

# Physics of Electronics:

## 6. Semiconductors

July – December 2008

# Contents overview

- Density of carriers in extrinsic semiconductors.
- Compensation doping.
- Electrical conduction in semiconductors.
- Continuity equation.
- Semiconductor measurements.

# Density of carriers in extrinsic sc

- General considerations:
  - Concentration of ionized impurities:

ionized impurities = impurity concentration  $\times$  prob. finding e/h at the imp. level

$$N_a^- = \frac{N_a}{1 + \exp[(E_A - E_F)/kT]} \quad \& \quad N_d^+ = N_d \left( 1 - \frac{1}{1 + \exp[(E_D - E_F)/kT]} \right)$$

- Neutrality condition

holes + ionized donors = electrons + ionized acceptors

$$p + N_d^+ = n + N_a^-$$

$$\rightarrow N_v \exp(-E_F/kT) + \frac{N_d}{1 + \exp[-(E_D - E_F)/kT]} = N_c \exp[-(E_g - E_F)/kT] + \frac{N_a}{1 + \exp[(E_A - E_F)/kT]}$$

which can be solved for  $E_F$

# Density of carriers in extrinsic sC

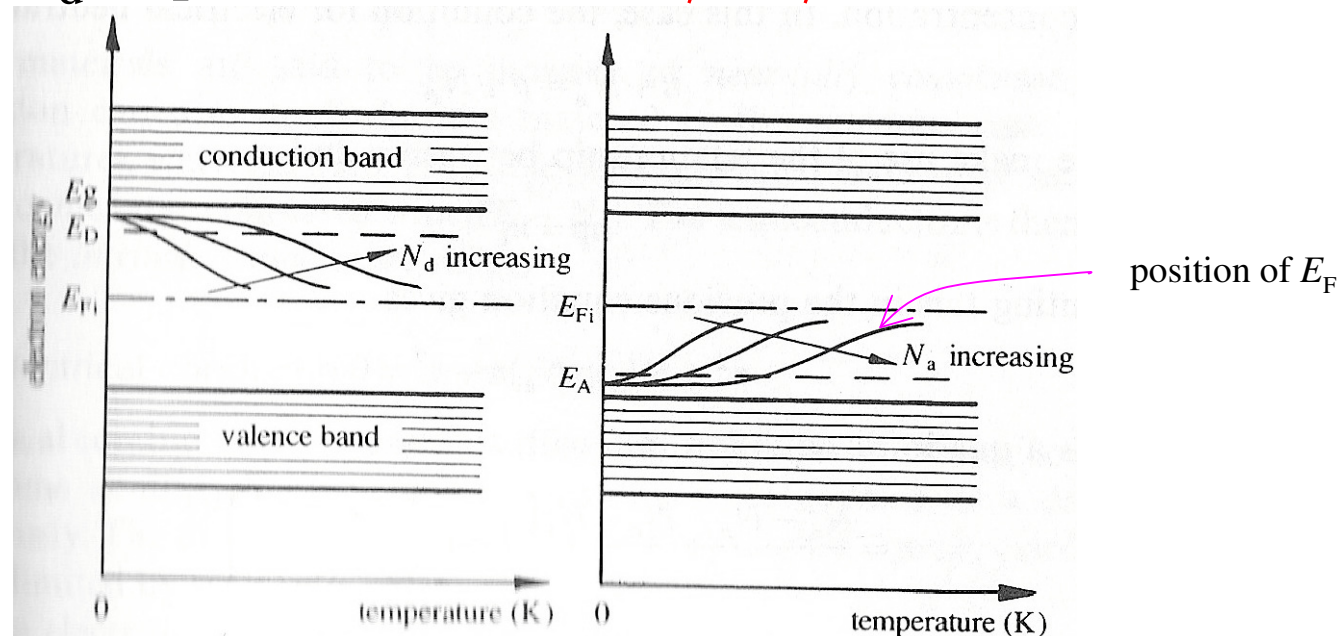
- n-type material:

