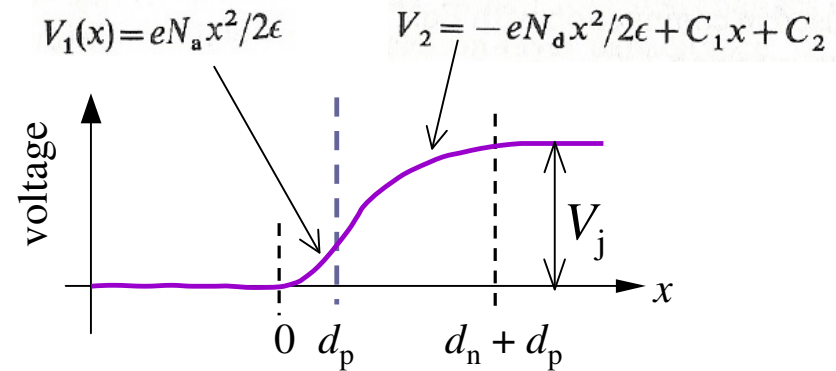
$$\Rightarrow \mathcal{E}_{\text{max,B}} = \left(\frac{2eV_{\text{BD}} N_d N_a}{\epsilon(N_a + N_d)} \right)^{1/2}$$



$$V_{\text{BD}} \propto \frac{N_a + N_d}{N_d N_a}$$

It can be tailored!!!

- Through the depletion layer we found:



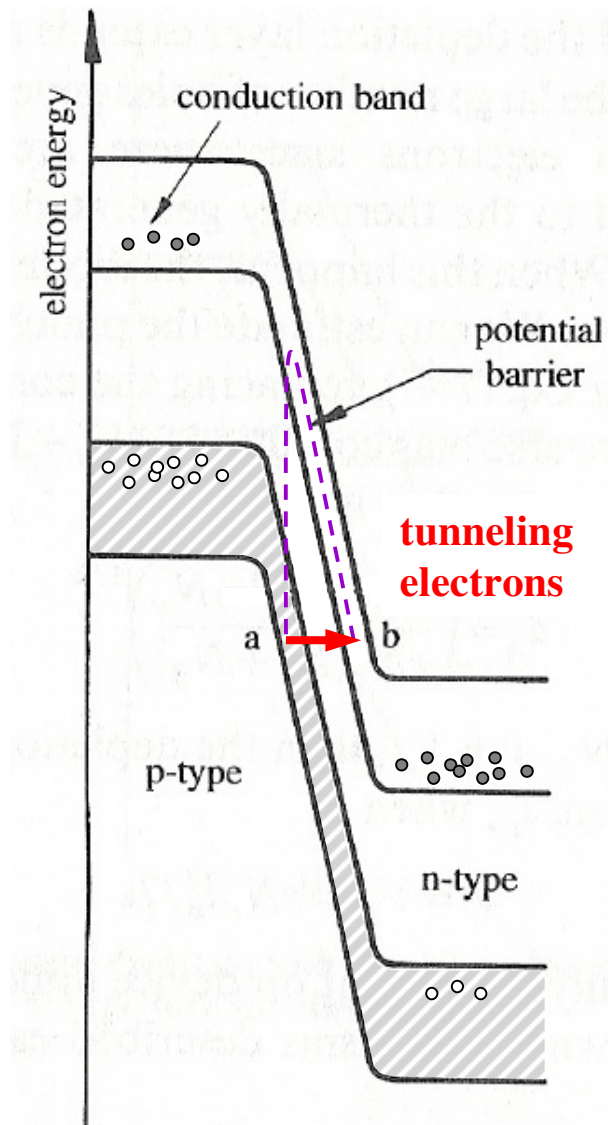
- Electric field in the depletion layer:

$$\frac{\partial V_2}{\partial x} = \mathcal{E} = \frac{-eN_d x}{\epsilon} + \frac{ed_p(N_a + N_d)}{\epsilon}$$

$$\Rightarrow \mathcal{E}_{\max} = ed_p N_a / \epsilon = \left(\frac{2eV_j N_d N_a}{\epsilon(N_a + N_d)} \right)^{1/2}$$

