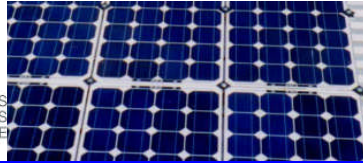




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UNIVERSIDAD DE CHILE



FI33A ELECTROMAGNETISMO

Clase 22

Campos Variables en el Tiempo-II

LUIS S. VARGAS
Area de Energía
Departamento de Ingeniería Eléctrica
Universidad de Chile



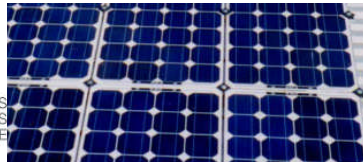
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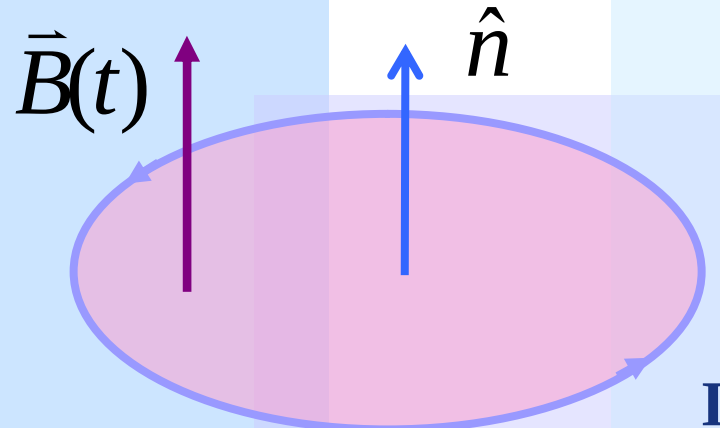
INDICE

- Ley de Faraday-Lenz
- 3ª ecuación de Maxwell
- Inductancia propia
- Inductancia mutua
- Corriente de Desplazamiento



Ley de Faraday-Lenz

Un campo magnético variable genera o induce una FEM

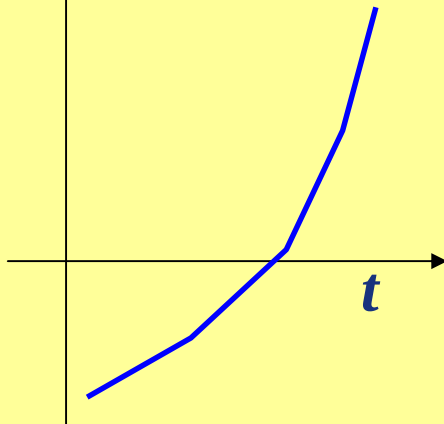


$$\varepsilon = -\frac{\partial \phi}{\partial t}$$

con

$$\phi = \iint_S \vec{B} \cdot d\vec{S}$$

$\phi(t)$



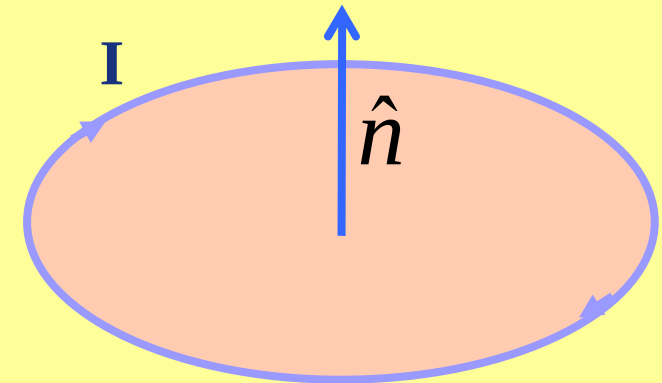
$$\dot{\phi} > 0 \Rightarrow \varepsilon < 0$$

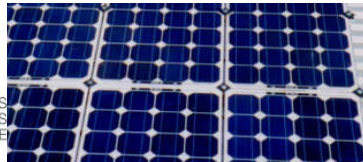
$\vec{B}$

crece

$\Rightarrow$

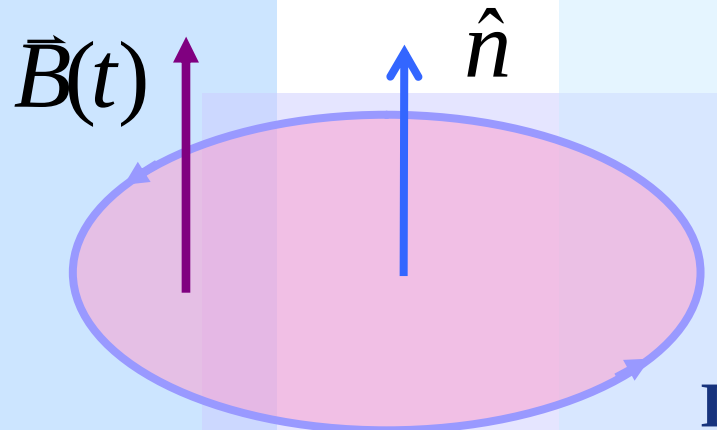
Corriente genera campo
opuesto al crecimiento





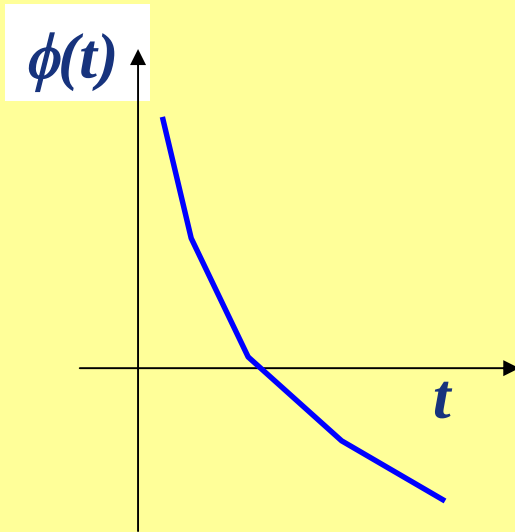
Ley de Faraday-Lenz

Un campo magnético variable genera o induce un FEM

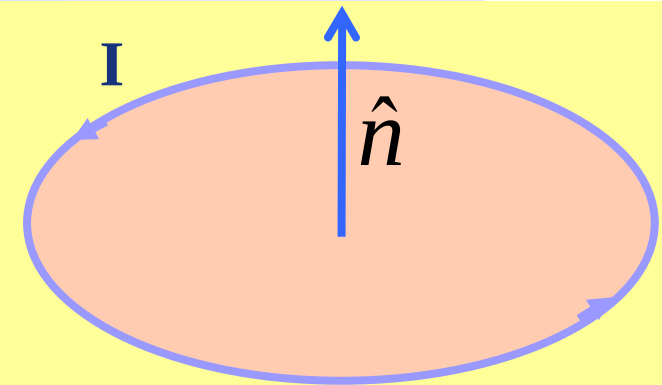


$$\varepsilon = - \frac{\partial \phi}{\partial t}$$

con $\phi = \iint_S \vec{B} \cdot d\vec{s}$

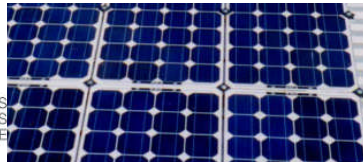


$$\dot{\phi} < 0 \Rightarrow \varepsilon > 0$$



$\vec{B}$ decrece

Corriente genera campo opuesto al decrecimiento



Modificación 3ª Ecuación de Maxwell

$$\Rightarrow \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

3ª Ecuación de Maxwell

$$\Rightarrow \nabla \times \left(\vec{E} + \frac{\partial \vec{A}}{\partial t} \right) = 0$$

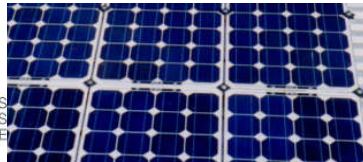
usando $\nabla \times (\nabla V) = 0$

$$\vec{E} + \frac{\partial \vec{A}}{\partial t} = -\nabla V$$

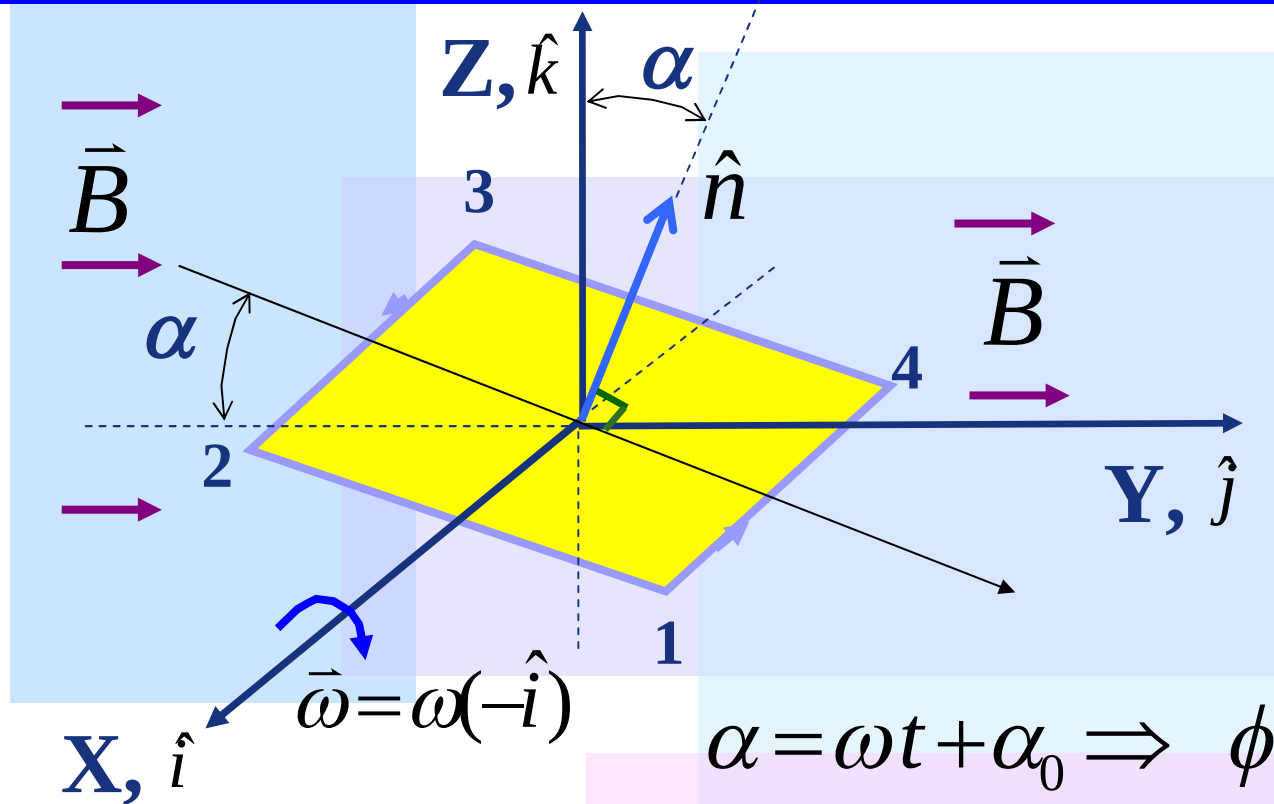
Origen
electrostático

$$\vec{E} = -\overbrace{\nabla V} - \underbrace{\frac{\partial \vec{A}}{\partial t}}$$