- $N_d \gg N_a$  &  $n \gg p \rightarrow p + N_d^+ = n + N_a^-$

- $$\rightarrow \exp(E_F/kT) = \frac{-1 + \{1 + (4N_d/N_c) \exp[(E_g - E_D)/kT]\}^{1/2}}{2 \exp(-E_D/kT)}$$

- p-type material:

- $N_a \gg N_d$  &  $p \gg n \rightarrow p + N_d^+ = n + N_a^-$



# Compensation doping

- Compensation: both donors and acceptors
  - Suppose all impurities are ionized ( $N_a = N_a^+$  &  $N_d = N_d^+$ )
  - Neutrality:  $n + N_a = p + N_d$   $\xrightarrow{np = n_i^2}$   $n^2 - (N_d - N_a)n - n_i^2 = 0$

$$\rightarrow n = \frac{N_d - N_a}{2} + \frac{N_d - N_a}{2} \left[ 1 + \left( \frac{2n_i}{N_d - N_a} \right)^2 \right]^{1/2} \quad (\text{Idem for } p)$$

- Extrinsic compensated  $n_i \ll |N_d - N_a|$

$$\text{for } N_d > N_a: \quad \begin{aligned} n &\simeq N_d - N_a \\ p &\simeq n_i^2 / (N_d - N_a) \end{aligned}$$

$$\text{for } N_a > N_d: \quad \begin{aligned} p &\simeq N_a - N_d \\ n &\simeq n_i^2 / (N_a - N_d) \end{aligned}$$

- Intrinsic compensated  $n_i \gg |N_d - N_a|$

$$n = n_i + (N_d - N_a)/2$$

$$p = n_i + (N_d - N_a)/2$$

# Electrical conduction in sC.

- Conductivity

- In a sC we have  $e^-$ 's and  $h^+$ 's moving:  $\sigma$

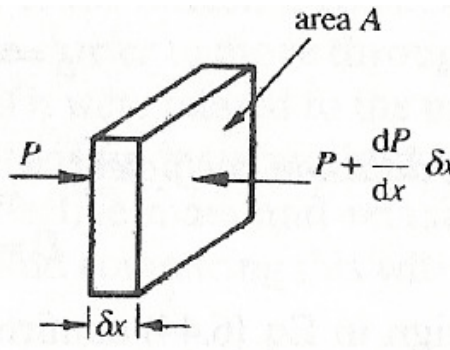
$$J_e = ne\mu_e \mathcal{E} \quad \& \quad J_h = pe\mu_h \mathcal{E} \quad \longrightarrow \quad J = J_e + J_h = \overbrace{e(n\mu_e + p\mu_h)}^{\sigma} \mathcal{E}$$

- In a n-type sC ( $n \gg p$ ):  $\sigma_n \simeq en\mu_e \simeq eN_d\mu_e$

- In a p-type sC ( $p \gg n$ ):  $\sigma_p \simeq ep\mu_h \simeq eN_a\mu_h$

- Diffusion (when density gradients are present):

- In a neutral gas with gradient in the  $x$  direction:



Force on the elemental volume

$$\left[ P(x) - \left( P + \frac{dP}{dx} \delta x \right) \right] A = -\frac{dP}{dx} \delta x A$$

Force on each particle

$$F_D = -\frac{1}{N} \frac{dP}{dx}$$

$$v_D = -\frac{\tau_r}{M} \frac{1}{N} \frac{dP}{dx} \quad \xrightarrow{P = NkT} \quad v_D = -\frac{\tau_r kT}{M} \frac{1}{N} \frac{dN}{dx}$$

(Note: In the final equation, the term  $\tau_r kT / M$  is circled with a dashed red line and labeled with a red 'D' above it.)

