

Ec. de mas en coord θ :

$$\delta M = \sigma \delta A \delta \theta$$

$$\sigma \equiv -\frac{1}{g} \frac{\partial p}{\partial \theta}$$

Ec. de cont.

$$\frac{1}{\delta M} \frac{D}{Dt} (\delta M) = \frac{1}{\sigma \delta A \delta \theta} \frac{D}{Dt} (\sigma \delta A \delta \theta) = 0$$

$$\frac{1}{\sigma} \frac{D\sigma}{Dt} + \frac{1}{\delta A} \frac{D}{Dt} (\delta A) + \frac{1}{\delta \theta} \frac{D}{Dt} (\delta \theta) = 0$$

$$\frac{1}{\sigma} \frac{D\sigma}{Dt} + \frac{1}{\delta A} \delta \frac{DA}{Dt} + \frac{1}{\delta \theta} \delta \frac{D\theta}{Dt} = 0$$

$$\frac{1}{\sigma} \frac{D\sigma}{Dt} + \nabla_{\theta} \cdot \vec{V} + \frac{\partial \theta}{\partial \theta} = 0$$

$$\frac{D\sigma}{Dt} + \sigma \nabla_{\theta} \cdot \vec{V} - \sigma \frac{\partial \theta}{\partial \theta} = 0$$

$$\frac{\partial \sigma}{\partial t} + \vec{V} \cdot \nabla_{\theta} \sigma + \sigma \frac{\partial \sigma}{\partial \theta} + \sigma \nabla_{\theta} \cdot \vec{V} + \sigma \frac{\partial \theta}{\partial \theta} = 0$$

$$\boxed{\frac{\partial \sigma}{\partial t} + \nabla_{\theta} \cdot (\sigma \vec{V}) + \frac{\partial}{\partial \theta} (\sigma \theta) = 0}$$

Ec. de momentum

$$\boxed{\frac{\partial \vec{V}}{\partial t} + \nabla_{\theta} \cdot \left(\frac{1}{2} \vec{V} \cdot \vec{V} + \psi \right) + (\vec{z}_{\theta} + \rho) \hat{k} \times \vec{V} = -\theta \frac{\partial \vec{V}}{\partial \theta} + k \cdot \nabla \times \vec{F}_r}$$

donde $\nabla_{\theta} = \hat{i} \frac{\partial}{\partial x}|_{\theta} + \hat{j} \frac{\partial}{\partial y}|_{\theta}$

$$\vec{z}_{\theta} = \hat{k} \cdot \nabla_{\theta} \times \vec{V}$$

$$\psi = c_p T + \phi = \phi + \pi \theta \quad \text{con } \pi = c_p \frac{T}{\theta}$$

Ec. hidrostática:

$$\boxed{\frac{\partial \psi}{\partial \theta} = \pi}$$

$$T = \theta \left(\frac{p}{p_0} \right)^{\kappa/c_p}$$

ψ : fn corriente de Montgomery

π : fn de Exner

Estos dos forman un sistema cerrado en $\vec{V}$, σ , ψ y p
si se conocen $\dot{\theta}$ y $\vec{F}_r$.

Ecuación de vorticidad:

Tomando $\hat{k} \cdot \nabla_\theta \times$ de la ec. de movimiento resulta

$$(1) \quad \boxed{\frac{D}{Dt} (f + \zeta_0) + (f + \zeta_0) \nabla_\theta \cdot \vec{V} = \hat{k} \cdot \nabla_\theta \times (\vec{F}_r - \dot{\theta} \frac{\partial \vec{V}}{\partial \theta})}$$

donde $\frac{D}{Dt} = \frac{\partial}{\partial t} + \vec{V} \cdot \nabla_\theta$ es la derivada siguiendo el movimiento horizontal sobre una superficie isentrópica.

Como $\frac{\partial}{\partial t} \left(\frac{1}{\sigma} \right) = -\frac{1}{\sigma^2} \frac{\partial \sigma}{\partial t}$

la ec. de continuidad puede ser escrita, multiplicándola por σ^2 :

$$\frac{1}{\sigma^2} \frac{\partial \sigma}{\partial t} + \frac{1}{\sigma^2} \nabla_\theta \cdot (\sigma \vec{V}) + \frac{1}{\sigma^2} \frac{\partial}{\partial \theta} (\sigma \dot{\theta}) = 0$$

$$-\frac{\partial}{\partial t} \left(\frac{1}{\sigma} \right) + \frac{1}{\sigma} \nabla_\theta \cdot \vec{V} + \frac{1}{\sigma^2} \nabla_\theta \sigma \cdot \vec{V} = -\frac{1}{\sigma^2} \frac{\partial}{\partial \theta} (\sigma \dot{\theta})$$

$$-\frac{\partial}{\partial t} \left(\frac{1}{\sigma} \right) - \vec{V} \cdot \nabla_\theta \left(\frac{1}{\sigma} \right) + \frac{1}{\sigma} \nabla_\theta \cdot \vec{V} = -\frac{1}{\sigma^2} \frac{\partial}{\partial \theta} (\sigma \dot{\theta})$$

$$(2) \quad \frac{D}{Dt} \left(\frac{1}{\sigma} \right) - \frac{1}{\sigma} \nabla_\theta \cdot \vec{V} = \frac{1}{\sigma^2} \frac{\partial}{\partial \theta} (\sigma \dot{\theta})$$

