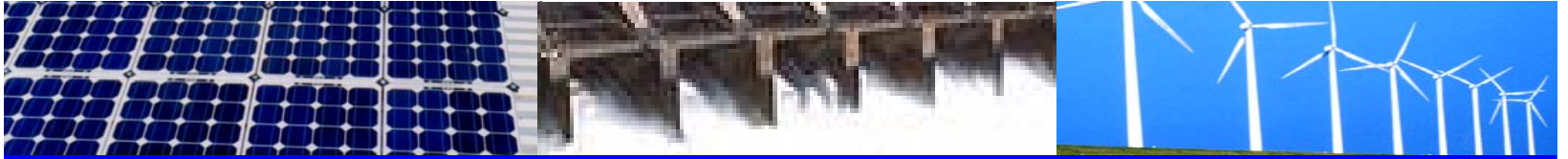




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FI33A ELECTROMAGNETISMO

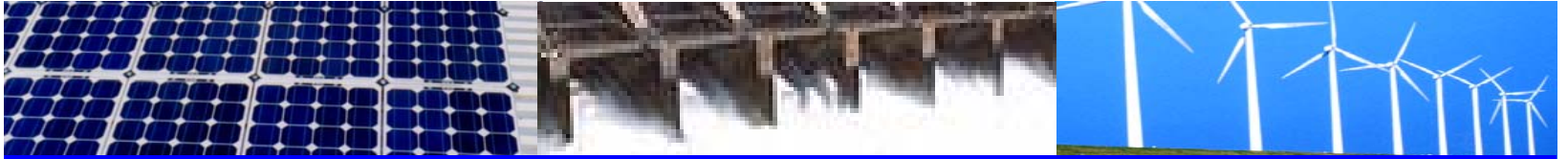
Clase 13

Corriente Eléctrica-III

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Area de Energía
Departamento de Ingeniería Eléctrica
Universidad de Chile



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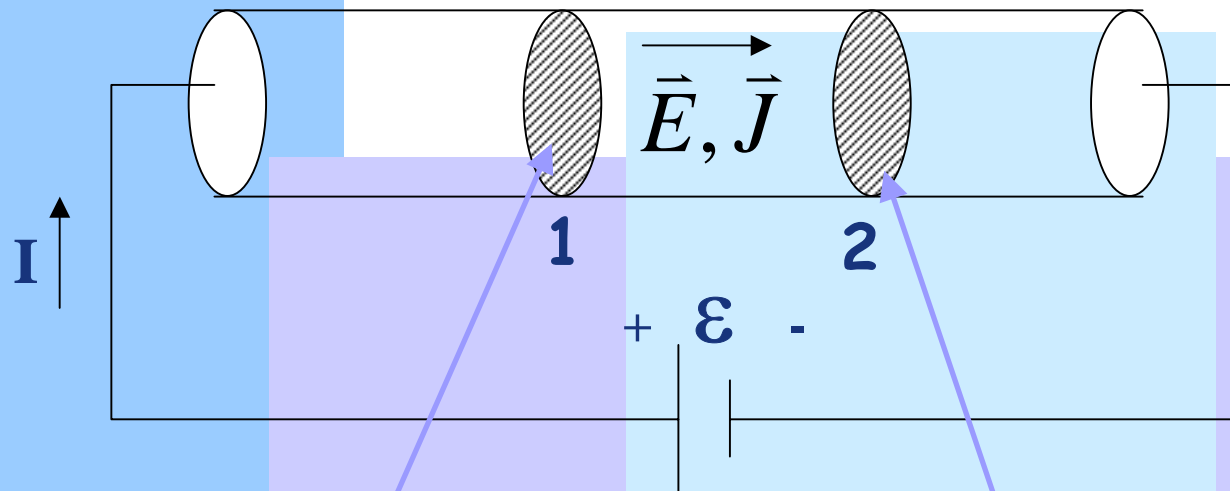