# Electrical conduction in sC.

- Diffusion of charge carriers

- In a sC with a concentration gradient of the  $e^-$  and  $h^+$ :

$$v_{De} = -D_e \frac{1}{n} \frac{dn}{dx} \quad \Rightarrow \quad J_{De} = -nev_D = eD_e \frac{dn}{dx}$$

$$v_{Dh} = -D_h \frac{1}{p} \frac{dp}{dx} \quad \Rightarrow \quad J_{Dh} = -eD_h \frac{dp}{dx}$$

- Diffusion coefficient is related with mobility:

$$D = \frac{\tau_r kT}{M} \xrightarrow{\mu = e\tau_r/m} \begin{matrix} D_e = (kT/e)\mu_e \\ D_h = (kT/e)\mu_h \end{matrix} \Rightarrow \boxed{D_e/\mu_e = D_h/\mu_h = kT/e}$$

- Total current flow

$$\boxed{J_e = ne\mu_e \mathcal{E} + eD_e \nabla n}$$

$$\boxed{J_h = pe\mu_h \mathcal{E} - eD_h \nabla p}$$

$$\xrightarrow{J_h = 0}$$

$$\mathcal{E} = D_h \nabla p / p \mu_h$$

current can be compensated at a particular electrical field

# Continuity Equation for Minority Carriers

- General considerations:
  - Describes how charge density varies in time.
  - Two main causes for the variation:
    - Rate of generation and loss
    - Drift of carriers
- No current flow (no gradient & no  $\mathcal{E}$ )
  - Let's consider a **p-type** material:

$$\frac{dn}{dt} = G - R$$

Variation = Generation – Recombination

In this case, only  $T$   
dependent

$$R = r n(t) p(t)$$



# Continuity Equation for Minority Carriers

- No current flow (no gradient & no  $\mathcal{E}$ )

- Consider first equilibrium:

$$dn/dt=0 \quad \Rightarrow \quad G(T)=rn_0p_0=rn_i^2$$

- Now consider variations from the equilibrium:

$$n = n_0 + \delta n \quad \longleftrightarrow \quad p = p_0 + \delta p \quad \text{with } \delta n = \delta p \quad \leftarrow \text{To maintain charge neutrality}$$

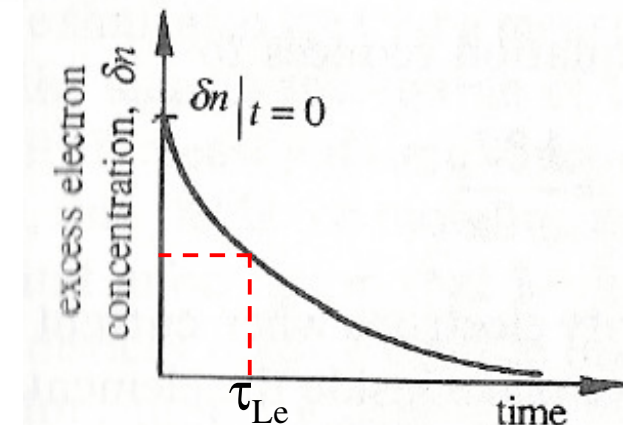
$$\Rightarrow \quad (dn/dt)|_{J=0} = [d(\delta n)/dt]|_{J=0} = G(T) - r(n_0 + \delta n)(p_0 + \delta p)$$

$$\Rightarrow \quad [d(\delta n)/dt]|_{J=0} = -\cancel{rn_0\delta p} - \cancel{rp_0\delta n} - \cancel{r\delta n\delta p} \quad \leftarrow \text{Considering } T=\text{const.}$$

$$\Rightarrow \quad [d(\delta n)/dt]|_{J=0} = -rp_0\delta n = -\delta n/\tau_{Le}$$