Debido a campo magnético
variable en el tiempo



Principio del generador



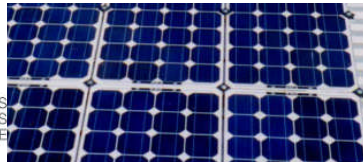
$$\phi = \iint_{S(t)} \vec{B} \cdot d\vec{s}$$

$$\phi = BA \sin \alpha$$

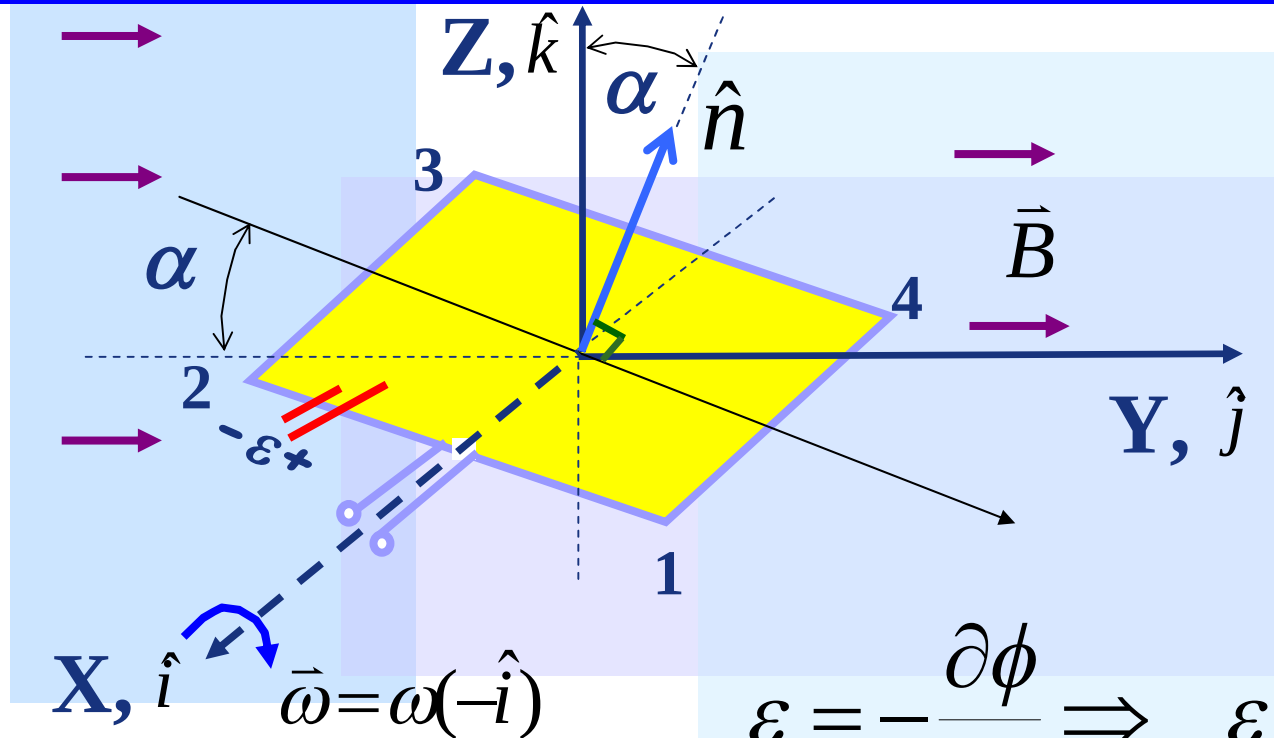
$$\alpha = \omega t + \alpha_0 \Rightarrow \phi = BA \sin(\omega t + \alpha_0)$$

**Ley de
Faraday-Lenz**

$$\varepsilon = -\frac{\partial \phi}{\partial t} \Rightarrow \varepsilon = B\omega \cos(\omega t + \alpha_0)$$



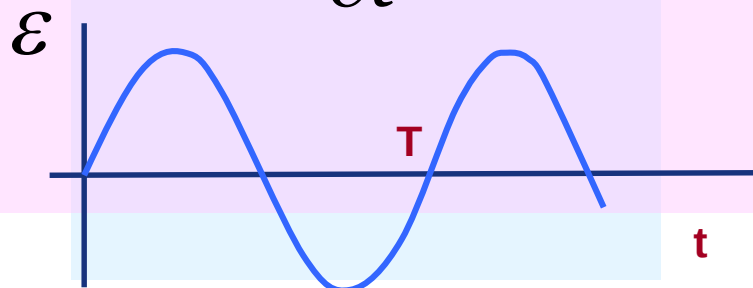
Principio del generador



$$\phi = \iint_{S(t)} \vec{B} \cdot d\vec{s}$$

$$\phi = B \sin \alpha$$

$$\varepsilon = -\frac{\partial \phi}{\partial t} \Rightarrow \varepsilon = B\omega \cos(\omega t + \alpha_0)$$

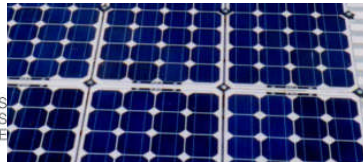


$$\omega = 2\pi f \Rightarrow T = \frac{1}{f}$$

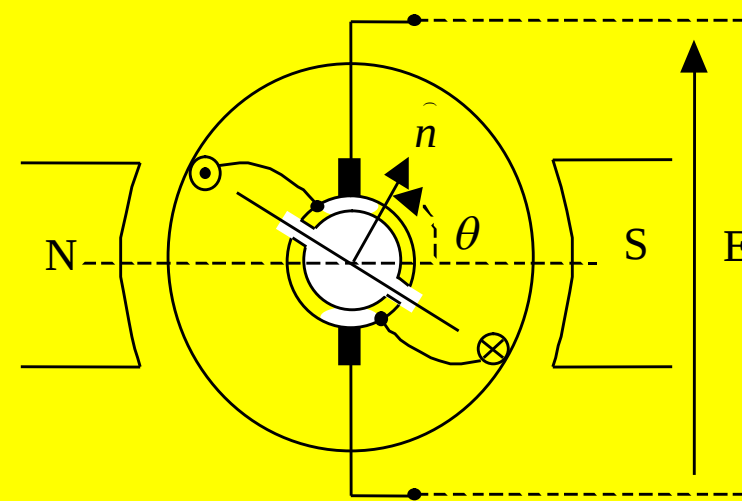
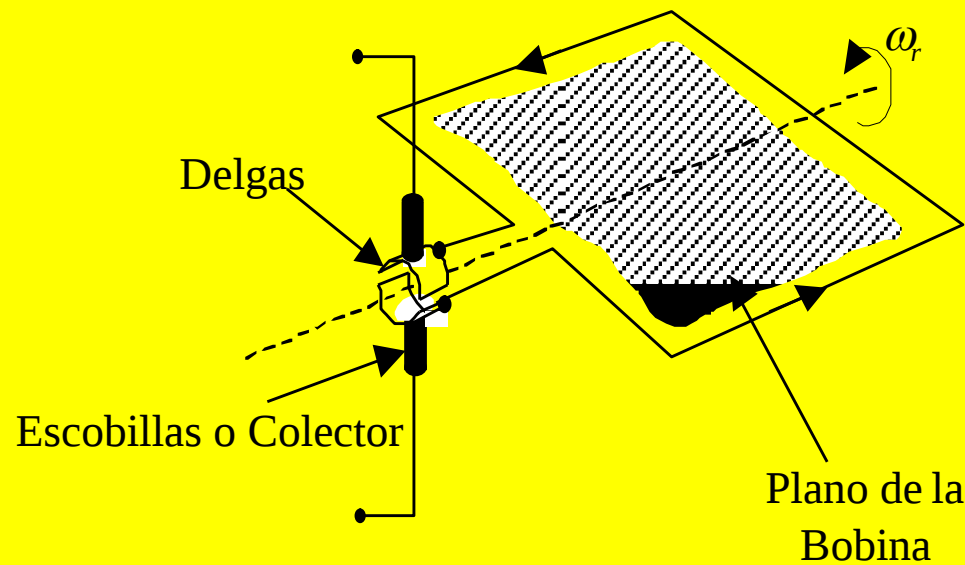


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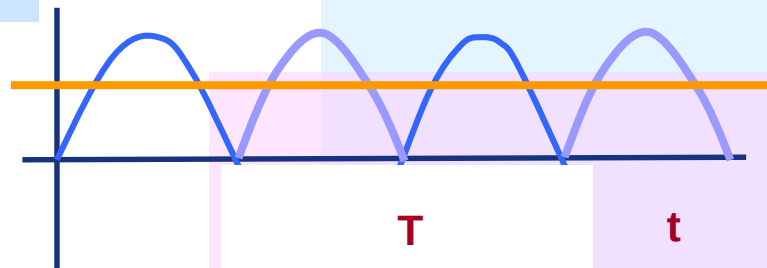
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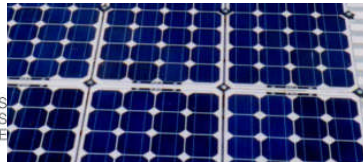
Principio del generador de Corriente Continua



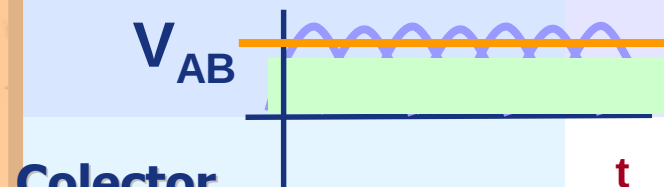
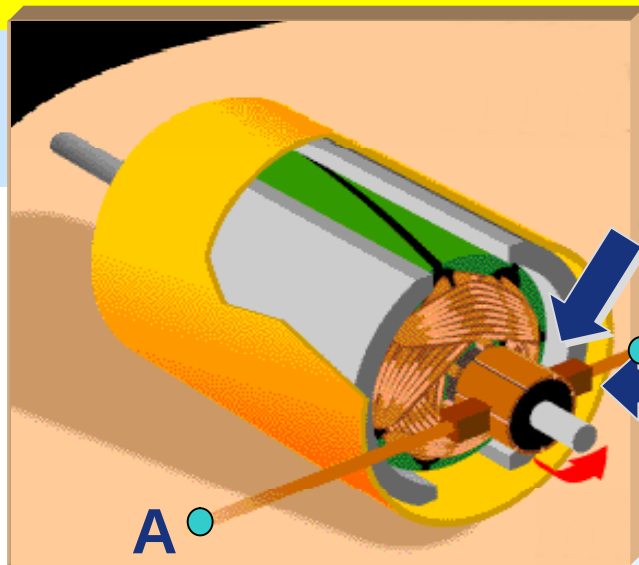
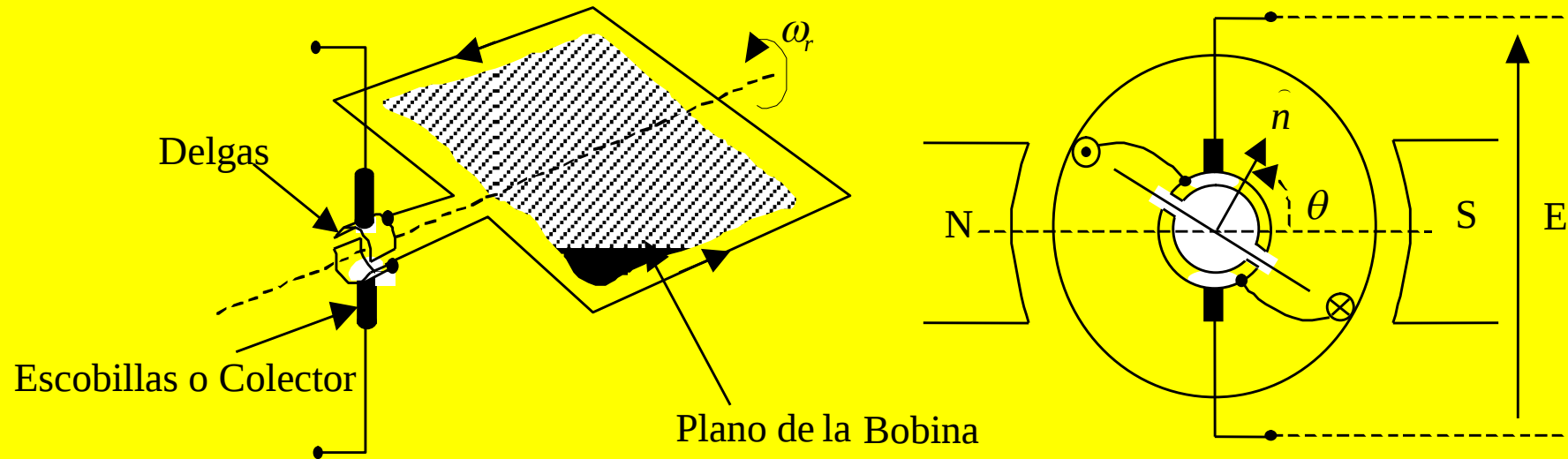
Valor
medio no
nulo



