

Physics of Electronics:

7. Junction Diodes

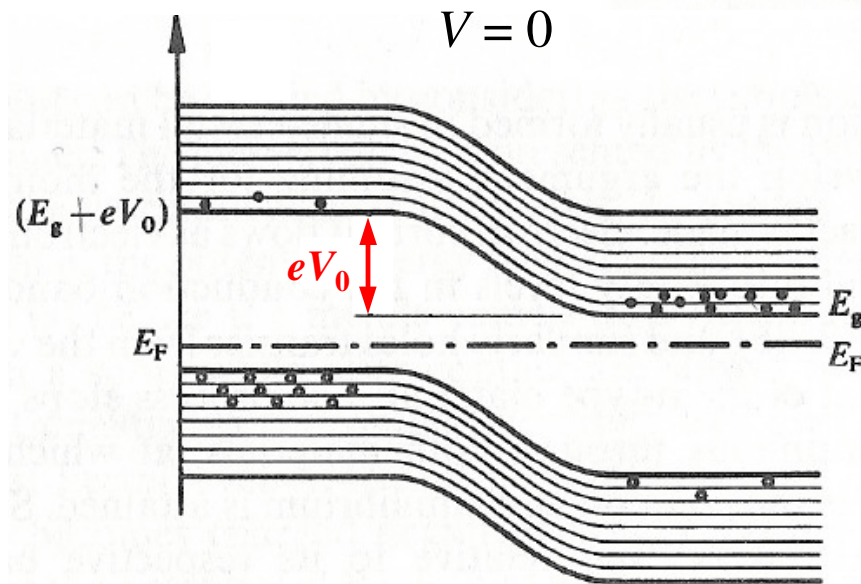
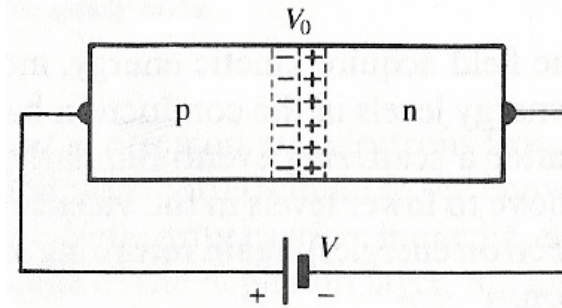
July – December 2008

Contents overview

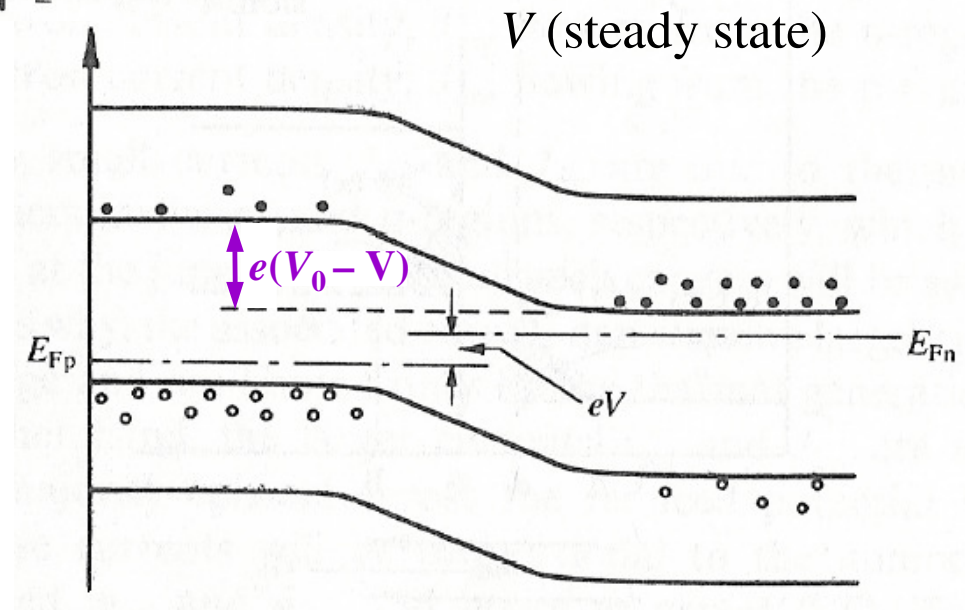
- Current flow in a pn junction with forward bias.
- Current flow in a pn junction with reversed bias.
- IV characteristics of a junction diode
- Electron and hole efficiencies.
- Pn junction with finite dimensions.
- Depletion-layer capacitance
- Small-signal equivalent circuit.