INDICE

- Efecto Joule
- Ecuación de Continuidad
- Condiciones de borde para J



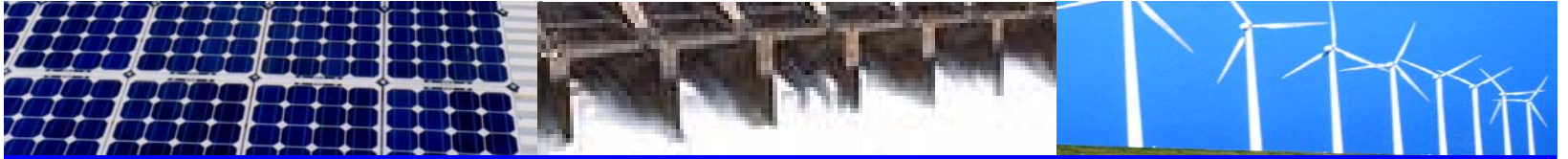
Efecto Joule



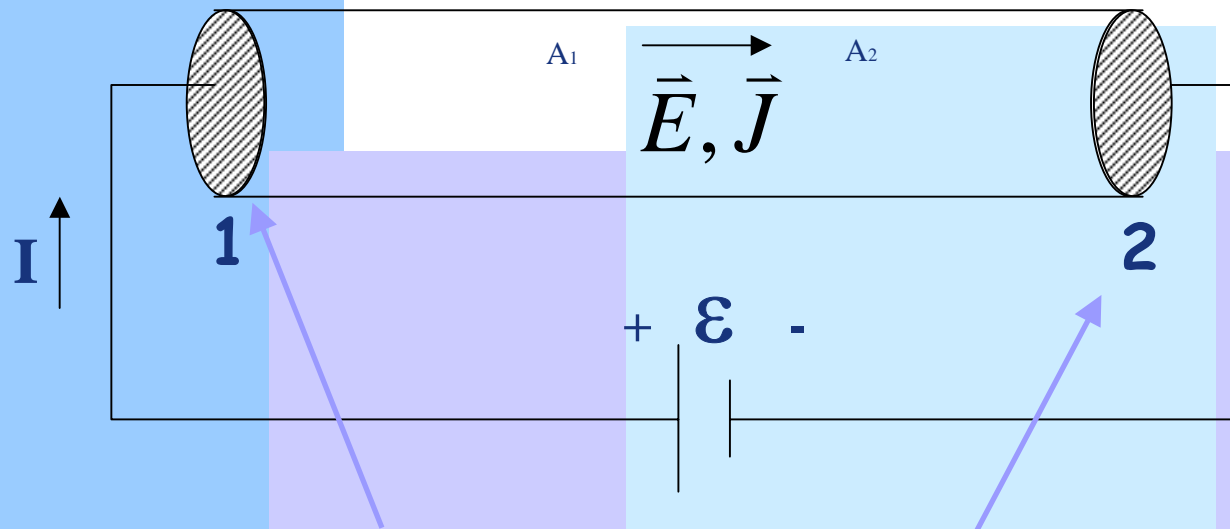
$$\Delta Q_1 = \Delta Q_2$$
$$\Delta U = \Delta Q(V_1 - V_2)$$