$$\omega = 2\pi f \Rightarrow T = \frac{1}{f}$$

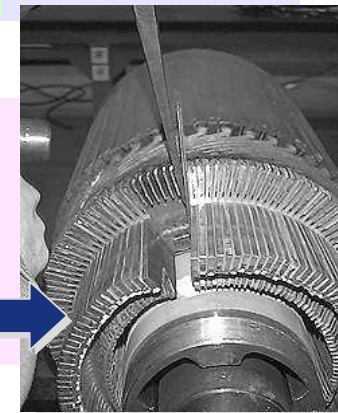


Generador de Corriente Continua



Colector
B
Escobillas

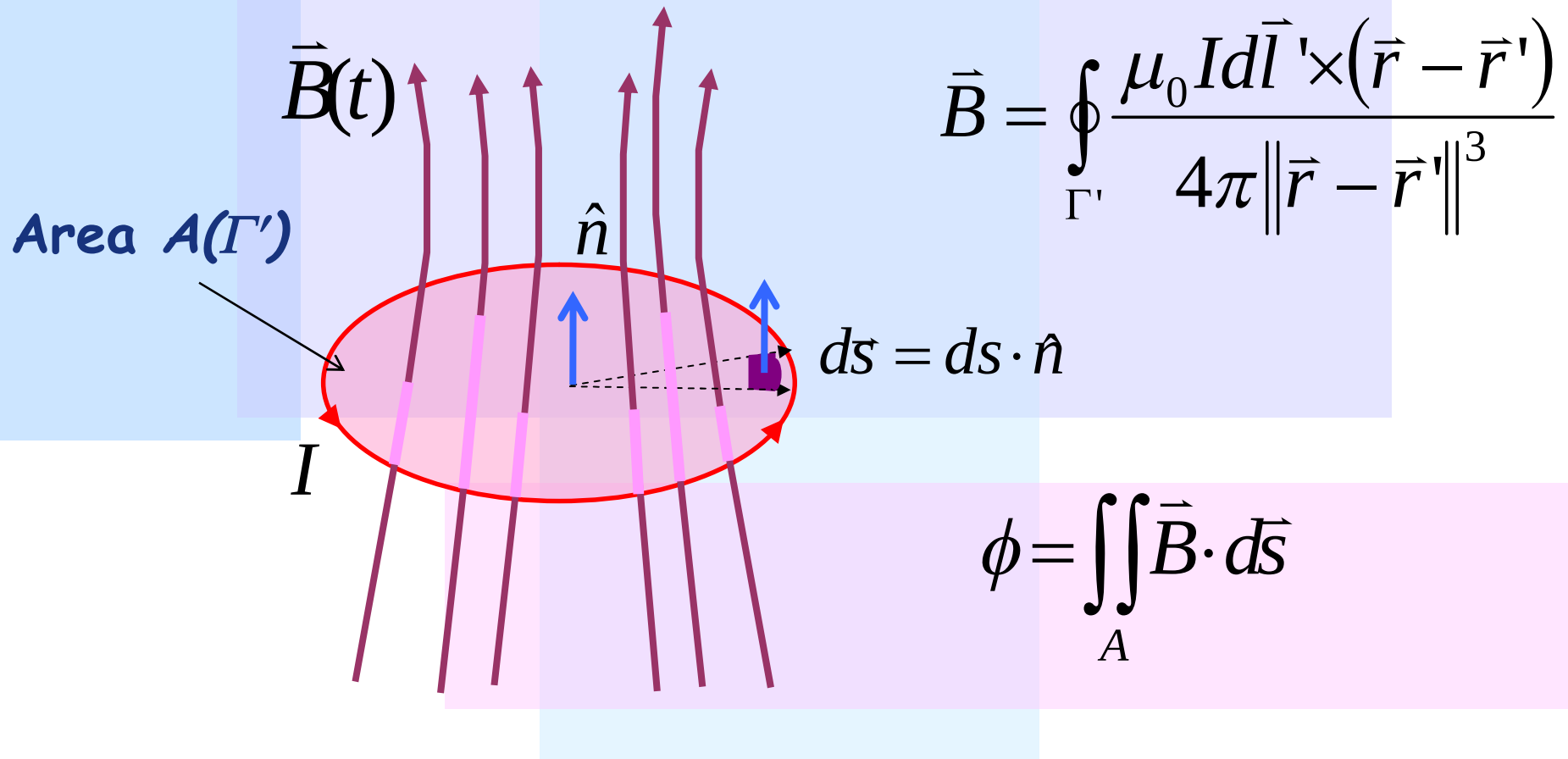
Colector real

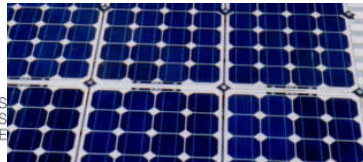




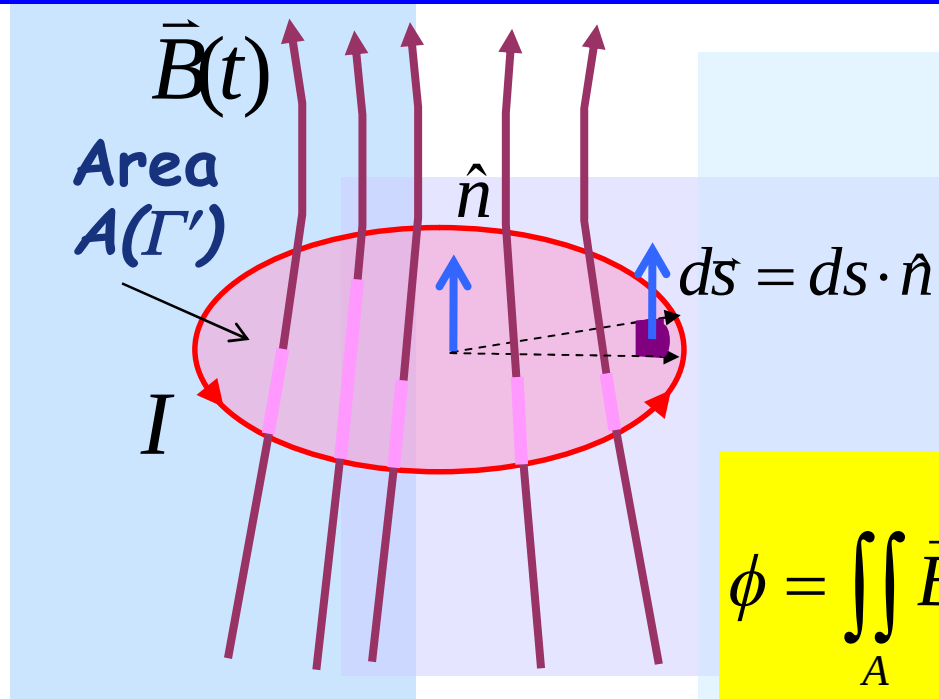
Inductancia Propia

Campo producido SOLO por I





Inductancia Propia



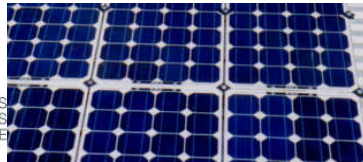
$$\vec{B} = \oint_{\Gamma'} \frac{\mu_0 I d\vec{l}' \times (\vec{r} - \vec{r}')}{4\pi \|\vec{r} - \vec{r}'\|^3}$$

$$\phi = \iint_A \vec{B} \cdot d\vec{s}$$

$$\phi = \iint_A \vec{B} \cdot d\vec{s} = \iint_A \left(\oint_{\Gamma'} \frac{\mu_0 I d\vec{l}' \times (\vec{r} - \vec{r}')}{4\pi \|\vec{r} - \vec{r}'\|^3} \right) \cdot d\vec{s}$$

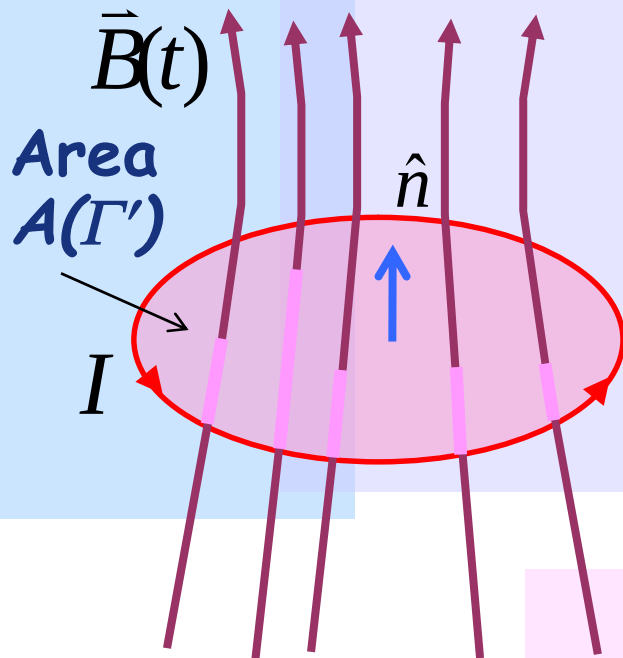
$$L \equiv \frac{\phi}{I} = \iint_A \left(\oint_{\Gamma'} \frac{\mu_0 d\vec{l}' \times (\vec{r} - \vec{r}')}{4\pi \|\vec{r} - \vec{r}'\|^3} \right) \cdot d\vec{s}$$

**Inductancia
propia del
circuito**



Inductancia Propia

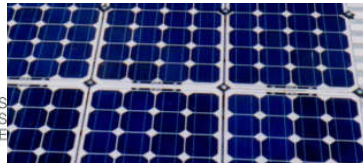
Campo producido SOLO por I



Inductancia propia del circuito

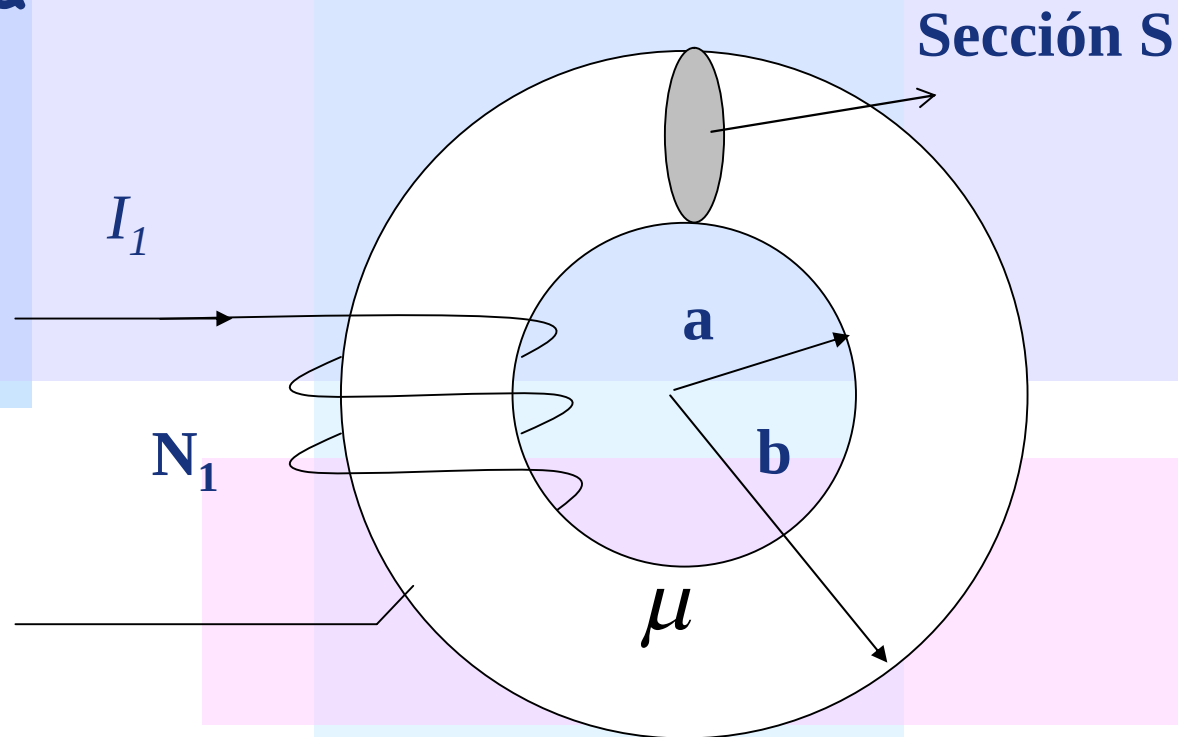
$$L \equiv \frac{\phi}{I} = \iint_A \left(\oint_{\Gamma'} \frac{\mu_0 d\vec{l}' \times (\vec{r} - \vec{r}')}{4\pi \|\vec{r} - \vec{r}'\|^3} \right) \bullet d\vec{s}$$

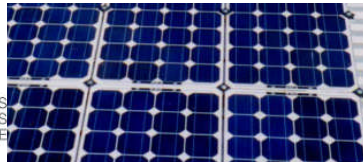
- NO depende de la corriente
- ni del flujo,
- Depende de la geometría
- $[L] = \text{Henry [H]}$



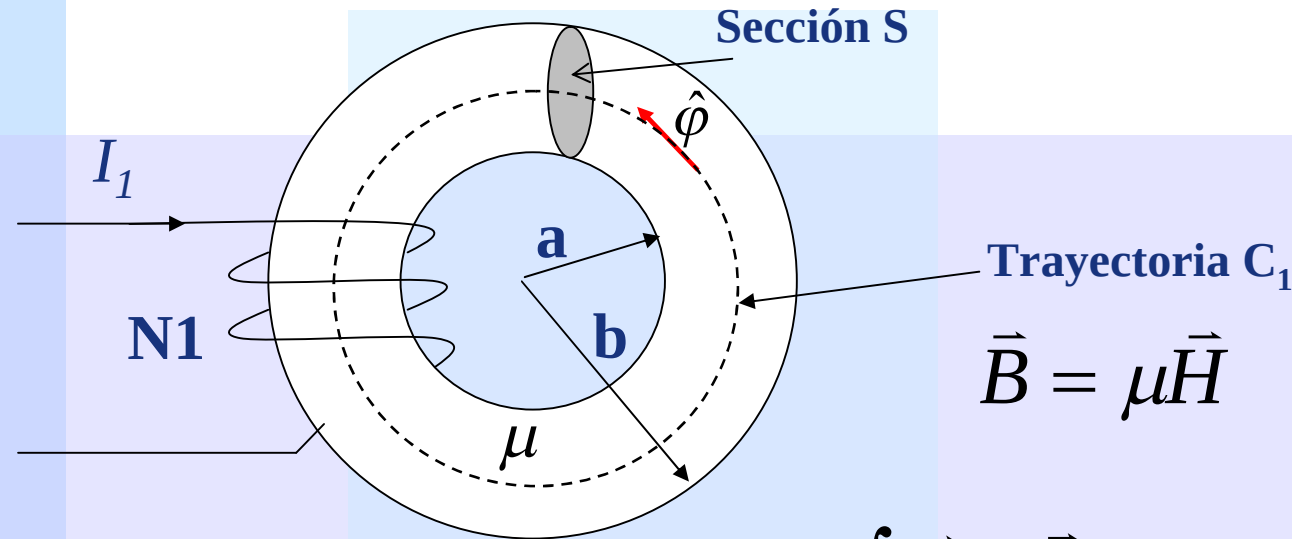
Inductancia propia

Ejemplo 1. Calcular la inductancia propia de la bobina de N_1 vueltas montada en el toroide de la figura