Pn junction with forward bias

- Consider a junction biased with a voltage V :



Same Fermi energy at both sides



Fermi energies separated by eV

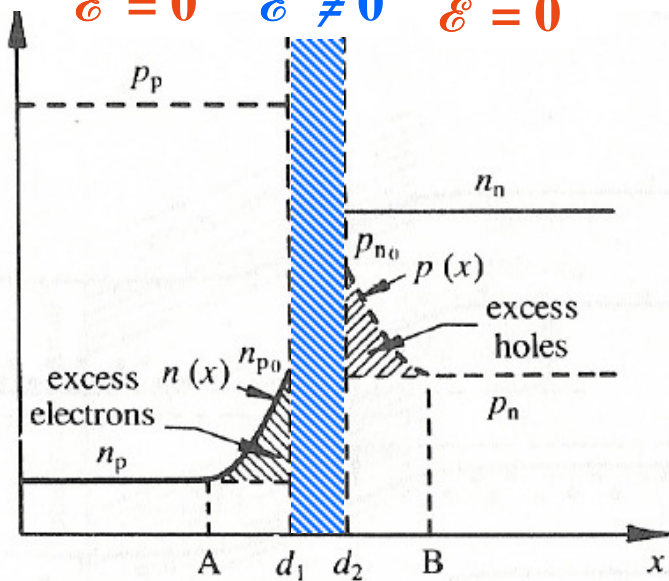
Pn junction with forward bias

- Diode equation:

- Continuity eq. for h^+ at $x \geq d_2$

V (steady state)

$\mathcal{E} = 0$ $\mathcal{E} \neq 0$ $\mathcal{E} = 0$



$$\frac{\partial(\delta p)}{\partial t} = -\frac{\delta p}{\tau_{Lh}} + \mu_h \mathcal{E}_x \frac{\partial(\delta p)}{\partial x} + D_h \frac{\partial^2(\delta p)}{\partial x^2}$$

$L_h^2 = \tau_{Lh} D_h$

$$\Rightarrow \delta p = C_1 \exp(-x/L_h) + C_2 \exp(+x/L_h)$$

$$\Rightarrow p(x) = (p_{n0} - p_n) \exp(-x/L_h) + p_n$$

• Since $\mathcal{E} = 0$, there is only diffusion:

$$J_{Dh} = -eD_h dp/dx$$

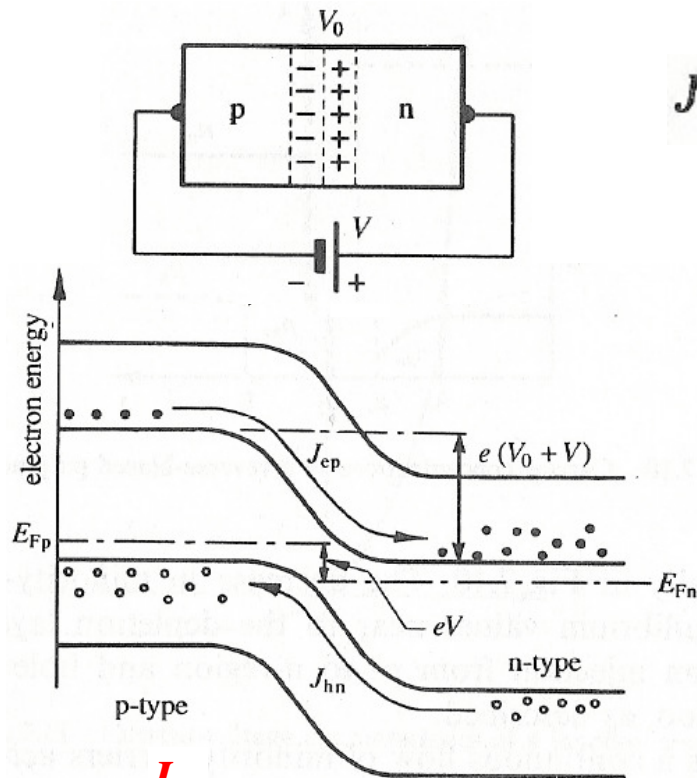
$$\Rightarrow J_h|_{d_2} = (eD_h/L_h)p_n [\exp(eV/kT) - 1]$$

- Continuity eq. for e^- at $x \leq d_1$: $J_e|_{d_1} = (eD_e n_p/L_e) [\exp(eV/kT) - 1]$

$$J = J_h + J_e = e \left(\frac{D_h p_n}{L_h} + \frac{D_e n_p}{L_e} \right) [\exp(eV/kT) - 1]$$