The Zener Diode

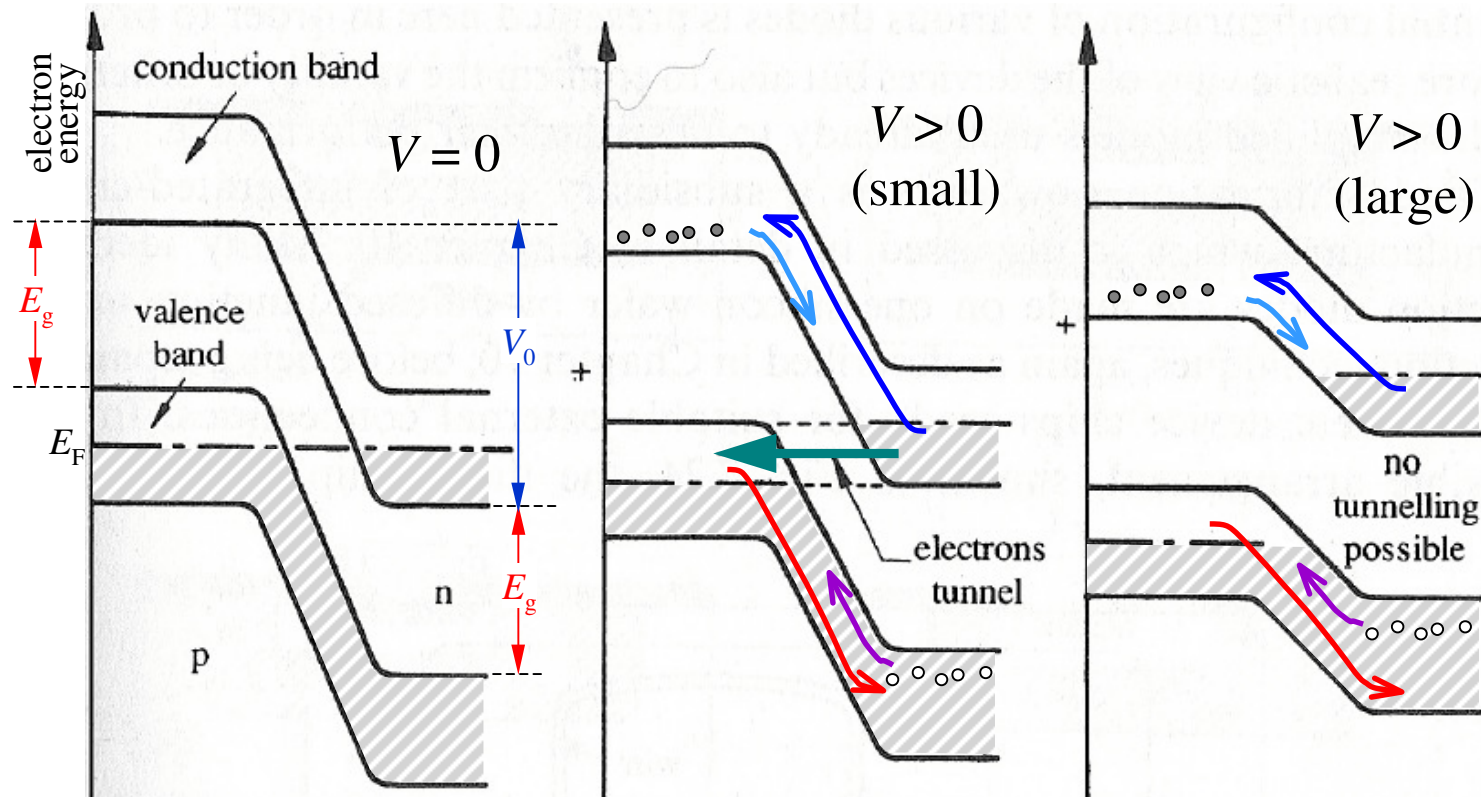
- Zener breakdown



- It occurs in highly doped junctions (thin depletion layer).
- Electrons in the valence band can tunnel to the conduction band.
- V_{BD} is T independent.

The Tunnel Diode

- Consider highly-doped pn junctions
 - Fermi energies cross the bands.
 - V_0 is large [$V_0 \simeq \frac{kT}{e} \log_e \left(\frac{N_d N_a}{n_i^2} \right)$] & d_p, d_n short [$d_p = \left(\frac{2\epsilon V_0 N_d}{e N_a (N_a + N_d)} \right)^{1/2}$]
 - Tunneling is favored.



The Tunnel Diode

- Consider highly-doped pn junctions
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 - Tunneling is favored (very fast response).
 - Extra current source at V small.