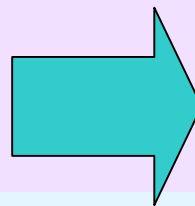
Inductancia Propia



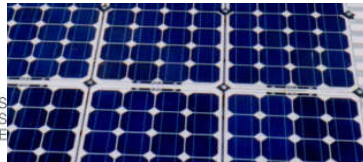
Calculamos el campo magnético

$$\oint_{C_1} \vec{H} \cdot d\vec{l} = I_{\text{enlazada}}$$

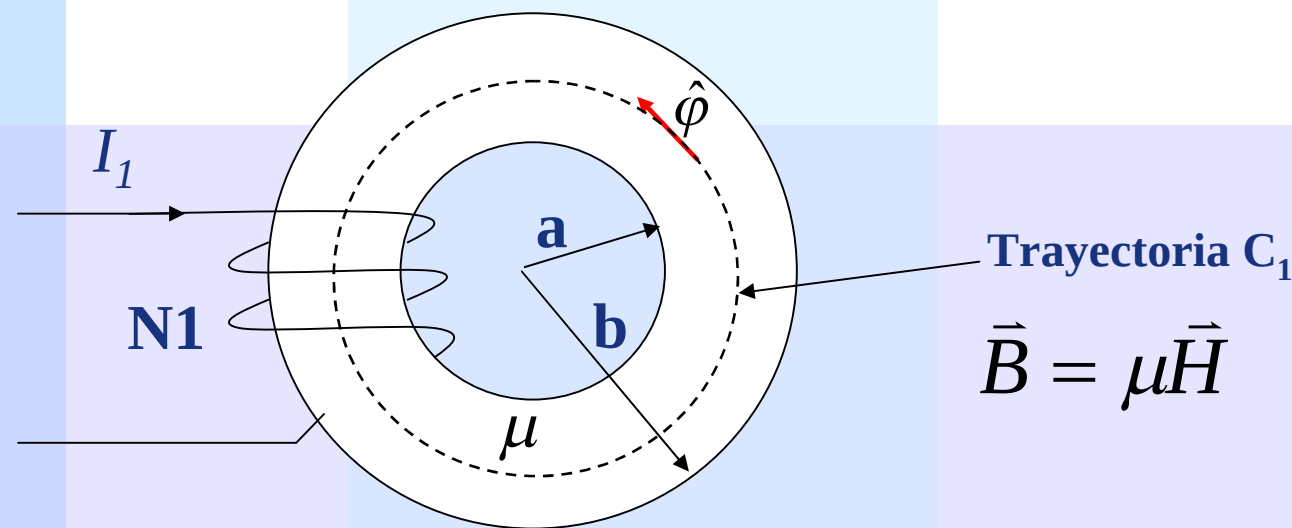
Vemos que $\vec{H} = H\hat{\phi}$



$$\oint_C \vec{H} \cdot d\vec{l} = \int_0^{2\pi} H(r)\hat{\phi} \cdot r d\phi \hat{\phi}$$
$$\oint_C \vec{H} \cdot d\vec{l} = 2\pi r H(r)$$



Inductancia Propia



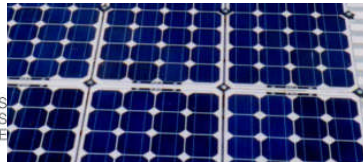
Trayectoria C_1

$$\vec{B} = \mu \vec{H}$$

Corriente total enlazada $I_{\text{enlazada}} = -N_1 I_1$

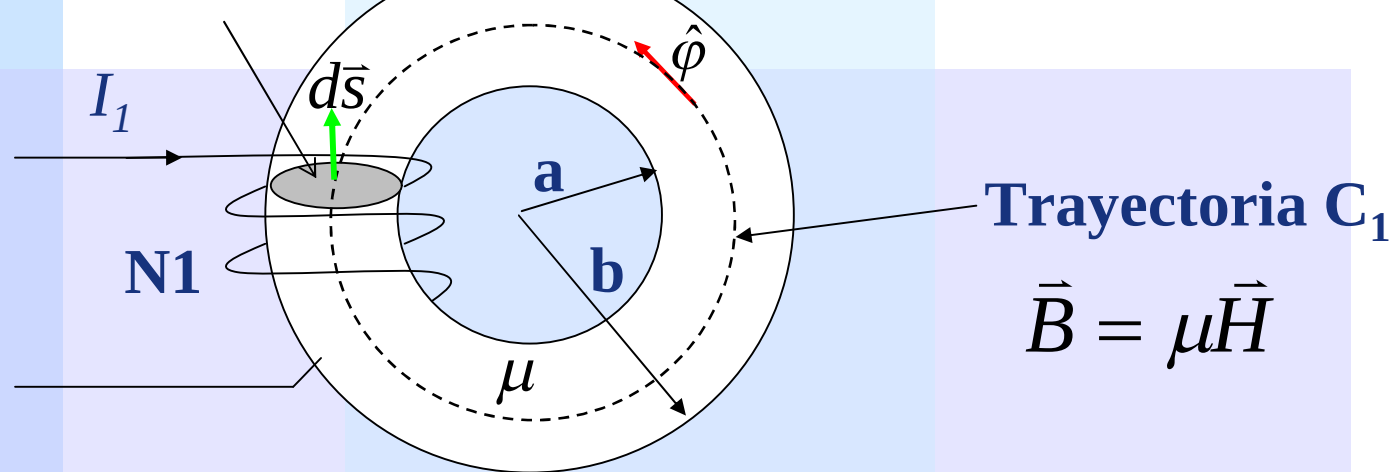
$$\Rightarrow r H(r) 2\pi = -N_1 I_1 \Rightarrow \vec{H}(r) = -\frac{N_1 I_1}{2\pi r} \hat{\phi}$$

En el punto medio $\vec{H} = -\frac{N_1 I_1}{\pi(a+b)} \hat{\phi} \Rightarrow \vec{B} = -\frac{\mu N_1 I_1}{\pi(a+b)} \hat{\phi}$



Inductancia Propia

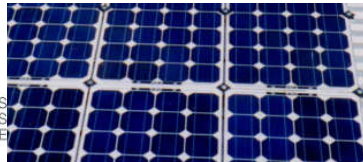
Sección A



Flujo enlazado por una vuelta $\phi = \iint_A \vec{B} \cdot d\vec{s}$

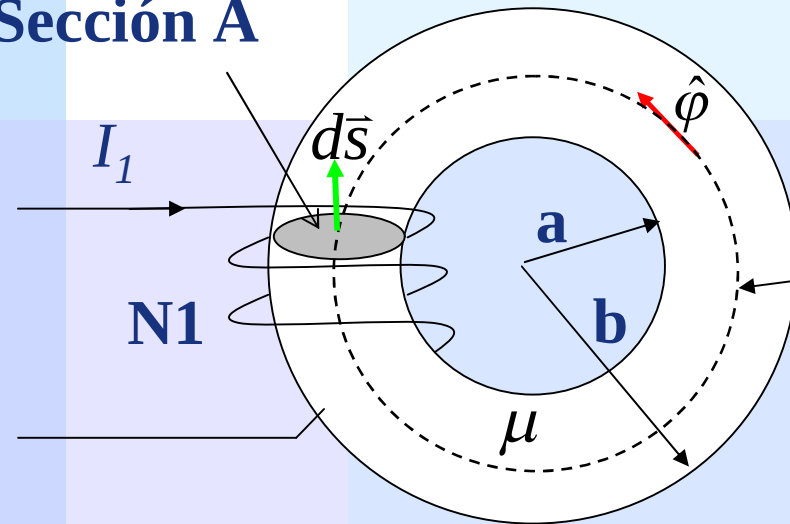
De la figura $d\vec{s} = ds(-\hat{\phi})$

$$\Rightarrow \iint_A \vec{B} \cdot d\vec{s} = \iint_A -\frac{\mu N_1 I_1}{\pi(a+b)} \hat{\phi} \cdot ds(-\hat{\phi})$$