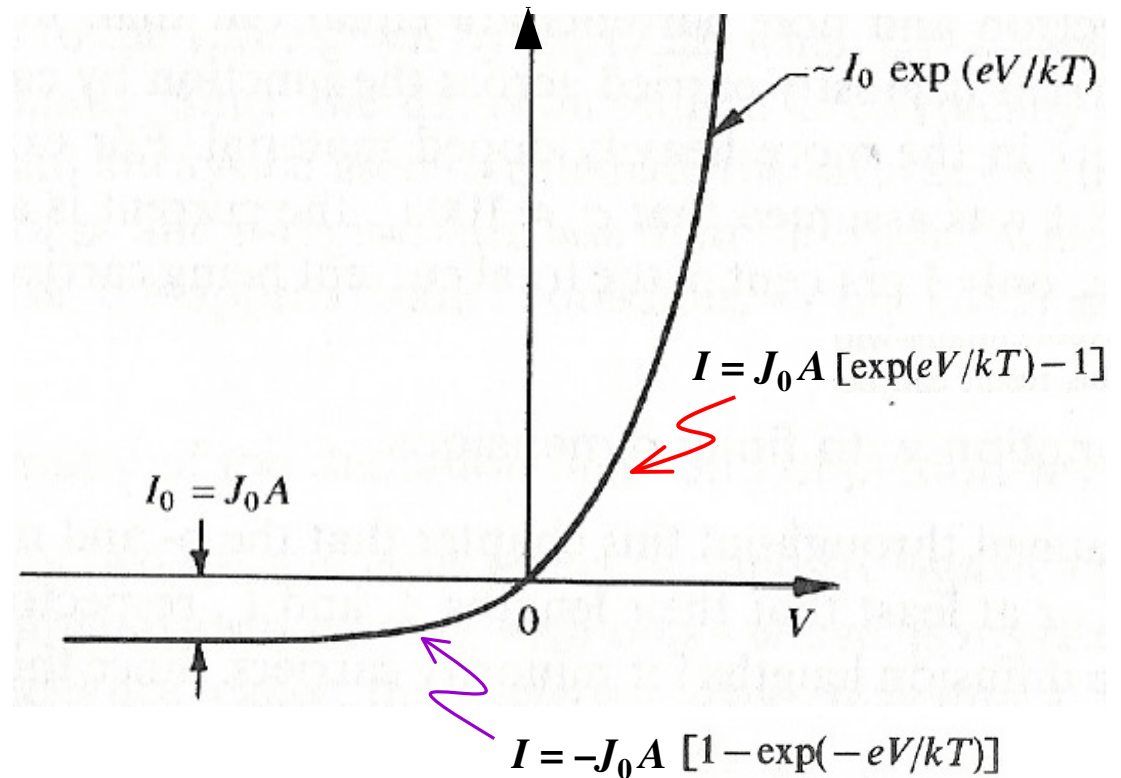
Pn junction with reversed bias

- Currents through the junction :
 - Applying diode equation (with $-V$):



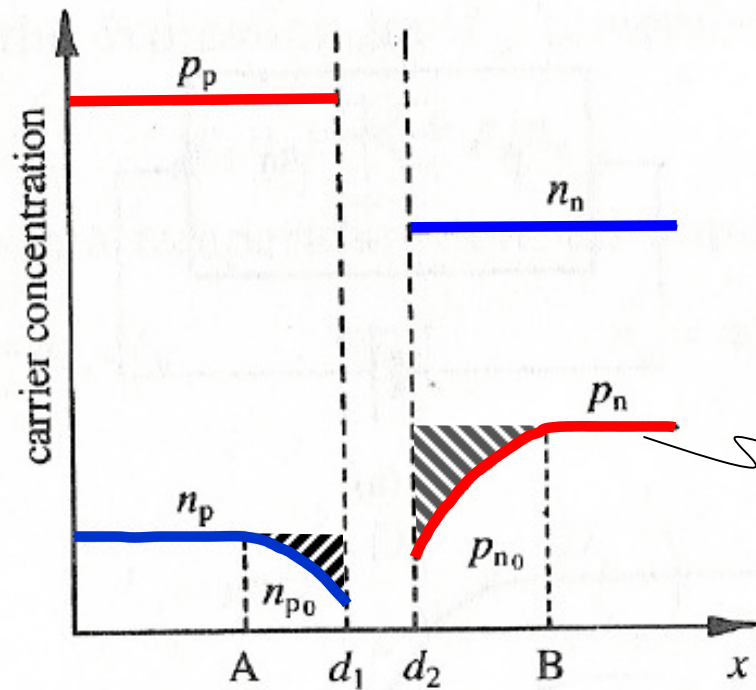
$$J = \cancel{J_{hp}}^0 - J_{hn} + \cancel{J_{en}}^0 - J_{ep} \approx -J_0$$

$$J = J_0 [\exp(-eV/kT) - 1] = -J_0 [1 - \exp(-eV/kT)]$$



Pn junction with reversed bias

- Decrease of minority carriers:
 - As before (with $-V$):



