

a) Sabemos que $\Delta S = \frac{\lambda}{T}$ con λ = calor latente de cambio de fase
 $\rightarrow T$ = temperatura en la cual se produce.

Luego, $\lambda_{\text{vap}} = T \Delta S = (100+273)(7.106 - 1.295,2) = 2,17 \cdot 10^6 \text{ J/kg}$.

b) Como P es constante,

$$\lambda_{\text{vap}} = \Delta h \quad \therefore \Delta h = h_{\text{vap}} - h_{\text{agua}} = 2,17 \cdot 10^6 \text{ J/kg} = 2,170 \text{ kJ/kg}$$

$$\therefore h_{\text{agua}} = (2.675,2 - 2,170) \text{ kJ/kg}$$

$$\therefore \boxed{h_{\text{agua}} = 505,2 \text{ kJ/kg}}$$

c) La integración de la ec. de Clausius-Clapeyron para el caso de la vaporización, con las aproximaciones $v_{\text{vap}} \gg v_{\text{agua}} \therefore v_{\text{vap}} - v_{\text{agua}} \approx v_{\text{vap}}$
y luego considerando el vapor como gas ideal: $v_{\text{vap}} = \frac{RT}{P}$, conduce a

$$\ln P = -\frac{\lambda_{\text{vap}}}{RT} + \text{cte}$$

Expresemos λ_{vap} en J/mol : 1 mol de agua = 18 g de agua $\therefore 1 \text{ g} = \frac{1}{18} \text{ mol}$

$$\therefore \lambda_{\text{vap}} = 2,17 \cdot 10^6 \text{ J/kg} = 2,17 \cdot 10^3 \text{ J/g} = 2,17 \cdot 10^3 \cdot 18 \text{ J/mol}$$

$$\therefore \boxed{\lambda_{\text{vap}} = 39.060 \text{ J/mol}}$$

$$\therefore \frac{\lambda_{\text{vap}}}{R} = \frac{39.060}{8,31} \approx 4.700 \quad \rightarrow \ln P = -\frac{4.700}{T} + \text{cte}$$

Determinamos la constante: $\ln 1 = -\frac{4.700}{373} + \text{cte} \rightarrow \boxed{\text{cte} = 12,6}$

Luego, $\boxed{\ln P = -\frac{4700}{T} + 12,6}$

d) $\ln \frac{1}{2} = -\frac{4700}{T} + 12,6 \rightarrow -0,7 - 12,6 = -\frac{4700}{T} \rightarrow T = 353,4 \text{ K}$
 $\boxed{t = 80,4^\circ\text{C}}$

$$\ln 2 = -\frac{4700}{T} + 12,6 \rightarrow T \approx 395 \text{ K} \rightarrow \boxed{t = 122^\circ\text{C}}$$

2. a) $\langle \epsilon_c \rangle = \frac{1}{2} kT$ (término cuadrático contribuye con $\frac{1}{2} kT$)

$$b) \langle \epsilon_p \rangle = \frac{\int_{-\infty}^{\infty} bx^4 e^{-\beta kx^4} dx}{\int_{-\infty}^{\infty} e^{-\beta kx^4} dx} = \frac{-\frac{\partial}{\partial \beta} \int_{-\infty}^{\infty} e^{-\beta bx^4} dx}{\int_{-\infty}^{\infty} e^{-\beta bx^4} dx} = -\frac{\partial}{\partial \beta} \ln \int_{-\infty}^{\infty} e^{-\beta bx^4} dx$$

$$\therefore \langle \epsilon_p \rangle = -\frac{\partial}{\partial \beta} \ln (kT) \cdot \beta^{-\frac{1}{4}} = -\frac{\partial}{\partial \beta} \ln(kT) + \frac{\partial}{\partial \beta} \left[\frac{1}{4} \ln \beta \right]$$

$$= \frac{1}{4} \cdot \frac{1}{\beta}$$

$$\therefore \boxed{\langle \epsilon_p \rangle = \frac{1}{4} kT}$$

c) Energía media total: $\langle \epsilon \rangle = \langle \epsilon_c \rangle + \langle \epsilon_p \rangle$