**Potencia es la derivada de la
Energía**

$$P = \frac{\Delta U}{\Delta t} = \frac{\Delta Q}{\Delta t}(V_1 - V_2)$$



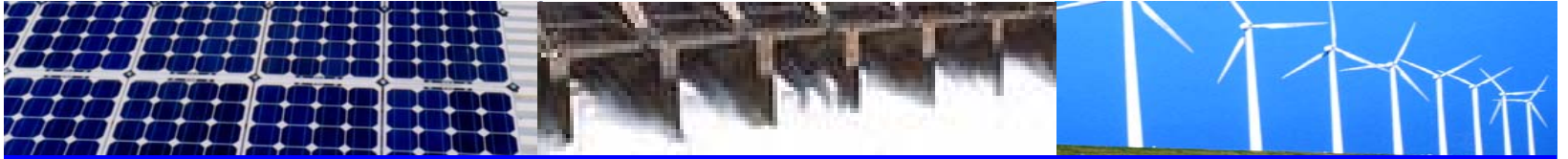
Efecto Joule



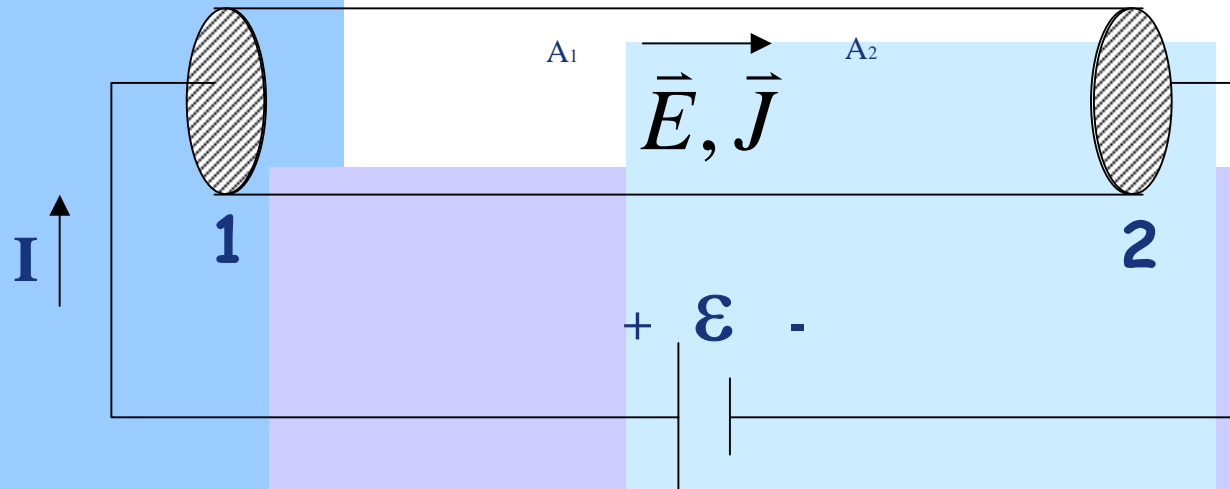
- Calor disipado
- Fem proporciona energía

Potencia es diferencia de potencial por corriente

$$\Rightarrow P = I\Delta V$$



Efecto Joule

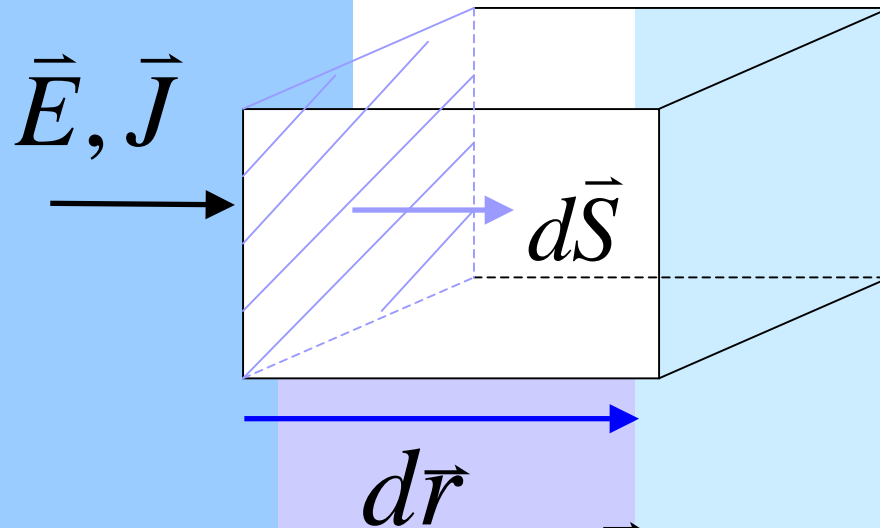


- Calor disipado
- Fem proporciona energía

$$\Delta V = RI \Rightarrow P = I \cdot R \cdot I = RI^2 \quad \text{ó} \quad P = \frac{\Delta V^2}{R}$$