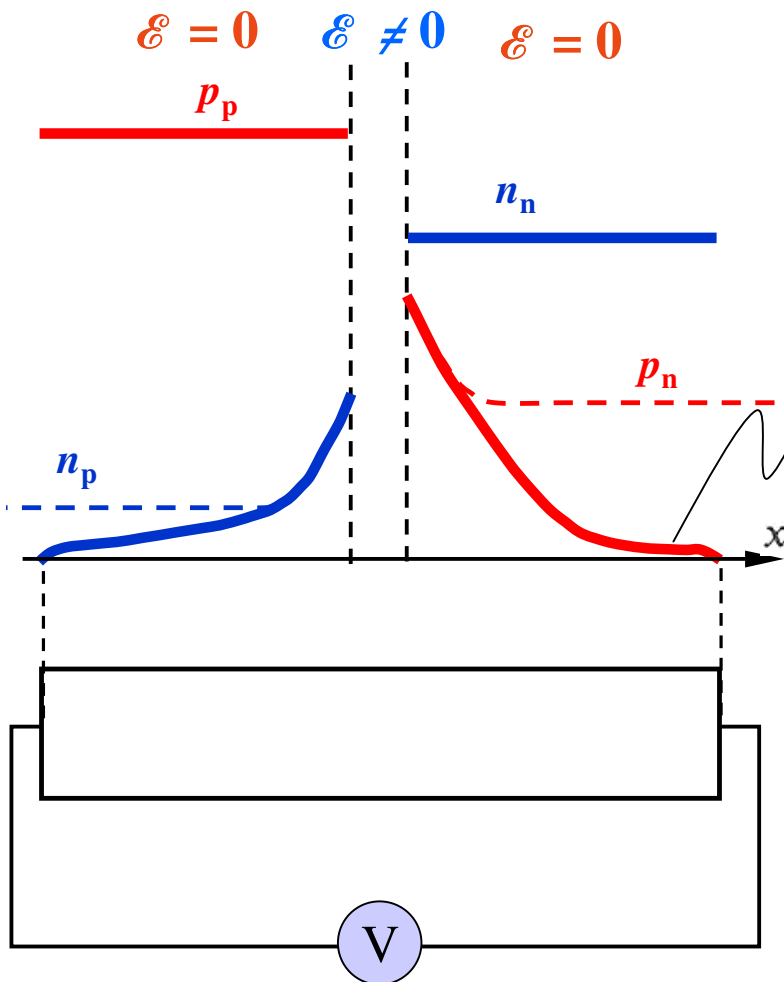
$$n_{p0} = n_p \exp(-eV/kT)$$

$$p_{n0} = p_n \exp(-eV/kT)$$

$$p(x) = p_n \{ [\exp(-eV/kT) - 1] \exp(-x/L_h) + 1 \}$$

Pn junction with finite dimensions

- Current in a finite junction:
 - From the continuity eq. we obtained



$$\delta p = C_1 \exp(-x/L_h) + C_2 \exp(+x/L_h)$$

$$\begin{aligned} p &= p_{n0} = p_n \exp(eV/kT) & \text{when } x=0 \\ p &= 0 & \text{when } x=l_n \end{aligned}$$

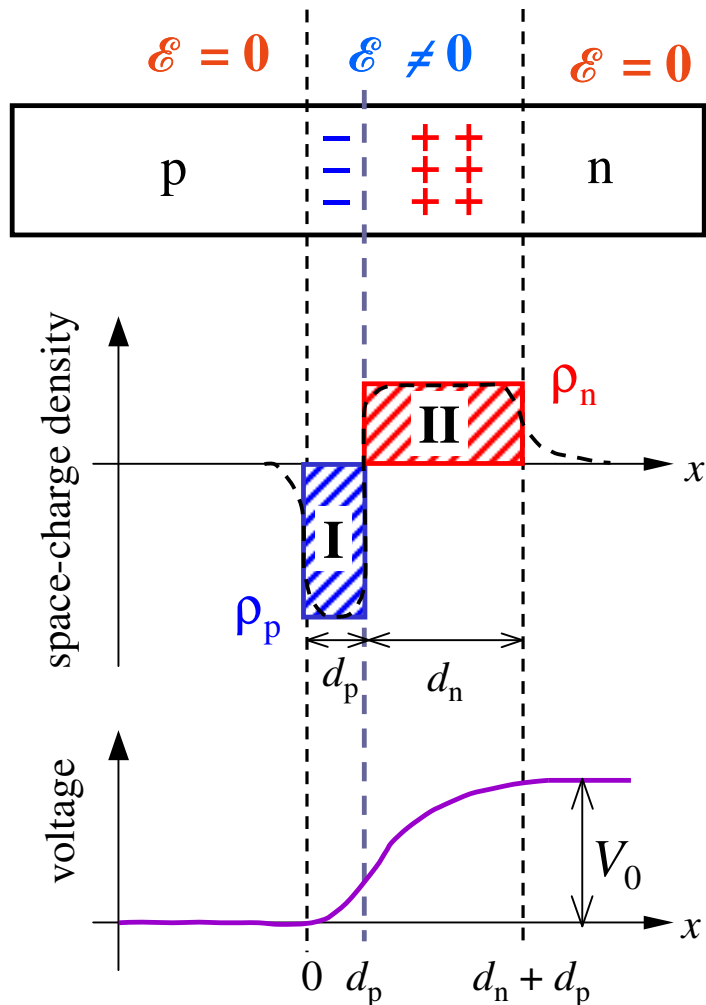
$$\begin{aligned} p(x) &= \left(\frac{p_n \exp(-l_n/L_h) + p_{n0} - p_n}{1 - \exp(-2l_n/L_h)} \right) \exp(-x/L_h) \\ &\quad - \left(\frac{p_n \exp(-l_n/L_h) + p_{n0} - p_n}{1 - \exp(-2l_n/L_h)} - (p_{n0} - p_n) \right) \exp(+x/L_h) + p_n \\ J_h &= -eD_h \left(\frac{dp(x)}{dx} \right) \Big|_{x=0} = \frac{eD_h}{L_h} p_n \tanh\left(\frac{l_n}{L_h}\right) [\exp(eV/kT) - 1] \end{aligned}$$