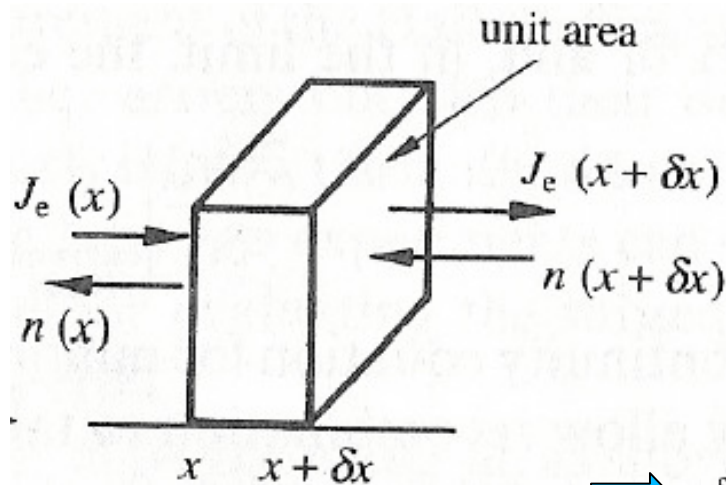
- For a **n-type** material:

$$d(\delta p)/dt = -\delta p/\tau_{Lh}$$



# Continuity Equation for Minority Carriers

- With current flow



$$J_e = -nev_D \rightarrow \begin{cases} n_{x+\delta x} = J_e(x+\delta x)/ev_D \\ n_x = J_e(x)/ev_D \end{cases}$$

$$\begin{aligned} n_{x+\delta x} - n_x &= \delta n = [J_e(x+\delta x) - J_e(x)]/ev_D \\ &= \frac{1}{ev_D} \left( J_e(x) + \frac{\partial J_e}{\partial x} \delta x - J_e(x) \right) \end{aligned}$$

$$\rightarrow \delta n = \frac{1}{ev_D} \frac{\partial J_e}{\partial x} \delta x \xrightarrow{v_D = \partial x / \partial t} \left. \frac{\partial(\delta n)}{\partial t} \right|_{\text{current flow}} = \frac{1}{e} \frac{\partial J_e}{\partial x}$$

- Total continuity equation

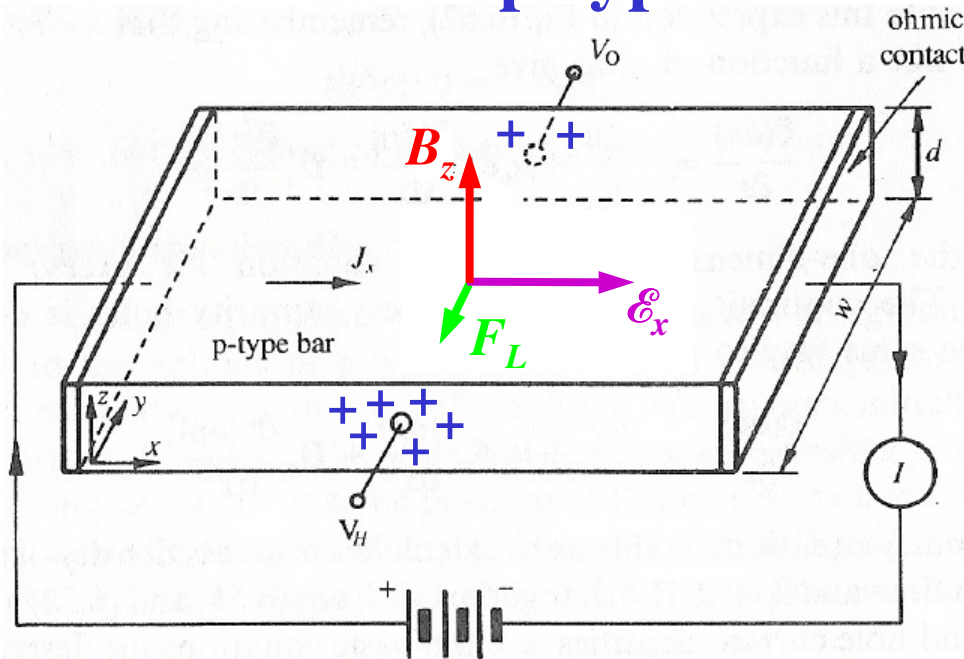
$$\left. \frac{\partial(\delta n)}{\partial t} \right|_{\text{total}} = -\frac{\delta n}{\tau_{Le}} + \frac{1}{e} \frac{\partial J_e}{\partial x} \quad (\text{idem for holes})$$