Inductancia Propia

Sección A



Trayectoria C_1

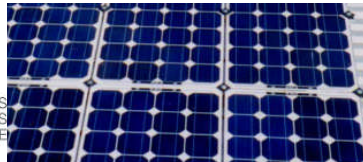
$$\vec{B} = \mu \vec{H}$$

Flujo enlazado por una vuelta

$$\phi = \frac{\mu N_1 I_1 A}{\pi(a+b)}$$

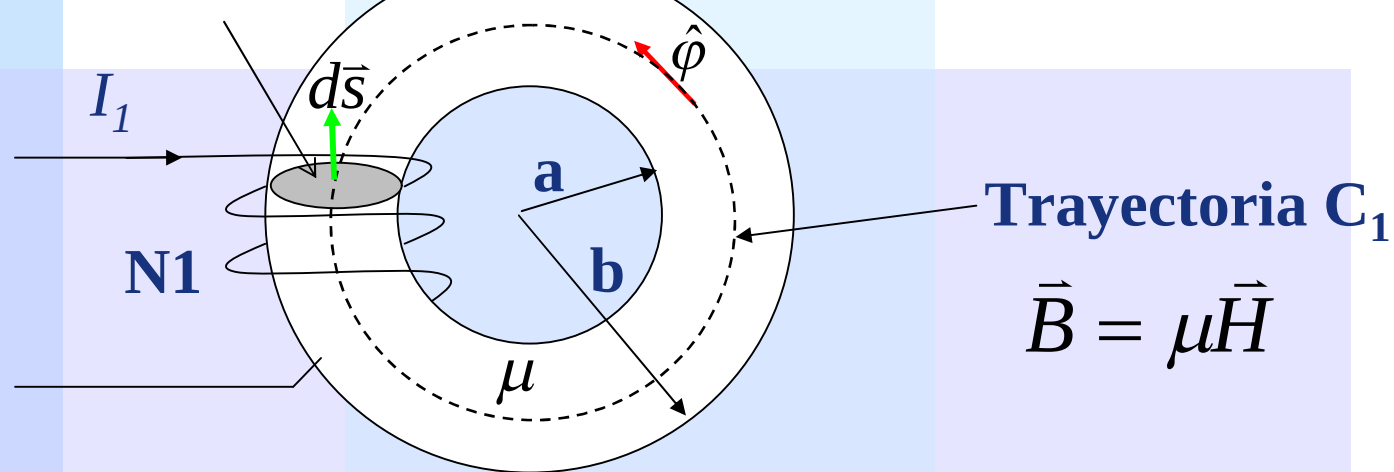
Flujo enlazado por todo el circuito

$$\phi_T = N_1 \frac{\mu N_1 I_1 A}{\pi(a+b)}$$



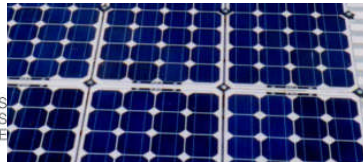
Inductancia Propia

Sección A



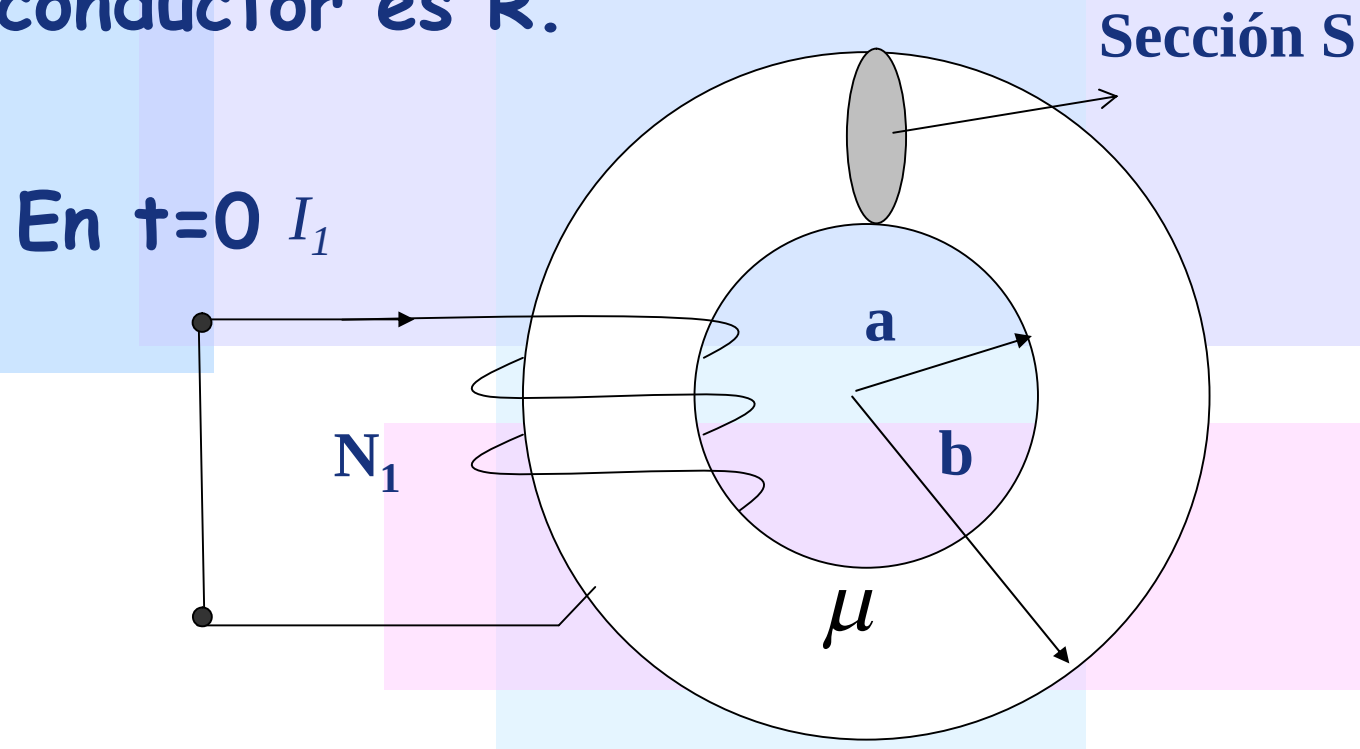
Flujo enlazado por todo el circuito $\phi_T = N_1 \frac{\mu N_1 I_1 A}{\pi(a+b)}$

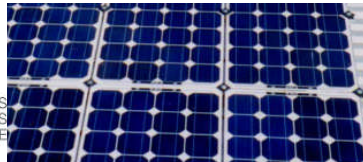
Inductancia propia del circuito $L \equiv \frac{\phi_T}{I_1} = \frac{\mu N_1^2 A}{\pi(a+b)}$



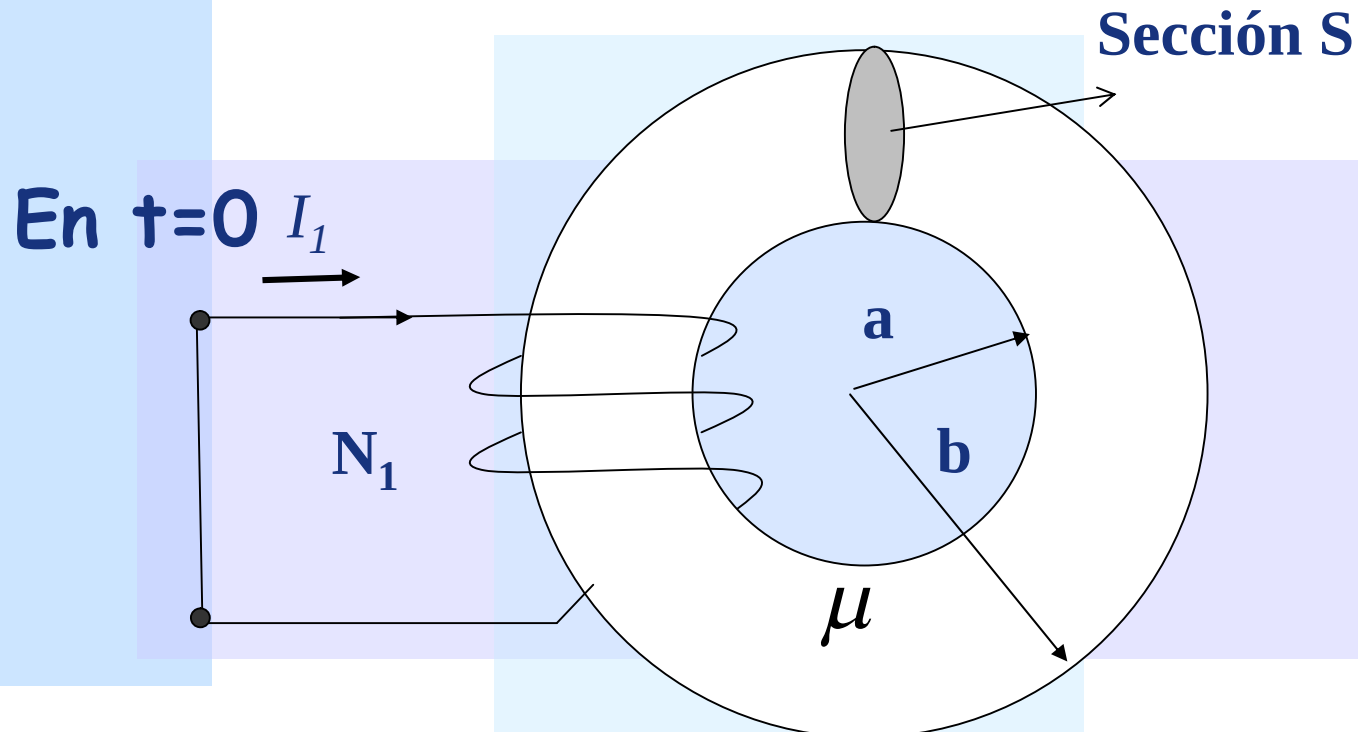
Inductancia propia

Suponga ahora que cuando circulaba una corriente I_1 se cortocircuita el conductor. Se pide calcular la corriente en función del tiempo si la resistencia del conductor es R .



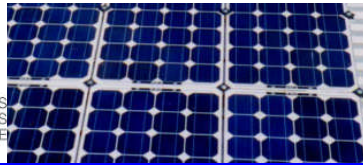


Inductancia propia

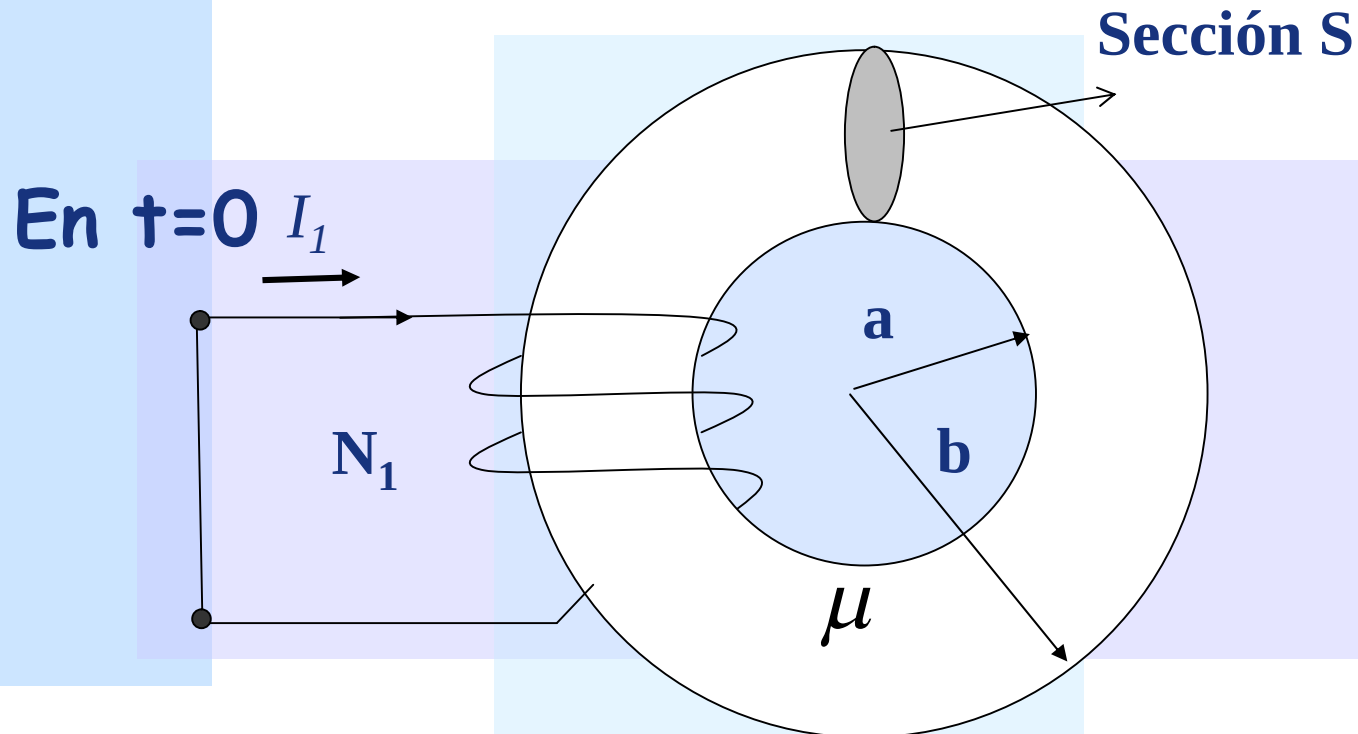


Se cumple $\phi = LI_1$ donde $L = \frac{\mu N_1^2 A}{\pi(a+b)}$

Si hay fem inducida cumple con $\varepsilon = -\frac{\partial \phi}{\partial t}$

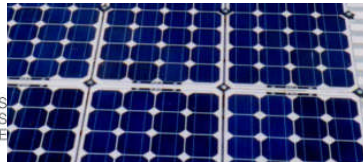


Inductancia propia

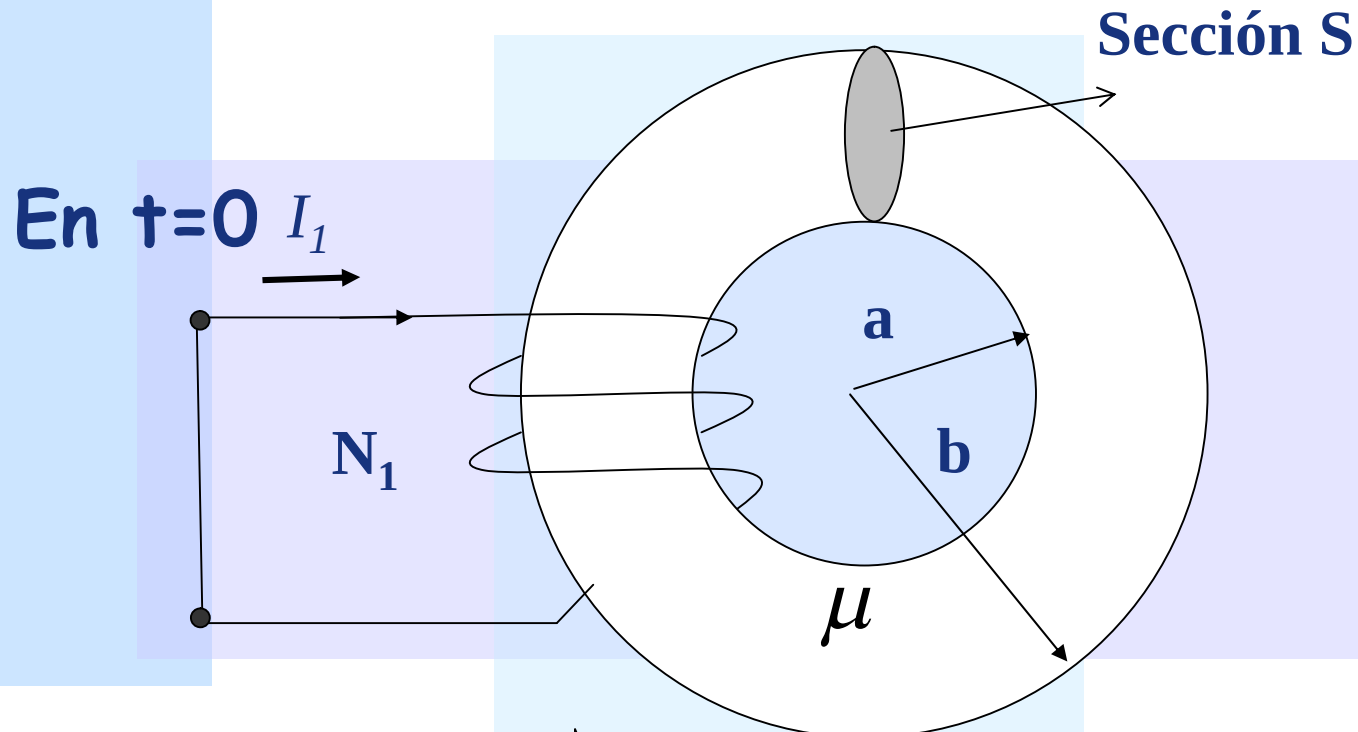


Si hay fem inducida cumple con $\varepsilon = -L \frac{\partial I}{\partial t}$

Por otra parte la resistencia impone que $\varepsilon = RI$



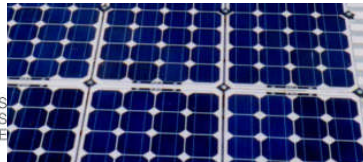
Inductancia propia



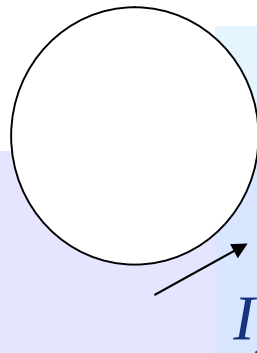
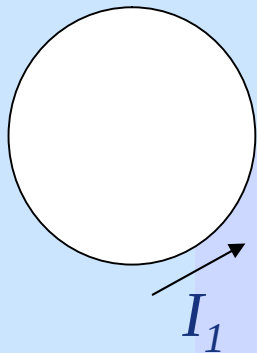
$$\Rightarrow -L \frac{\partial I}{\partial t} = RI \quad \Rightarrow \frac{\partial I}{\partial t} + \frac{R}{L} I = 0 \quad \Rightarrow I(t) = Ae^{-t/\tau}$$
$$\tau = L/R$$