Idem for the electrons. Therefore the saturation current is:

$$J_0 = e \left[\frac{D_h p_n}{L_h} \tanh\left(\frac{l_n}{L_h}\right) + \frac{D_e n_p}{L_e} \tanh\left(\frac{l_p}{L_e}\right) \right]$$

Depletion-layer capacitance

- The depletion layer forms a capacitance (important for high frequency applications)



Using Poisson relation ($\partial^2 V / \partial x^2 = -\rho / \epsilon$) and continuity for voltage and its derivative:

$$\text{I: } V_1(x) = eN_a x^2 / 2\epsilon$$

$$\text{II: } V_2 = -eN_d x^2 / 2\epsilon + C_1 x + C_2$$

$$C_1 = ed_p(N_a + N_d) / \epsilon \quad C_2 = -\frac{ed_p^2(N_a + N_d)}{2\epsilon}$$

Applying $\mathcal{E} = 0$ at $x \geq d_n + d_p$: $d_p / d_n = N_d / N_a$

Applying $V = V_0$ at $x \geq d_n + d_p$ & using previous relation:

$$d_p = \left(\frac{2\epsilon V_0 N_d}{eN_a(N_a + N_d)} \right)^{1/2}$$

$$d_n = \left(\frac{2\epsilon V_0 N_a}{eN_d(N_a + N_d)} \right)^{1/2}$$

Depletion-layer capacitance

- Capacitance of the depletion layer:
 - Charge per unit area accumulated at the depletion layer

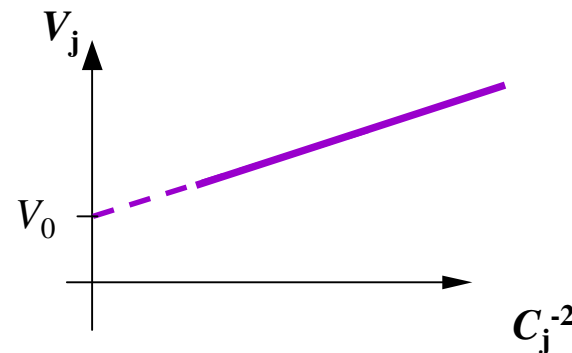
$$Q_j = eN_d d_n = eN_a d_p \quad \xrightarrow{V_j = V_0 \pm V} \quad Q_j = \left(\frac{2\epsilon e V_j N_a N_d}{N_a + N_d} \right)^{1/2}$$

- The capacitance per unit area ($C_j = dQ_j/dV_j$) is then:

$$C_j = \left(\frac{\epsilon e N_a N_d}{2(N_a + N_d)} \right)^{1/2} \frac{1}{V_j^{1/2}} \quad \longrightarrow \quad \boxed{C_j = \epsilon / (d_p + d_n)}$$

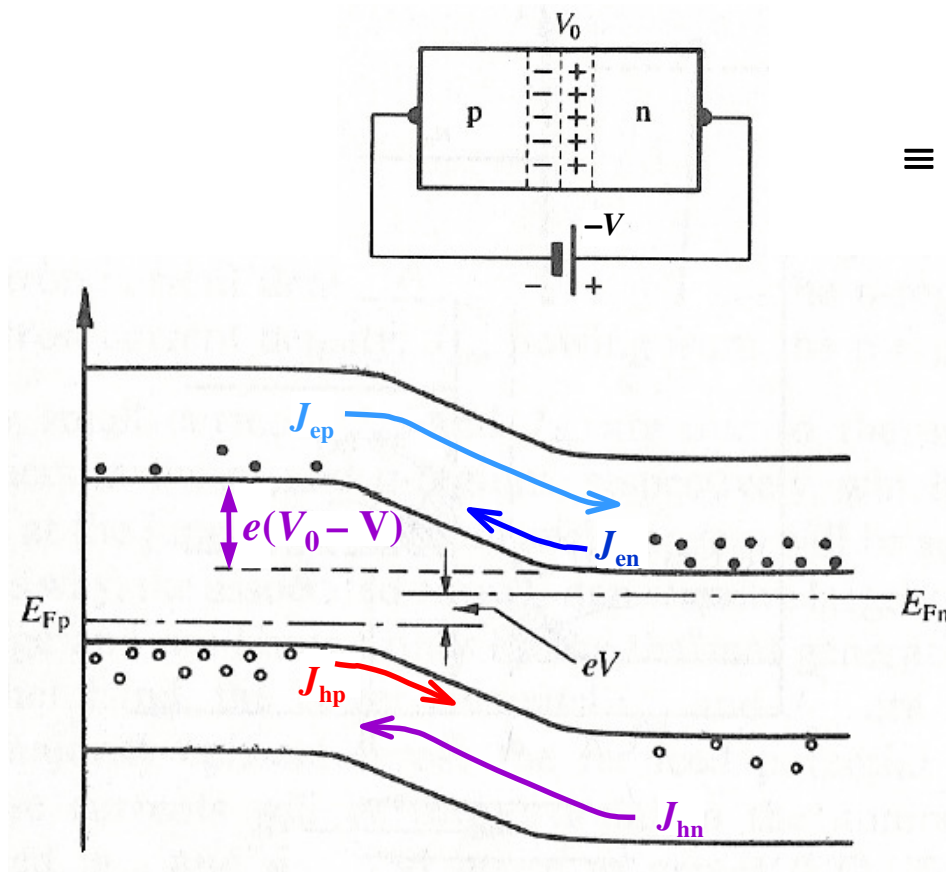
$$\downarrow N_a \gg N_d$$

$$V_j = V_0 + V \simeq \frac{1}{C_j^2} \left(\frac{\epsilon e N_d}{2} \right)^{1/2}$$