Entonces multiplicando $\frac{1}{\sigma} \times (1) + (f + \zeta_0) \times (2)$:

$$\frac{D}{Dt} \left(\frac{f + \zeta_0}{\sigma} \right) = \frac{\partial}{\partial t} \left(\frac{f + \zeta_0}{\sigma} \right) + \vec{V} \cdot \nabla \left(\frac{f + \zeta_0}{\sigma} \right) = \frac{f + \zeta_0}{\sigma^2} \frac{\partial}{\partial \theta} (\sigma \dot{\theta}) + \frac{1}{\sigma} \hat{k} \cdot \nabla_\theta \times (\vec{F}_r - \dot{\theta} \frac{\partial \vec{V}}{\partial \theta})$$

$$\boxed{\frac{DP}{Dt} = \frac{P}{\sigma} \frac{\partial}{\partial \theta} (\sigma \dot{\theta}) + \frac{1}{\sigma} \hat{k} \cdot \nabla_\theta \times (\vec{F}_r - \dot{\theta} \frac{\partial \vec{V}}{\partial \theta})}$$

donde $P = \frac{f + \zeta_0}{\sigma} = (f + \zeta_0) \left(-g \frac{\partial \theta}{\partial p} \right)$ vorticidad (potencial de Ertel)

Notar que si el calentamiento diabático y el roce son nulos, el movimiento horizontal sobre una superficie isentrópica conserva la vortic. pot. de Ertel.

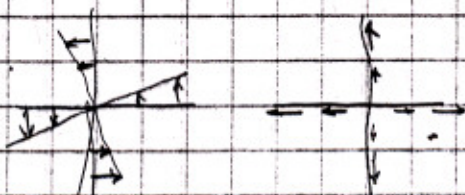
La ecuación de vorticidad (1) puede ser escrita en la forma

$$\frac{\partial \zeta_0}{\partial t} = -\nabla_\theta \cdot [(f + \zeta_0) \vec{V}] + \hat{k} \cdot \nabla_\theta \times [\vec{F}_r - \dot{\theta} \frac{\partial \vec{V}}{\partial \theta}]$$

y recordando que para un vector $\vec{A}$ arbitrario:

$$\hat{k} \cdot \nabla_0 \times \vec{A} = \nabla_0 \cdot (\vec{A} \times \hat{k})$$

se puede escribir



$$\frac{\partial \xi_0}{\partial t} = -\nabla_0 \cdot \left[(\xi_0 + f) \vec{V} - \left(\vec{F}_r - \dot{\phi} \frac{\partial \vec{V}}{\partial \phi} \right) \times \hat{k} \right]$$

expresión que muestra que la vorticidad isentrópica en un punto fijo se origina sólo en la divergencia del vector

$$(f + \xi_0) \vec{V} - \left(\vec{F}_r - \dot{\phi} \frac{\partial \vec{V}}{\partial \phi} \right) \times \hat{k}$$

Sobre una superficie isentrópica, que no intercepte la $z=0$, la vorticidad media es constante, pues al integrar sobre ella la divergencia, transformada en flujo a través de un perímetro nulo, integra a cero, dando

$$\frac{\partial}{\partial t} \oint \xi_0 dA = 0$$

$$\bar{\xi}_0 = \text{cte}$$

Además:

$$\bar{\xi} = \hat{k} \cdot \nabla \times \vec{V} = \frac{1}{a \cos \phi} \left[\frac{\partial v}{\partial \lambda} - \frac{\partial (u \cos \phi)}{\partial \phi} \right]$$

$$\frac{1}{4\pi a^2} \int_0^{2\pi} \int_{-\pi/2}^{\pi/2} \xi a^2 \cos \phi d\phi d\lambda = \frac{1}{4\pi a} \int_0^{2\pi} \int_{-\pi/2}^{\pi/2} \left[\frac{\partial v}{\partial \lambda} - \frac{\partial (u \cos \phi)}{\partial \phi} \right] d\phi d\lambda$$

$$\bar{\xi} = \frac{1}{4\pi a} \left\{ \int_{-\pi/2}^{\pi/2} \left[v \right]_{\lambda=0}^{\lambda=2\pi} d\phi - \int_0^{2\pi} \left[u \cos \phi \right]_{\phi=-\pi/2}^{\phi=\pi/2} d\lambda \right\}$$

$$\bar{\xi} = \frac{1}{4\pi a} \{ 0 - 0 \} = 0$$

La vorticidad en la $z=0$ isentrópica es redistribuida por el movimiento, pero no la crea ni la destruye.

Es en fun. de Exner: $\Pi = C_p \frac{T}{\Theta} = C_p \left(\frac{p}{p_0} \right)^{R/C_p}$

i) 1er Pr. Td:

$$Tds = dq = C_p dT - \alpha dp \quad \text{con} \quad ds = C_p d(\ln \Theta)$$

$$ds = C_p \frac{1}{\Theta} d\Theta$$

$$Tds = C_p \frac{T}{\Theta} d\Theta = \Pi d\Theta$$

$$\boxed{Tds = \Pi d\Theta} \quad \text{1er Pr. Td}$$

notando que $\Pi\Theta = C_p T$
 $d(\Pi\Theta) = C_p dT$ y el 1er Pr. Td:

$$Tds = d(\Pi\Theta) - \alpha dp$$

$$-\alpha dp = Tds - d(\Pi\Theta) = \Pi d\Theta - d(\Pi\Theta)$$

$$\boxed{\alpha dp = \Theta d\Pi}$$

ii) Ec. hidrostática

$$\alpha dp = -d\phi$$

$$\boxed{\Theta d\Pi = -d\phi} \quad \text{con}$$

$$\psi = \phi + C_p T = \phi + \Pi\Theta$$

$$\boxed{\frac{\partial \psi}{\partial \Theta} = \Pi}$$

iii) Ec. estado

$$p\alpha = RT$$

$$\text{con} \quad T = \frac{\Theta \Pi}{C_p}$$

$$\boxed{p\alpha = \frac{R}{C_p} \Pi\Theta}$$