$$\text{But } J_e = ne\mu_e \mathcal{E} + eD_e \nabla n \xrightarrow{n = n_0 + \delta n}$$

$$\frac{\partial(\delta n)}{\partial t} = -\frac{\delta n}{\tau_{Le}} + \mu_e \mathcal{E}_x \frac{\partial(\delta n)}{\partial x} + D_e \frac{\partial^2(\delta n)}{\partial x^2}$$

# Semiconductor Measurements

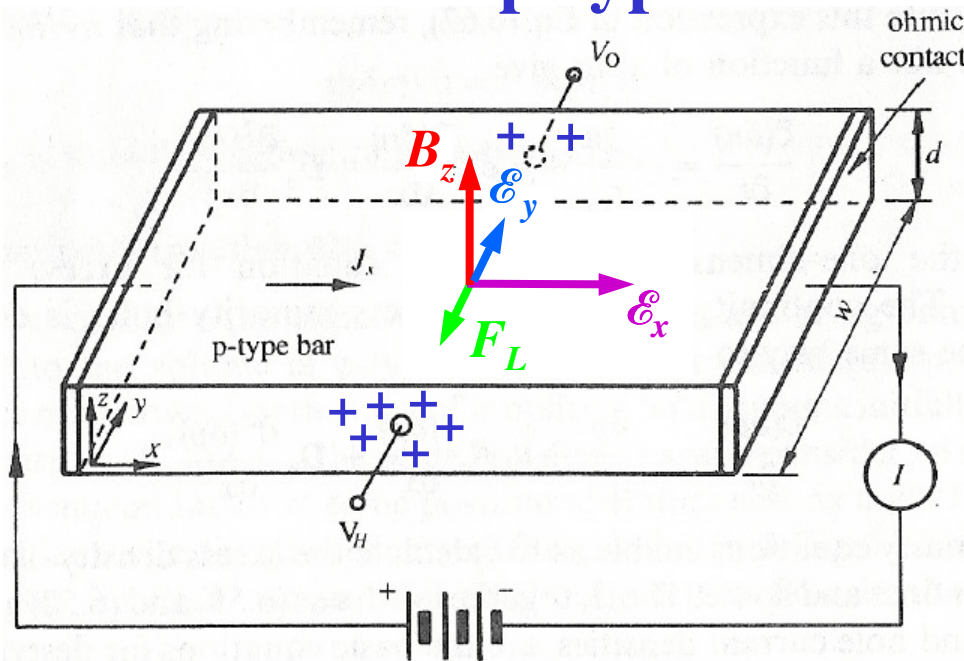
- Hall effect
  - Consider **p-type** material