Efecto Joule



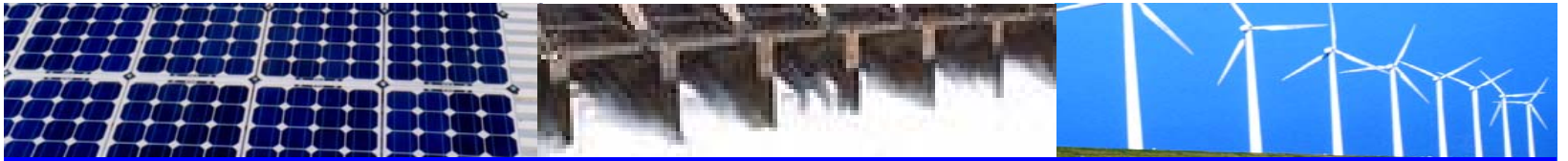
$$dP = I \Delta V$$

$$dP = \underbrace{(\vec{J} \cdot d\vec{S})}_I \cdot \underbrace{(\vec{E} \cdot d\vec{r})}_{\Delta V}$$

$$d\vec{S} \cdot d\vec{r} = dv \Rightarrow dP = \vec{J} \cdot \vec{E} \cdot dv$$

**Potencia disipada
en volumen Ω**

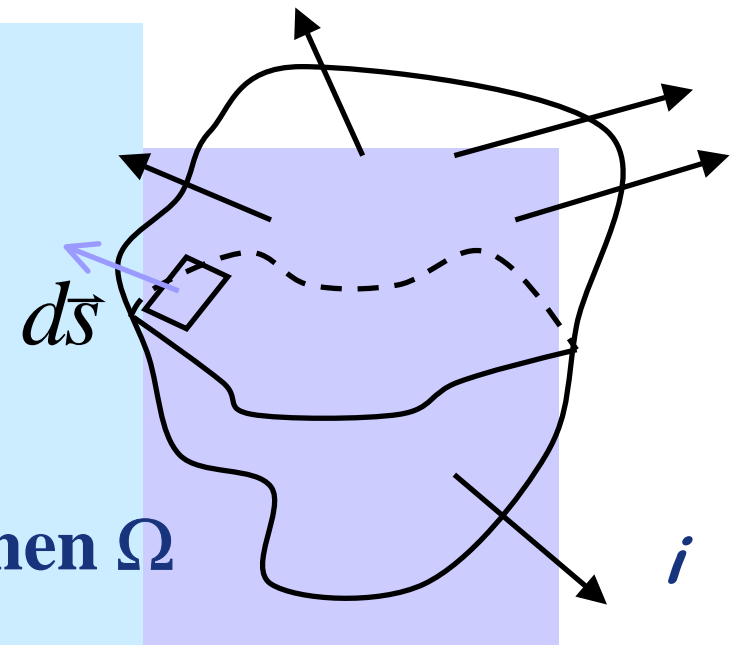
$$\therefore P = \iiint_{\Omega} \vec{J} \cdot \vec{E} dv$$



Ecuación de Continuidad

Corriente saliendo de volumen Ω

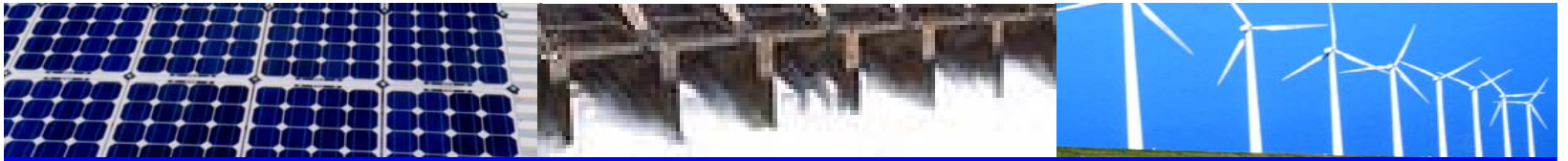
$$I_{salida} = \oiint_{S(\Omega)} \vec{J} \cdot d\vec{S}$$