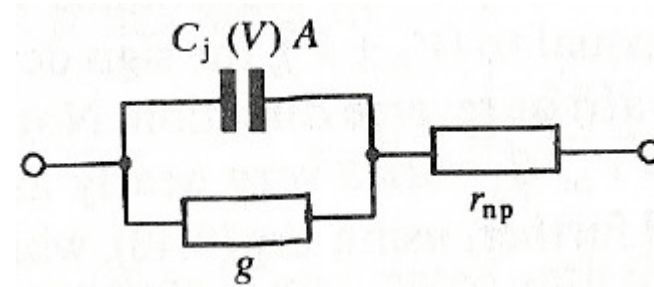


Small-AC-signal equivalent circuit

- Reversed biased junction



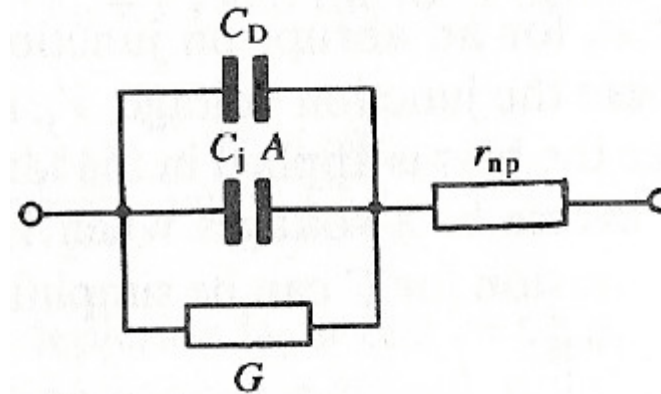
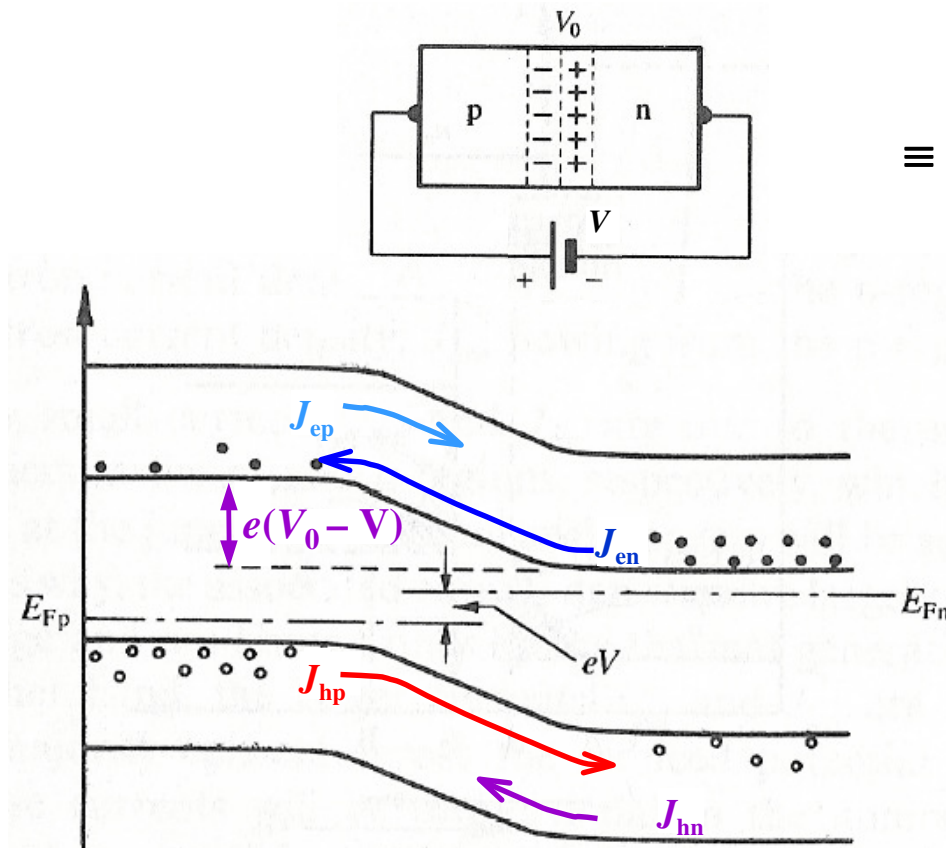
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r_{np} : Resistance of the junction @ $V=0$
 $C_j(V)$: Depletion layer capacitance
 g : Conductance of the reverse flow