$$F_L = ev \times B \Rightarrow |F_L| = e v_{Dx} B_z$$

# Semiconductor Measurements

- Hall effect
  - Consider **p-type** material

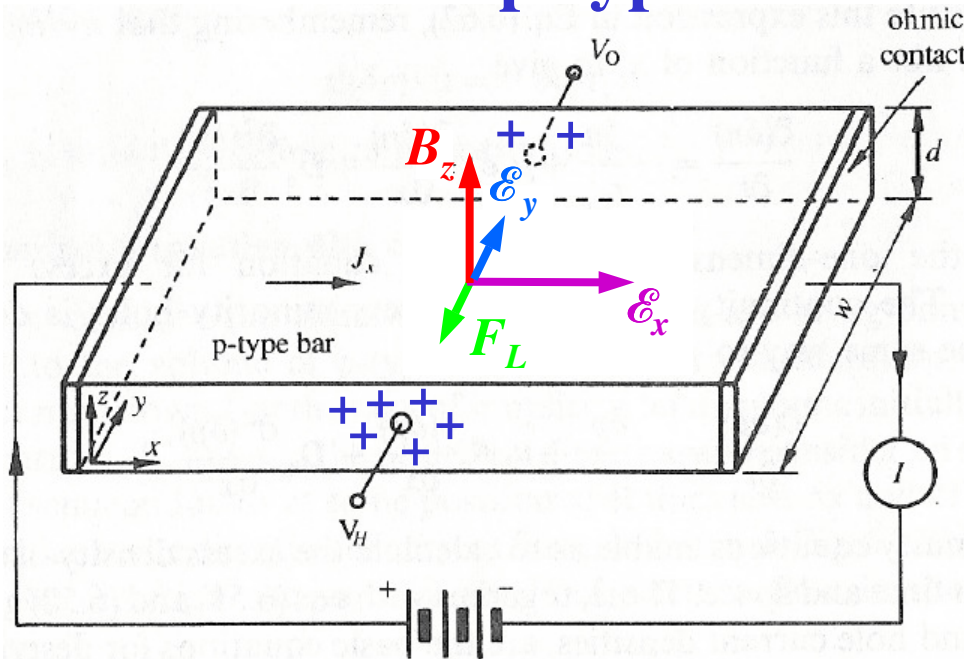


$$F_L = ev \times B \Rightarrow |F_L| = e v_{Dx} B_z$$

$$\Rightarrow e\mathcal{E}_y = F_L = ev_{Dx} B_z$$

# Semiconductor Measurements

- Hall effect
  - Consider **p-type** material



$$F_L = ev \times B \Rightarrow |F_L| = e v_{Dx} B_z$$

$$\begin{aligned} \Rightarrow \left. \begin{aligned} e\mathcal{E}_y &= F_L = ev_{Dx} B_z \\ &\& J_x \simeq pev_{Dx} \end{aligned} \right\} \mathcal{E}_y = J_x B_z / pe \end{aligned}$$

The Hall coefficient is defined:

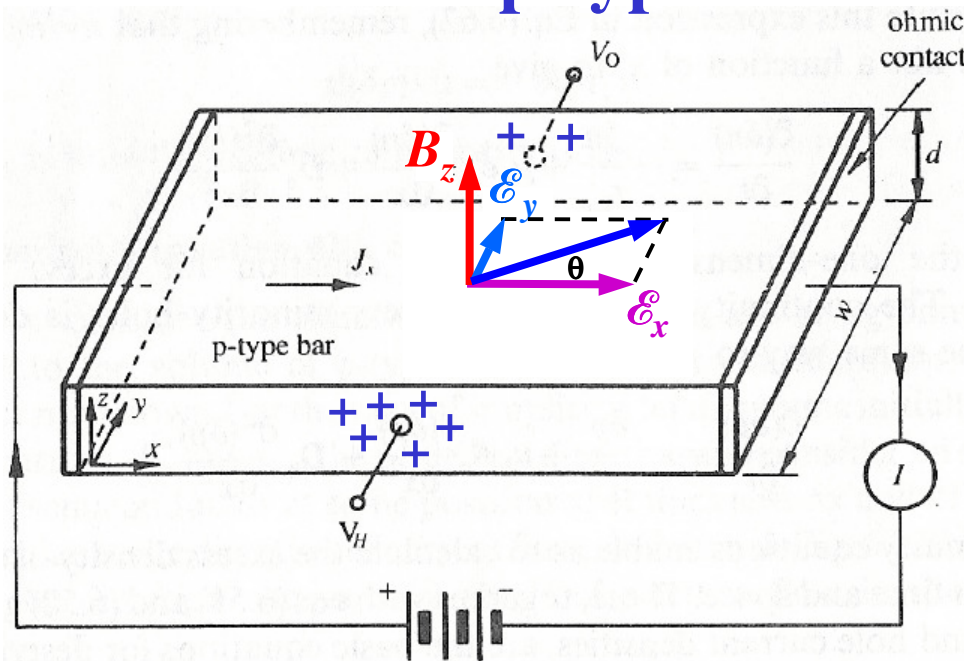
$$R_H = \mathcal{E}_y / J_x B_z = 1/pe$$