Q_{in} a la carga contenida en el volumen Ω

$$I_{salida} = - \frac{dQ_{in}}{dt}$$

Corriente que sale corresponde a la variación de carga encerrada en el volumen

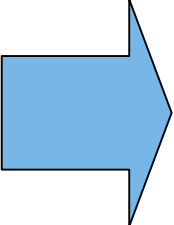


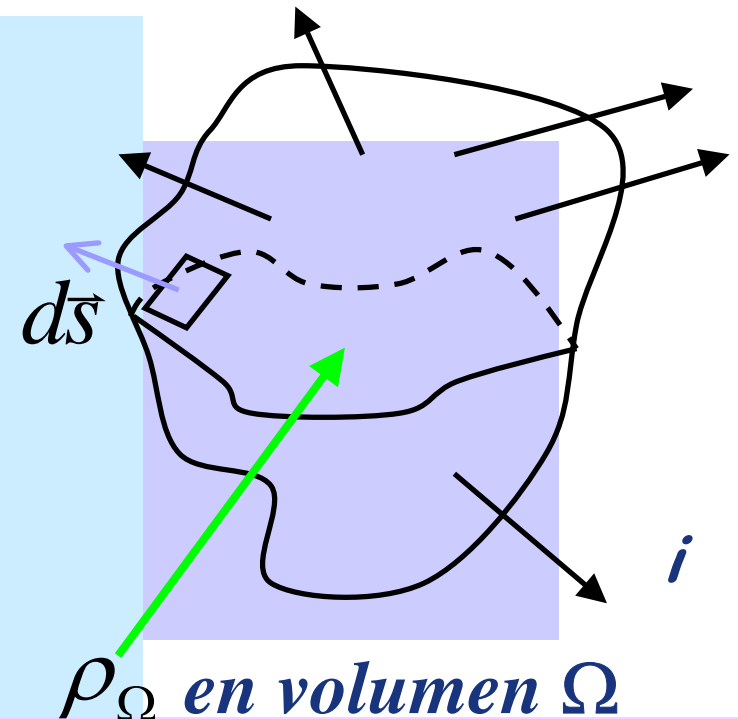
Ecuación de Continuidad

$$Q_{in} = \iiint_{\Omega} \rho_{\Omega}(\vec{r}) dV \quad (5.30)$$

$$I_{salida} = -\frac{\partial}{\partial t} \iiint_{\Omega} \rho_{\Omega}(\vec{r}) dV \quad (5.31)$$

volumen Ω es fijo (no depende de t)


$$I_{salida} = -\iiint_{\Omega} \left[\frac{\partial}{\partial t} \rho_{\Omega}(\vec{r}) \right] dV$$