En $t=0$ $I(t=0) = I_1$

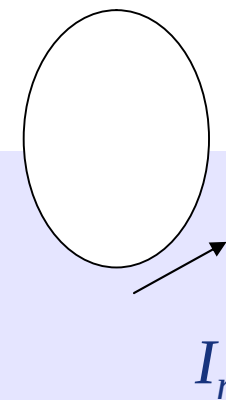
$$\therefore I(t) = I_1 e^{-t/\tau}$$



Inductancia mutua



.....



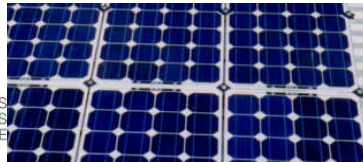
n circuitos

Sea ϕ_{jk} el flujo magnético que atraviesa el circuito j debido SOLO a la corriente que circula por el circuito k

Inductancia mutua entre el circuito j y k

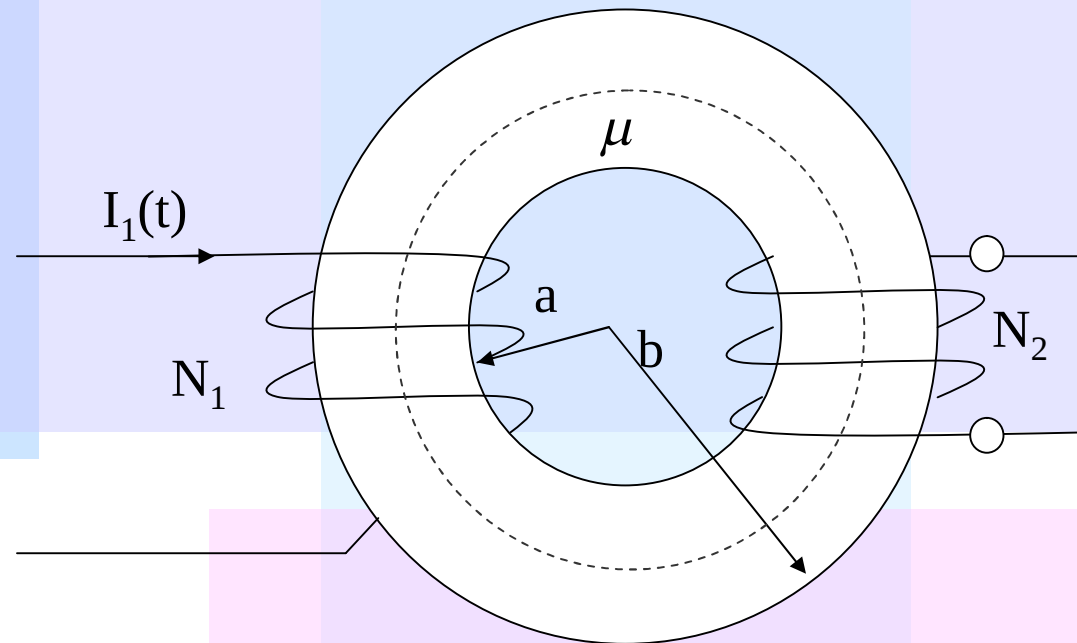
$$L_{jk} = \frac{\phi_{jk}}{I_k}$$

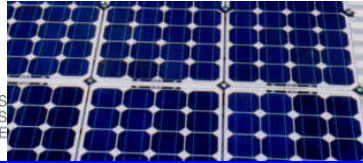
Se cumple **$L_{jk} = L_{kj}$** **PROBARLO!**



Inductancia mutua

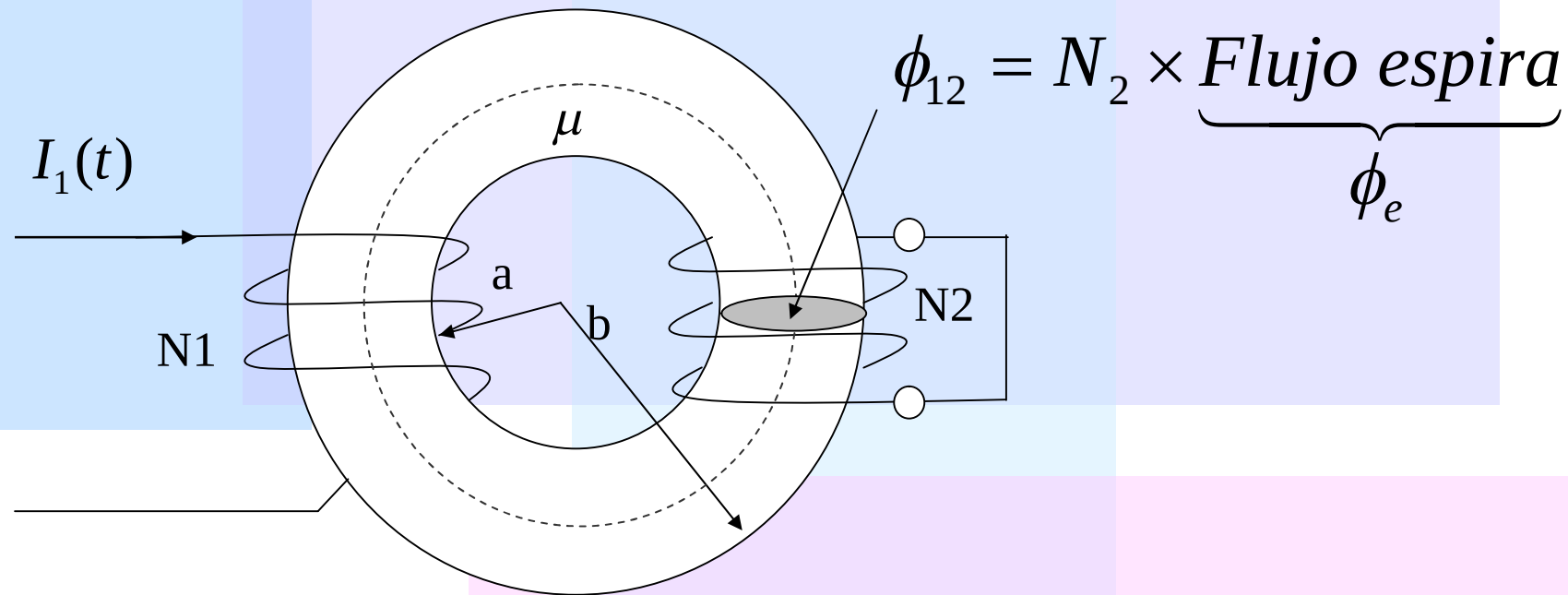
Ejemplo 2. Calcular la inductancia mutua entre los circuitos montados en el toroide de la figura



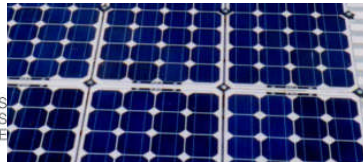


Inductancia mutua

Ejemplo 2. Calcular la inductancia mutua entre los circuitos montados en el toroide de la figura

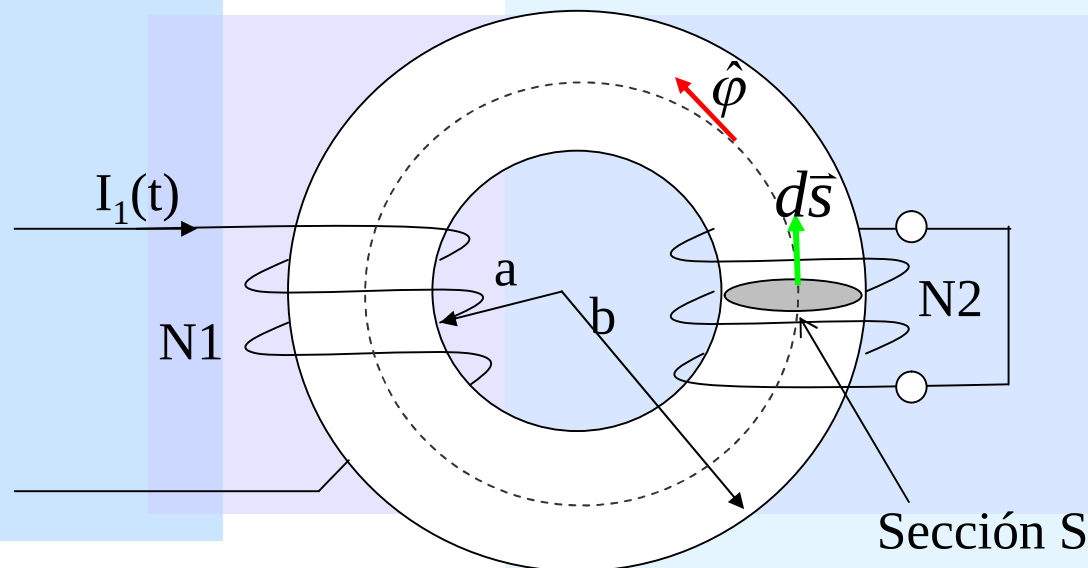


Inductancia mutua entre el circuito 1 y 2 $L_{12} \equiv \frac{\phi_{12}}{I_1}$



Inductancia mutua

Campo producido por I_1 es $\vec{B} = -\frac{\mu N_1 I_1}{\pi(a+b)} \hat{\phi}$

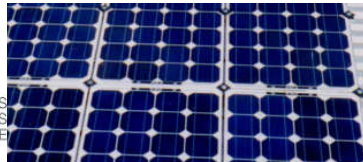


Flujo en S es

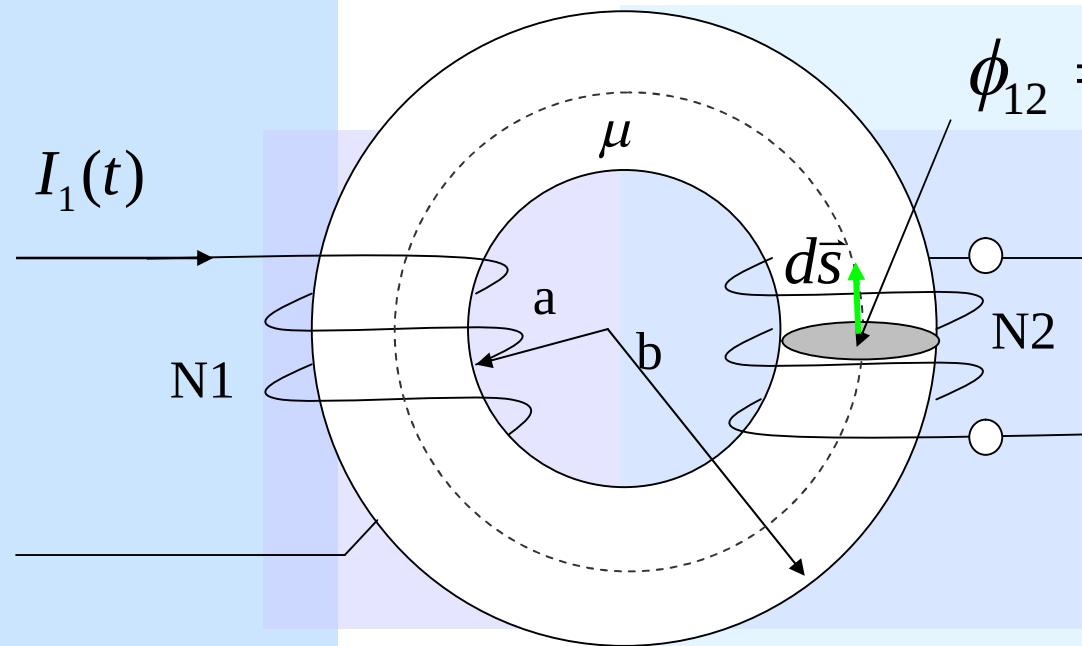
$$\phi_e = \iint_S \vec{B} \cdot d\vec{s}$$