- If  $V_H$  and  $I$  are measured:

$$\left. \begin{aligned} V_H &= \mathcal{E}_y w \\ I &= J_x w d \end{aligned} \right\} R_H = \frac{V_H}{w I B_z} w d = \frac{V_H d}{I B_z} = \frac{1}{pe}$$

# Semiconductor Measurements

- Hall effect
  - Consider **p-type** material



The Hall angle is defined:

$$\tan \theta = \mathcal{E}_y / \mathcal{E}_x$$

$$\Rightarrow \tan \theta = \frac{J_x B_z}{pe} \frac{\sigma}{J_x} = \mu_h B_z$$

$$\Rightarrow \mu_h = R_H \sigma$$

- For a **n-type** material:

$$R_{He} = -1/ne$$



# Semiconductor Measurements

- Hall effect
  - Consider both type of carriers to be present:

$$\begin{array}{l} v_{Dh} = \mu_H \mathcal{E}_x \\ v_{De} = -\mu_H \mathcal{E}_x \end{array} \quad \Rightarrow \quad \begin{array}{l} F_h = -e(v_{Dh} \times B) = -ev_{Dh} B_z \\ F_e = e(v_{De} \times B) = -ev_{De} B_z \end{array}$$

- A net current is created in the y direction

$$\sigma \mathcal{E}_y = e(pv_{yh} - nv_{ye})$$

- Now, using the expressions for the Hall angle

$$v_{yh} = \mu_h \mathcal{E}_y = \mu_h (\mathcal{E}_x \tan \theta) = \mu_h (\mathcal{E}_x \mu_h B_z) \quad \& \quad v_{ye} = \mu_e^2 \mathcal{E}_x B_z$$

$$\Rightarrow \sigma \mathcal{E}_y = e(p\mu_h^2 - n\mu_e^2) \mathcal{E}_x B_z$$

- Hall coefficient

$$R_H = \frac{\mathcal{E}_y}{J_x B_z} = \frac{e(p\mu_h^2 - n\mu_e^2)}{\sigma^2}$$

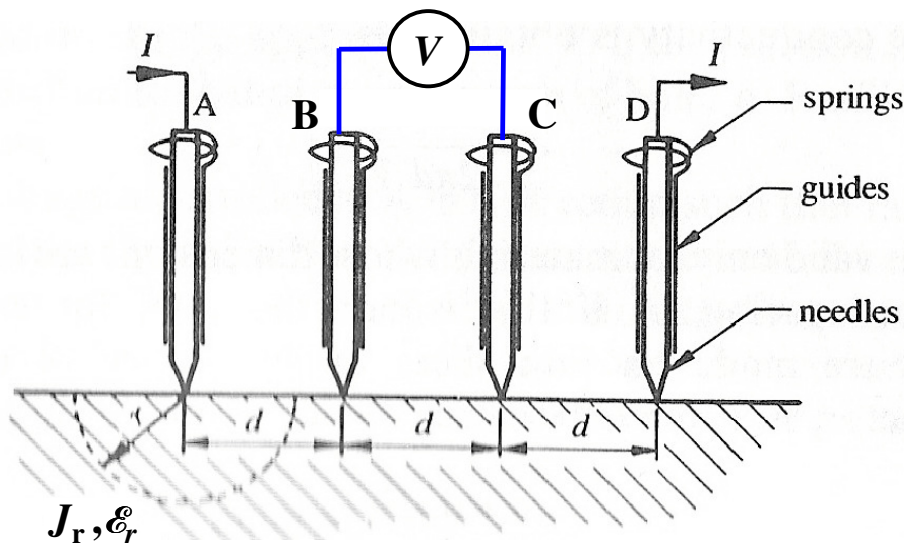
$$\sigma = e(n\mu_e + p\mu_h)$$



$$R_H = \frac{p\mu_h^2 - n\mu_e^2}{e(p\mu_h + n\mu_e)^2}$$

# Semiconductor Measurements

- Four-point probe method for conductivity meas.
  - Consider  $d \ll L$



For the current leaving A:

$$J_r = I / (2\pi r^2) \Rightarrow \mathcal{E}_r = J / \sigma = I / (2\pi \sigma r^2)$$

Potential at a distance  $a$  from A:

$$V_a = \int_{-\infty}^a \mathcal{E}_r \, dr = -\frac{I}{2\pi\sigma} \int_{-\infty}^a \frac{1}{r^2} \, dr = \frac{I}{2\pi\sigma a}$$

$$\Rightarrow V_{BC} = \frac{I}{2\pi\sigma d} - \frac{I}{2\pi\sigma(2d)} = \frac{I}{4\pi\sigma d}$$

Identically for the current entering D, then:

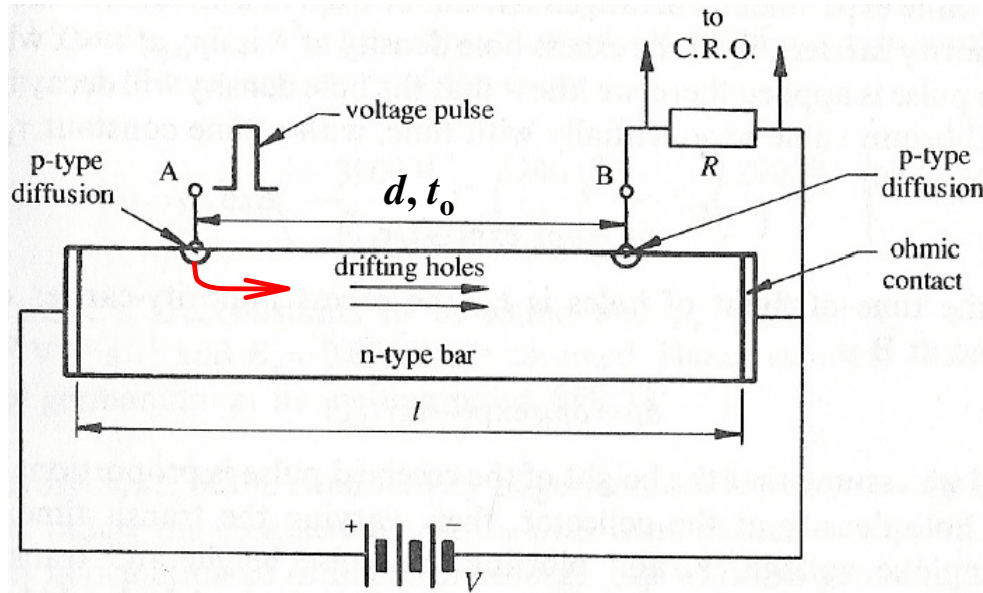
$$V = 2V_{BC} = \frac{I}{(2\pi\sigma d)}$$

$$\Rightarrow \sigma = \frac{1}{2\pi d} \frac{I}{V}$$



# Semiconductor Measurements

- Minority carrier life-time and mobility.



Mobility:

$$v_D = d/t_0 \quad \Rightarrow \quad \mu E = \mu V/l = d/t_0$$

$$\Rightarrow \quad \boxed{\mu = ld/t_0 V}$$

Life-time:

$$d(\delta p)/dt = -\delta p/\tau_{Lh} \quad \Rightarrow \quad \delta p = \delta p_0 \exp(-t/\tau_{Lh})$$