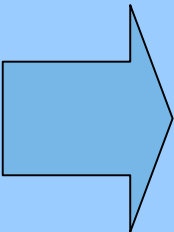


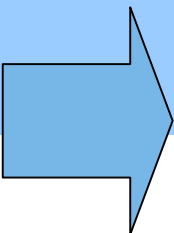


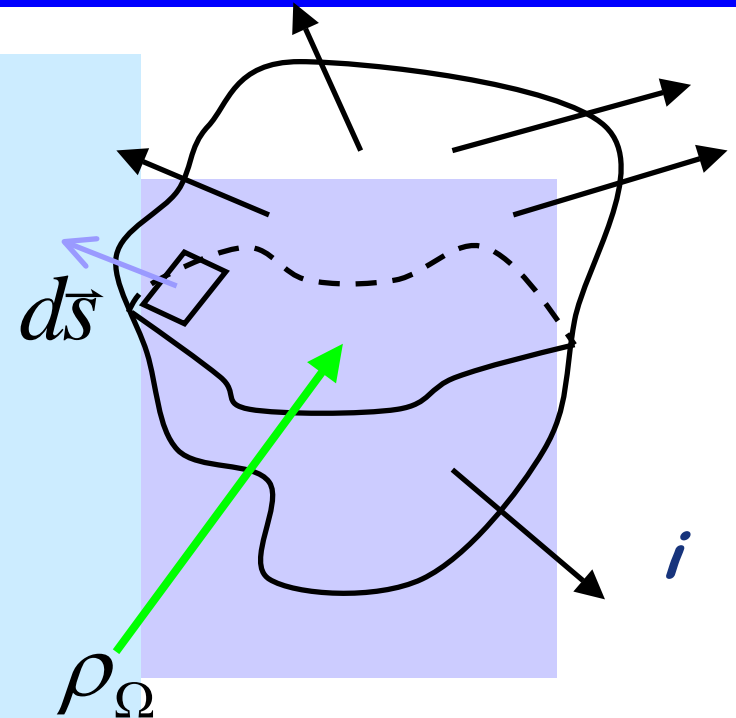
Ecuación de Continuidad

teníamos

$$-\frac{dQ_{in}}{dt} = I_{salida}$$

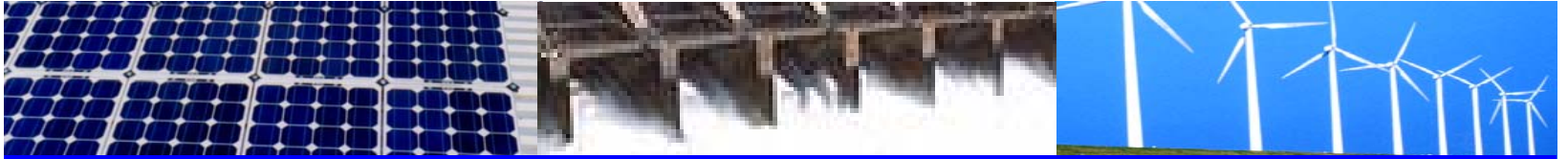

$$-\iiint_{\Omega} \left[\frac{\partial}{\partial t} \rho_{\Omega}(\vec{r}) \right] dV = \oiint_{S(\Omega)} \vec{J} \cdot d\vec{S}$$


$$-\iiint_{\Omega} \left[\frac{\partial}{\partial t} \rho_{\Omega}(\vec{r}) \right] dV = \iiint_{\Omega} \nabla \cdot \vec{J} dV$$



$$\Rightarrow \nabla \cdot \vec{J} + \frac{\partial}{\partial t} \rho(\vec{r}) = 0$$

Ecuación de continuidad



Ecuación de Continuidad

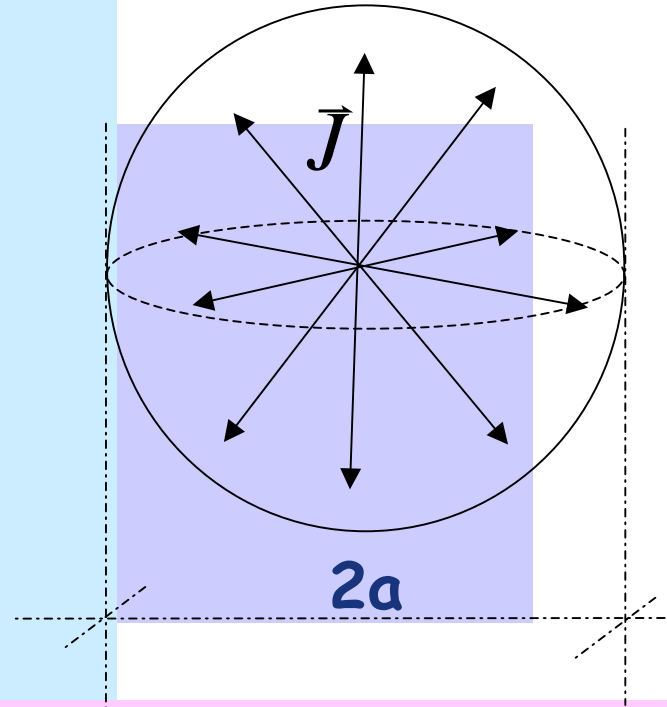
EJEMPLO

Calcular corriente total saliendo
del círculo de radio a si $\underline{j} = J_0 \underline{r}$