Small-AC-signal equivalent circuit

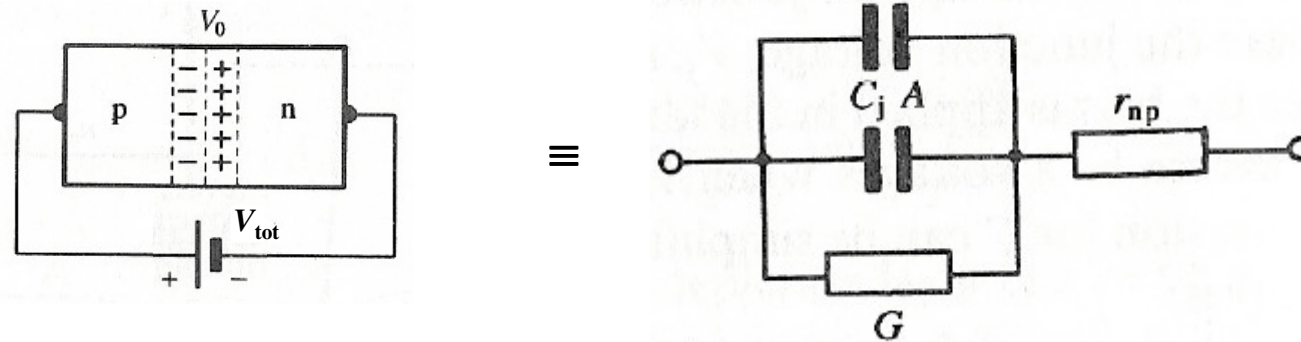
- Forward biased junction



- r_{np} : Resistance of the junction @ $V=0$
- $C_j(V)$: Depletion layer capacitance
- G : Conductance of the forward flow
- C_D : Diffusion capacitance

Small-AC-signal equivalent circuit

- Forward biased junction



$$\left. \begin{aligned} V_{tot} &= V + V_1 \exp(i\omega t) \\ p_{n0} &= p_n \exp(eV/kT) \end{aligned} \right\} \Rightarrow p_{n0} = p_n \exp\left(\frac{e}{kT}[V + V_1 \exp(j\omega t)]\right)$$

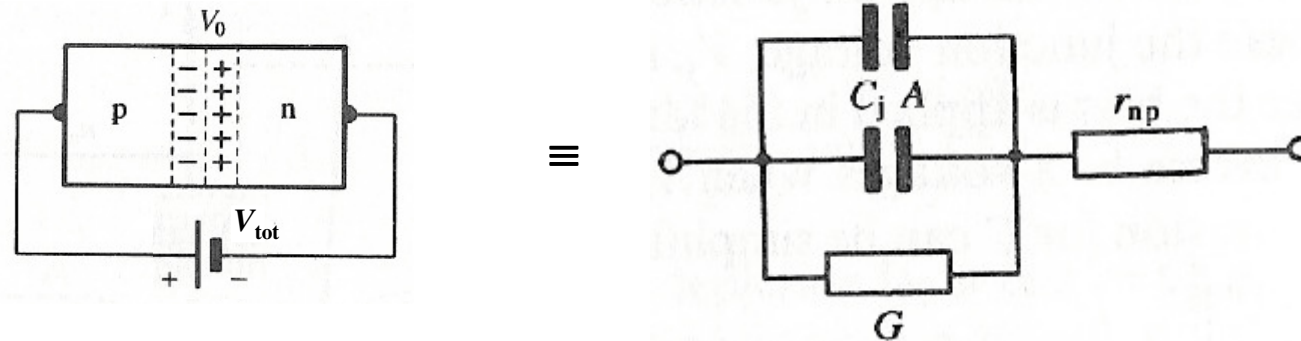
$$V_1 \ll V_0 \Rightarrow p_{n0} = \underbrace{p_n \exp\left(\frac{eV}{kT}\right)}_{\text{DC}} \underbrace{\left(1 + \frac{eV_1}{kT} \exp(j\omega t)\right)}_{\text{AC}}$$

- Therefore we can assume:

$$\delta p = p(x, t) - p_n = \underbrace{p_0(x)}_{\text{DC}} + \underbrace{p_1(x) \exp(j\omega t)}_{\text{AC}} - p_n$$

Small-AC-signal equivalent circuit

- Forward biased junction



- Applying continuity equation:

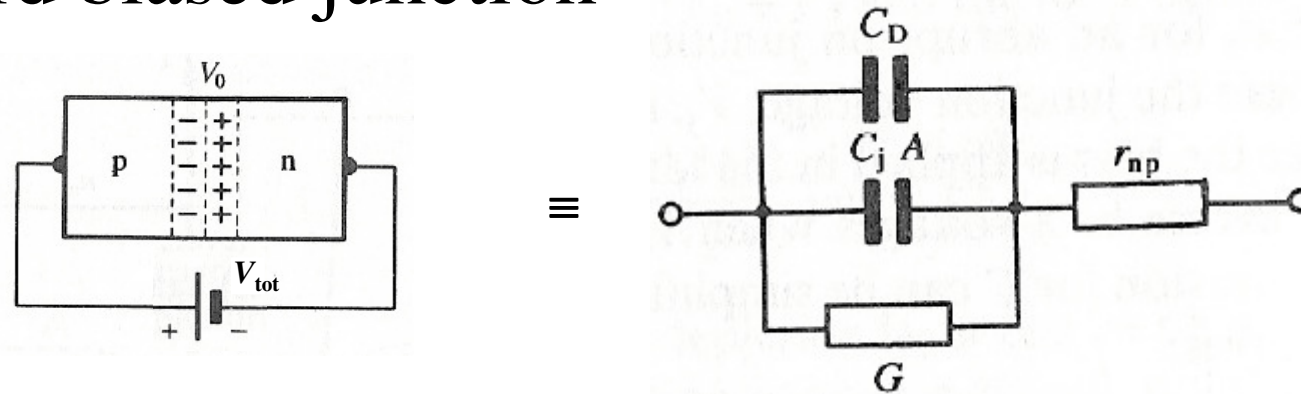
$$\frac{\partial(\delta p)}{\partial t} = -\frac{\delta p}{\tau_{Lh}} + \mu_h \cancel{E_x}^0 \frac{\partial(\delta p)}{\partial x} + D_h \frac{\partial^2(\delta p)}{\partial x^2} \quad \Rightarrow \quad j\omega p_1 = \frac{-p_1}{\tau_{Lh}} + D_h \frac{\partial^2 p_1}{\partial x^2}$$

$$\Rightarrow \frac{\partial^2 p_1}{\partial x^2} = \frac{1 + j\omega\tau_{Lh}}{L_h^2} p_1 \quad \Rightarrow \quad p_1 = C \exp\left(-\frac{(1 + j\omega\tau_{Lh})^{1/2} x}{L_h}\right)$$

- But $p_1(x=0) = p_{n0} \Rightarrow p_1 = p_n \left(\frac{eV_1}{kT}\right) \exp\left(\frac{eV}{kT}\right) \exp\left(-\frac{(1 + j\omega\tau_{Lh})^{1/2} x}{L_h}\right) \exp(j\omega t)$

Small-AC-signal equivalent circuit

- Forward biased junction



- Alternating currents of holes injected into n-region:

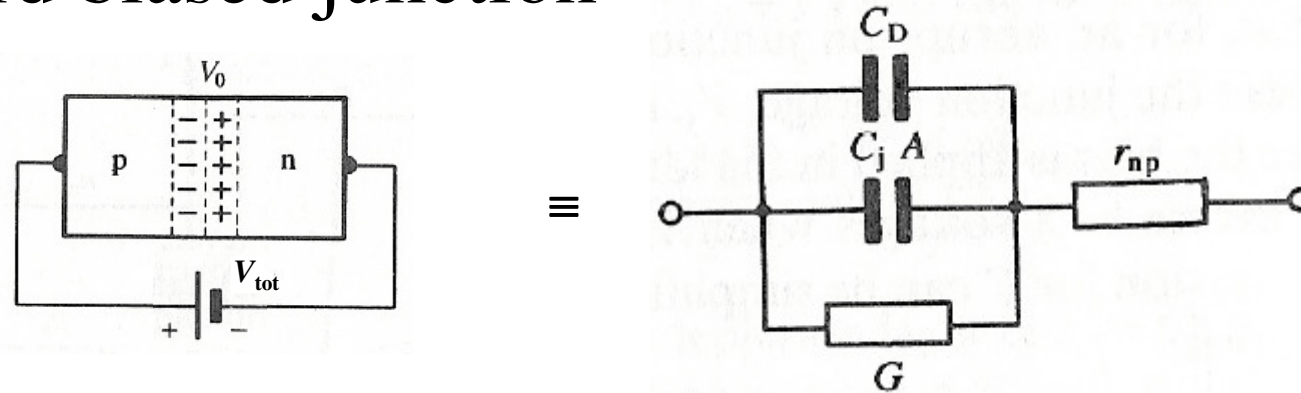
$$J_{h1} = -eD_h \left(\frac{dp(x)}{dx} \right) \Big|_{x=0} = \frac{(1 + i \omega \tau_{Lh})^{1/2}}{L_h} \frac{p_n D_h e^2 V_1}{kT} \exp\left(\frac{eV}{kT}\right) \exp(j\omega t)$$

- Alternating currents of electrons injected into p-region:

$$J_{e1} = -eD_e \left(\frac{dn(x)}{dx} \right) \Big|_{x=0} = \frac{(1 + i \omega \tau_{Le})^{1/2}}{L_e} \frac{n_p D_e e^2 V_1}{kT} \exp\left(\frac{eV}{kT}\right) \exp(j\omega t)$$

Small-AC-signal equivalent circuit

- Forward biased junction



- Total alternating currents of minority injected carriers:

$$J_1 = J_{h1} + J_{e1}$$

- Admittance ($Y_1 = J_1/V_1$) for $\omega\tau \ll 1$:

$$Y_1 \cong \underbrace{\frac{e^2}{kT} \exp\left(\frac{eV}{kT}\right) \left(\frac{D_h p_n}{L_h} + \frac{D_e n_p}{L_e}\right)}_G + \underbrace{\frac{j\omega e^2}{2kT} \left(\frac{D_h p_n \tau_{Lh}}{L_h} + \frac{D_e n_p \tau_{Le}}{L_e}\right)}_{C_D}$$