De la figura $d\vec{s} = ds \hat{\phi}$

$$\Rightarrow \iint_S \vec{B} \cdot d\vec{s} = \iint_S -\frac{\mu N_1 I_1}{\pi(a+b)} \hat{\phi} \cdot ds \hat{\phi} \quad \therefore \phi_e = -\frac{\mu N_1 I_1 S}{\pi(a+b)}$$



Inductancia mutua

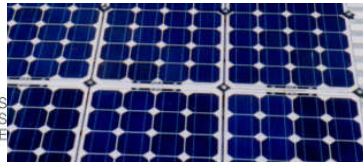


$$\phi_{12} = N_2 \times \underbrace{\text{Flujo espira}}_{\phi_e}$$

$$\phi_{12} = -\frac{\mu N_2 N_1 I_1 S}{\pi(a+b)}$$

Inductancia mutua entre el circuito 1 y 2 $L_{12} \equiv \frac{\phi_{12}}{I_1}$

$$\therefore L_{12} = -\frac{\mu N_2 N_1 S}{\pi(a+b)}$$



Corriente de Desplazamiento

$$\nabla \times \vec{H} = \vec{J}$$

4ª Ecuación de Maxwell

$$\nabla \cdot (\nabla \times \vec{H}) = \nabla \cdot \vec{J}$$

Tomando la divergencia

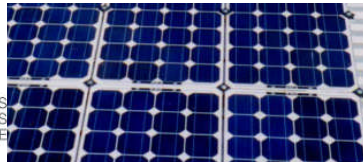
Nulo

$$0 = \nabla \cdot \vec{J}$$

Pero por la ecuación de
continuidad se cumple

$$\nabla \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0$$

Luego hay una contradicción !?#@



Corriente de Desplazamiento

Cuando no hay corriente $\vec{J} = 0$, pero SI hay campos eléctricos variables se encuentra experimentalmente la relación

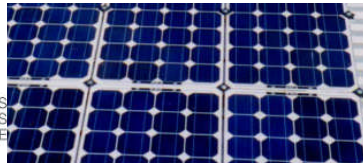
$$\nabla \times \vec{H} = \frac{\partial \vec{D}}{\partial t}$$

Así, el termino $\frac{\partial \vec{D}}{\partial t}$ debe sumarse a la 4ª ecuación, lo que da finalmente:

$$\nabla \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$$

4ª Ecuación de Maxwell

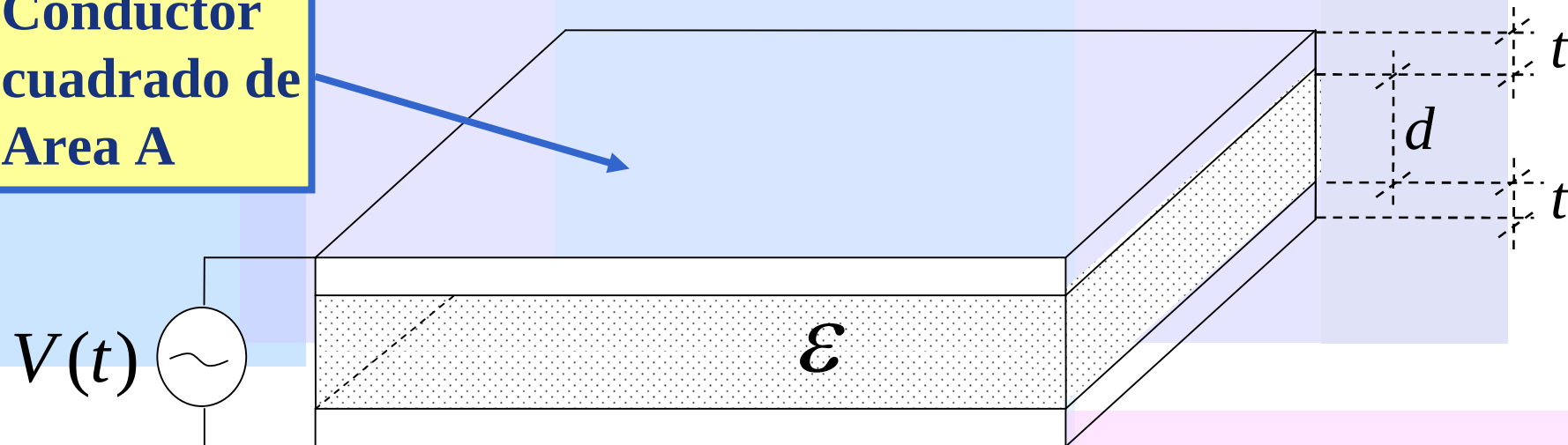
Corriente de
desplazamiento



Corriente de Desplazamiento

Ejemplo 3. Calcular corriente de desplazamiento en el interior del condensador alimentado por senoide

**Conductor
cuadrado de
Area A**

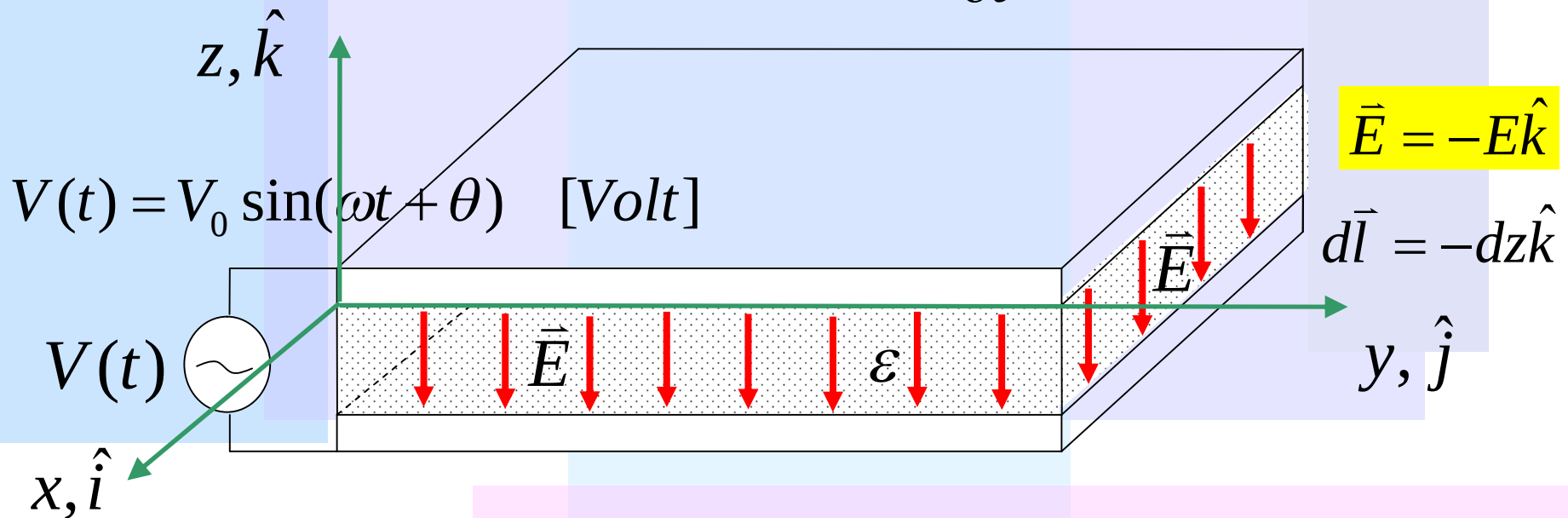


$$V(t) = V_0 \sin(\omega t + \theta) \quad [\text{Volt}]$$

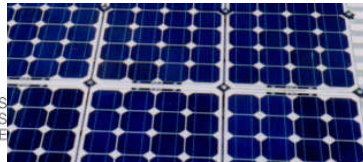


Corriente de Desplazamiento

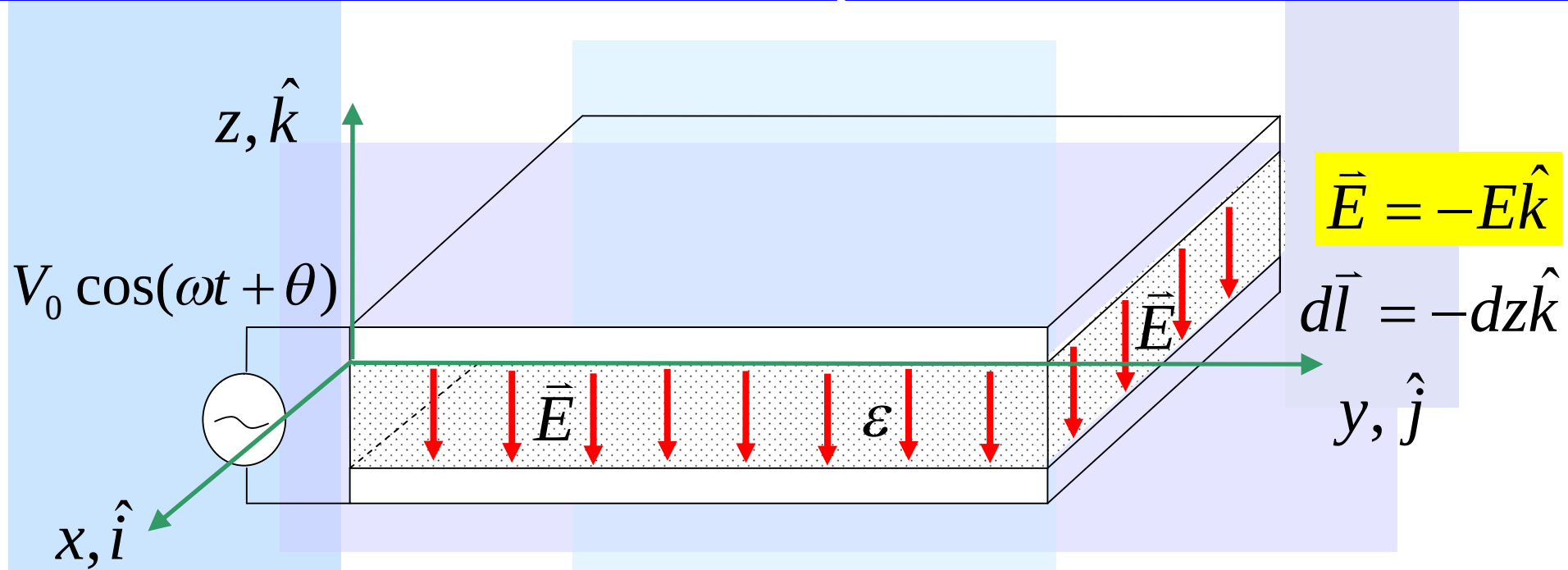
Corriente de desplazamiento $\frac{\partial \vec{D}}{\partial t}$ usamos $\vec{D} = \epsilon \vec{E}$



Además $V(z_1) - V(z_2) = \int_{z=z_1}^{z=z_2} \vec{E} \cdot d\vec{l} \Rightarrow V(t) = - \int_{z=0}^{z=-d} \vec{E} \cdot d\vec{l} = Ed$