$$= \frac{1}{2} kT + \frac{1}{4} kT$$

$$\therefore \boxed{\langle \epsilon \rangle = \frac{3}{4} kT}$$

d) Para 1 mol (N_0 osciladores): $\mu = \frac{3}{4} N_0 kT$

$$\therefore C_V = \left(\frac{\partial \mu}{\partial T} \right)_V = \frac{3}{4} N_0 k \rightarrow \boxed{C_V = \frac{3}{4} R}$$

$$3.- \quad N_i = \frac{N}{Z} g_i e^{-\beta \epsilon_i} \quad (1) \quad ; \quad Z = \sum_i g_i e^{-\beta \epsilon_i} \quad ; \quad \beta = \frac{1}{kT}$$

a) En nuestro sistema, $g_1 = g_2 = 1$. Estados no degenerados.

$$\therefore Z = e^{-\beta \epsilon_1} + e^{-\beta \epsilon_2} = e^{-\beta \epsilon_1} (1 + e^{-\beta(\epsilon_2 - \epsilon_1)}).$$

$$\text{Si hacemos } \epsilon = \epsilon_2 - \epsilon_1 > 0 \rightarrow \boxed{Z = e^{-\beta \epsilon_1} (1 + e^{-\beta \epsilon})} \quad (2)$$

$$\text{Luego, } N_1 = \frac{N e^{-\beta \epsilon_1}}{e^{-\beta \epsilon_1} (1 + e^{-\beta \epsilon})} \quad \text{y} \quad N_2 = \frac{N e^{-\beta \epsilon_2}}{e^{-\beta \epsilon_1} (1 + e^{-\beta \epsilon})} = \frac{N e^{-\beta \epsilon}}{1 + e^{-\beta \epsilon}}$$

$$\therefore \boxed{P_1 = \frac{N_1}{N} = \frac{1}{1 + e^{-\beta \epsilon}}} \quad \text{y} \quad \boxed{P_2 = \frac{N_2}{N} = \frac{e^{-\beta \epsilon}}{1 + e^{-\beta \epsilon}}} \quad P_1 + P_2 = 1. \quad \checkmark$$

$$b) \quad U = N_1 \epsilon_1 + N_2 \epsilon_2 = \frac{N (\epsilon_1 + \epsilon_2 e^{-\beta \epsilon})}{1 + e^{-\beta \epsilon}} = \frac{N [\epsilon_1 + (\epsilon_1 + \epsilon) e^{-\beta \epsilon}]}{1 + e^{-\beta \epsilon}}$$

$$\therefore U = \cancel{\frac{N \epsilon_1 (1 + e^{-\beta \epsilon})}{(1 + e^{-\beta \epsilon})}} + \frac{N \epsilon e^{-\beta \epsilon}}{1 + e^{-\beta \epsilon}}$$

$$\text{pero } U = N \epsilon_1 + \frac{N e^{-\beta \epsilon} \epsilon}{1 + e^{-\beta \epsilon}} \Rightarrow \boxed{U = N \left[\epsilon_1 + \frac{\epsilon}{1 + e^{\beta \epsilon}} \right]} \quad (3)$$

Alternativamente puede obtenerse de $U = -N \frac{\partial \ln Z}{\partial \beta}$

$$c) \quad C_v = \left(\frac{\partial U}{\partial T} \right)_V = \frac{\partial U}{\partial \beta} \frac{\partial \beta}{\partial T} = -\frac{1}{kT^2} \frac{\partial U}{\partial \beta} = k \beta^2 \frac{\partial U}{\partial \beta}$$

$$\therefore C_v = \frac{k \beta^2 N \epsilon^2 e^{\beta \epsilon}}{(1 + e^{\beta \epsilon})^2} \quad \text{pero} \quad \boxed{C_v = N k \frac{\beta^2 \epsilon^2 e^{\beta \epsilon}}{(1 + e^{\beta \epsilon})^2}} \quad (4)$$