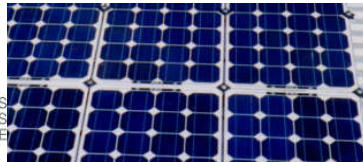


Corriente de Desplazamiento

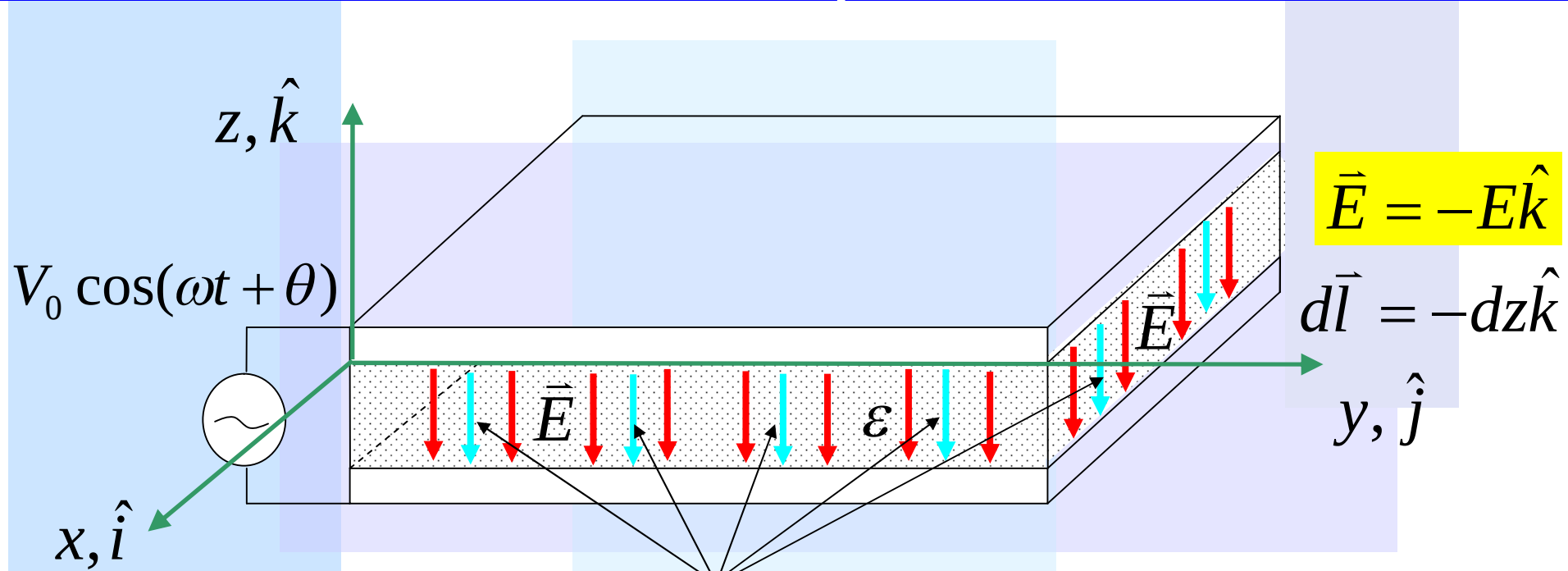


$$\Rightarrow \vec{E} = -\frac{V_0}{d} \sin(\omega t + \theta) \hat{k} \Rightarrow \vec{D} = -\frac{\epsilon V_0}{d} \sin(\omega t + \theta) \hat{k}$$

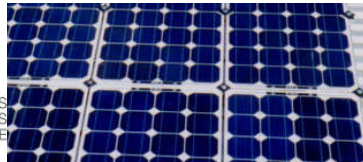
$$\therefore \frac{\partial \vec{D}}{\partial t} = -\frac{\omega \epsilon V_0}{d} \cos(\omega t + \theta) \hat{k}$$



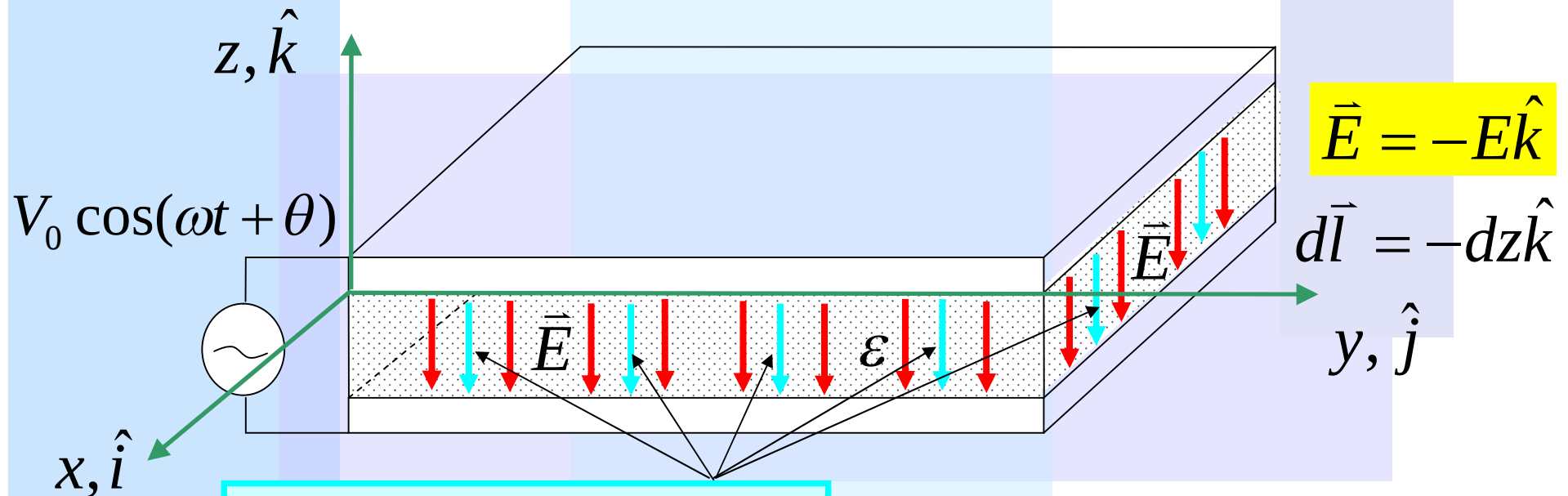
Corriente de Desplazamiento



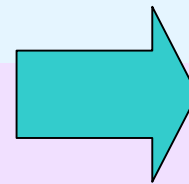
$$\frac{\partial \vec{D}}{\partial t} = -\frac{\omega \epsilon V_0}{d} \cos(\omega t + \theta) \hat{k} \equiv \vec{J}_D$$



Corriente de Desplazamiento

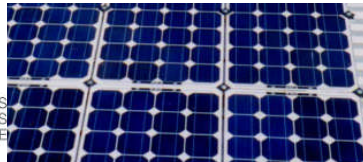


$$\vec{J}_D = -\frac{\omega \epsilon V_0}{d} \cos(\omega t + \theta) \hat{k}$$

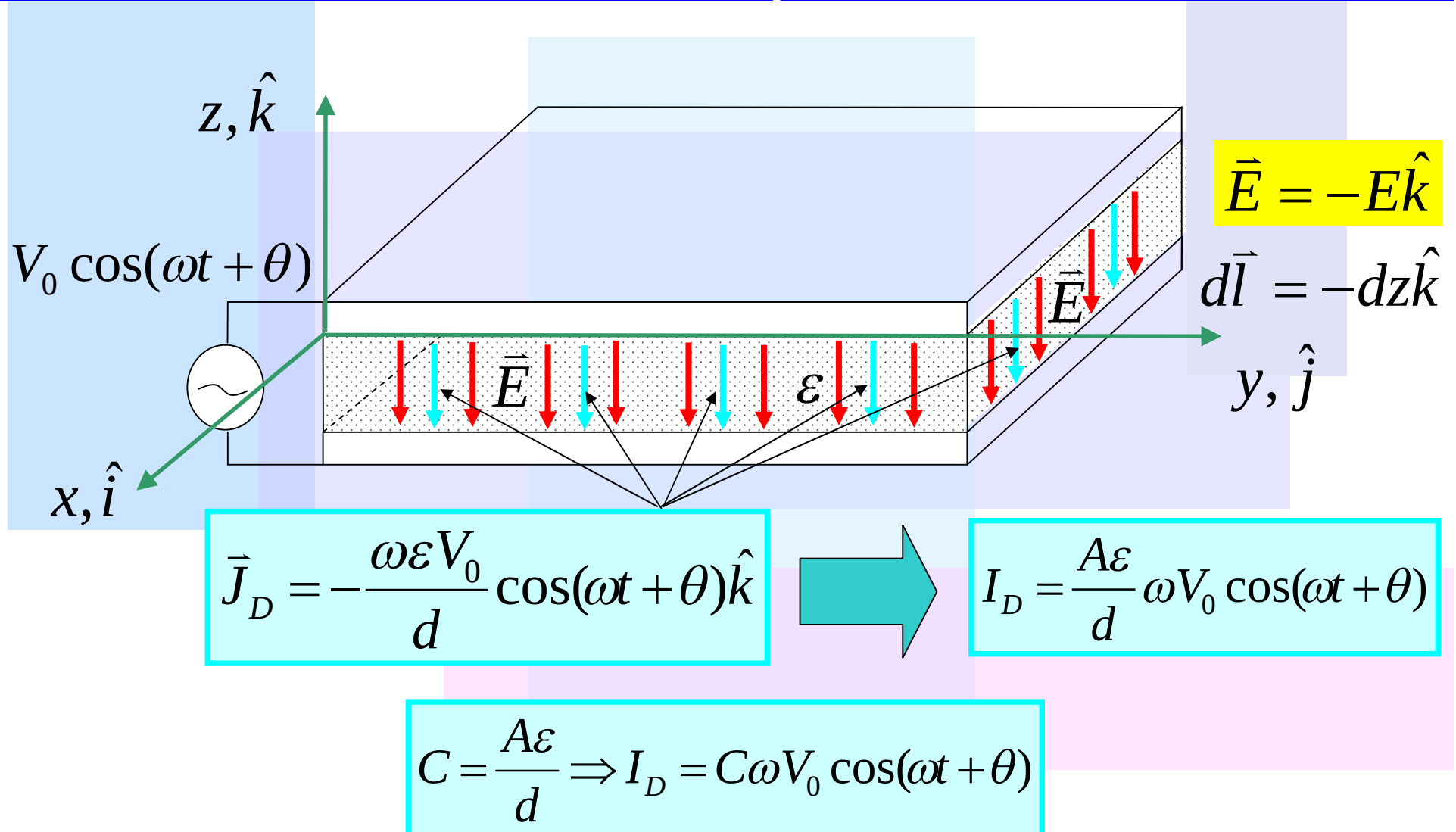


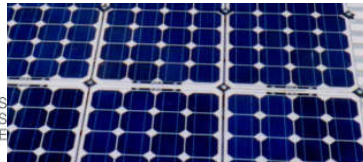
$$I_D = \iint_A \vec{J}_D \cdot d\vec{s}$$

$$I_D = \iint_A -\frac{\omega \epsilon V_0}{d} \cos(\omega t + \theta) \hat{k} \cdot d\vec{s}(-\hat{k}) = \frac{A \epsilon}{d} \omega V \cos(\omega t + \theta)$$

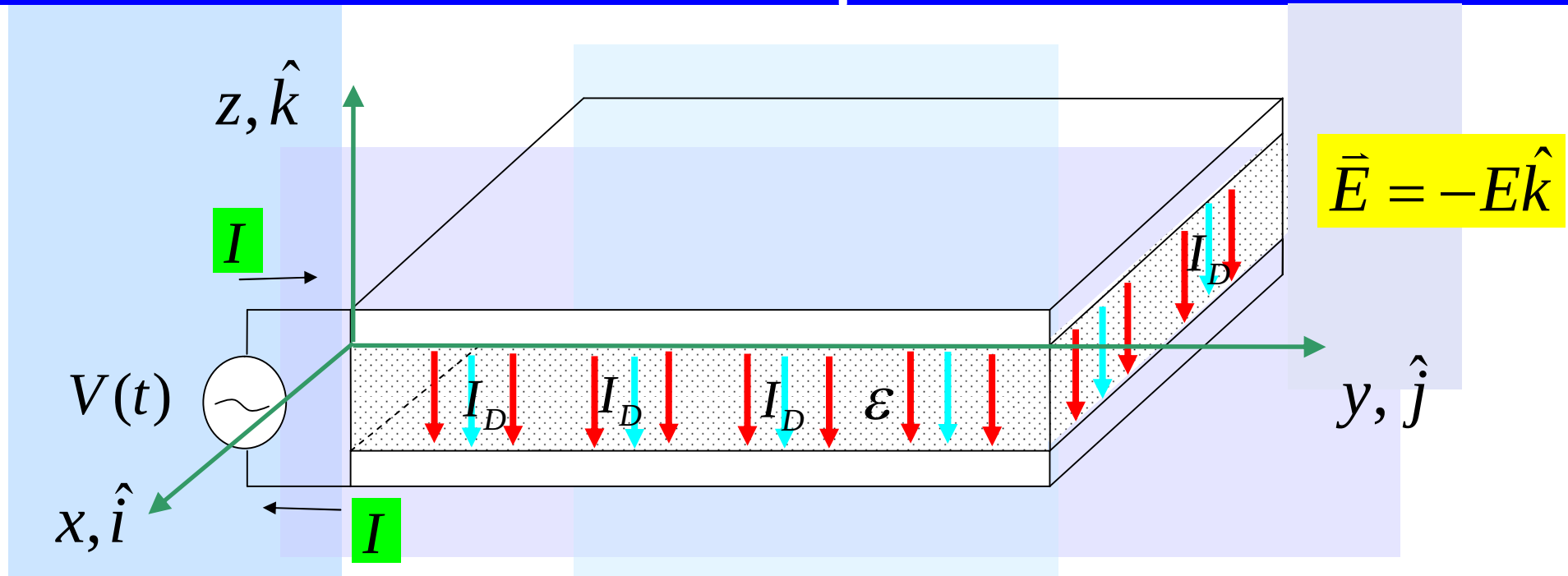


Corriente de Desplazamiento



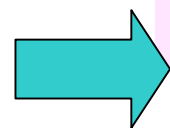


Corriente de Desplazamiento



Ecuación característica del condensador $Q = CV$

Ecuación de corriente $I = \frac{dQ}{dt} = C \frac{dV}{dt}$



$$I = C\omega V_0 \cos(\omega t + \theta)$$

$\therefore$

$$I = I_D$$