$$I = \oiint_S J_0 a \hat{r} \cdot d\vec{s}$$

$$I = 4\pi a^2 J_0 a$$

$$I = 4\pi a^3 J_0$$





Ecuación de Continuidad

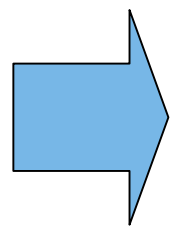
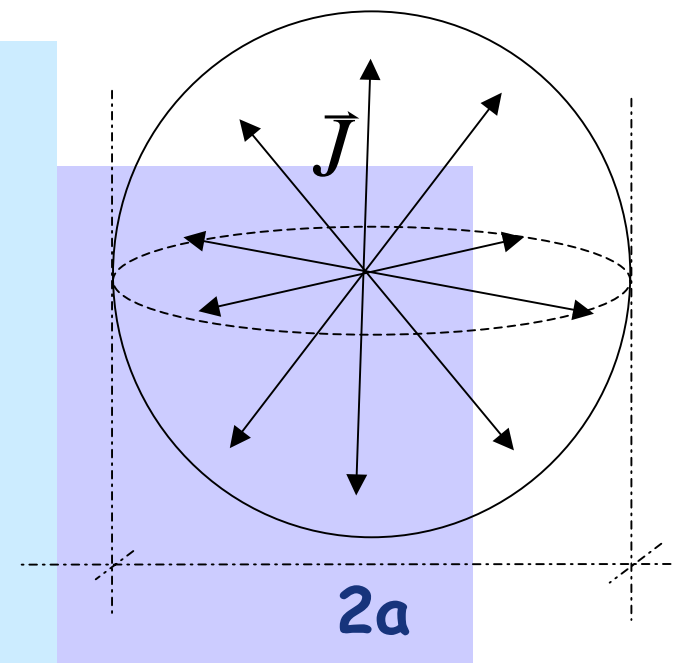
EJEMPLO: Si se tiene una densidad
ed corriente

$$\vec{J} = J_0 \vec{r}$$

Calcular densidad de carga en volumen

$$\nabla \cdot \vec{J} + \frac{\partial}{\partial t} \rho(\vec{r}) = 0$$

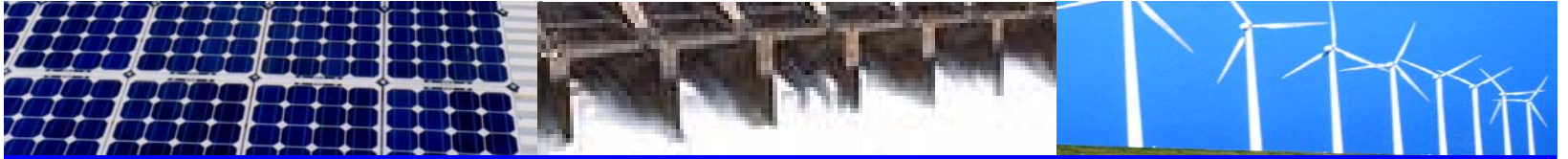
$$\nabla \bullet \vec{J} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 J_0 r \hat{r}) = \frac{1}{r^2} J_0 3r^2 = 3J_0$$



$$\frac{\partial}{\partial t} \rho(\vec{r}) = -3J_0$$



$$\rho(\vec{r}, t) = -3J_0 t + \rho_0(\vec{r})$$

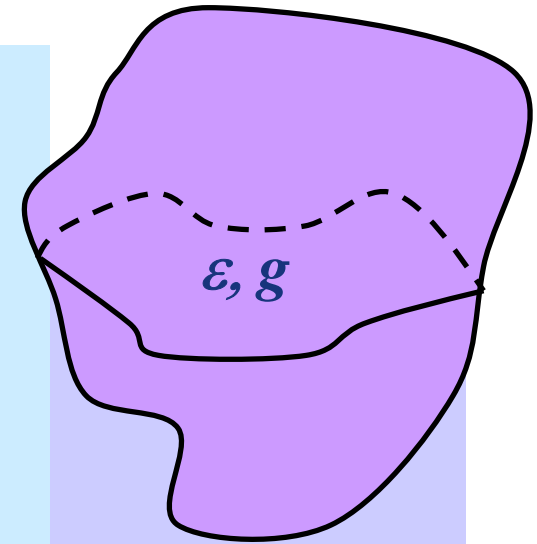


Ecuación de Continuidad en Medios materiales

$$\nabla \cdot \vec{J} + \frac{\partial}{\partial t} \rho(\vec{r}) = 0$$

$$\vec{J} = g \vec{E} \Rightarrow \nabla \cdot \vec{J} = g \nabla \cdot \vec{E}$$

$$\vec{D} = \varepsilon \vec{E} \Rightarrow \nabla \cdot \vec{J} = \frac{g}{\varepsilon} \nabla \cdot \vec{D} = \frac{g}{\varepsilon} \rho(t)$$



volumen Ω

$$\frac{g}{\varepsilon} \rho(t) + \frac{\partial \rho(t)}{\partial t} = 0 \Rightarrow \rho(t) = \rho_0 e^{-t/T_R}$$

$$T_R = \varepsilon / g$$

Constante de relajación



Ecuación de Continuidad en Medios materiales

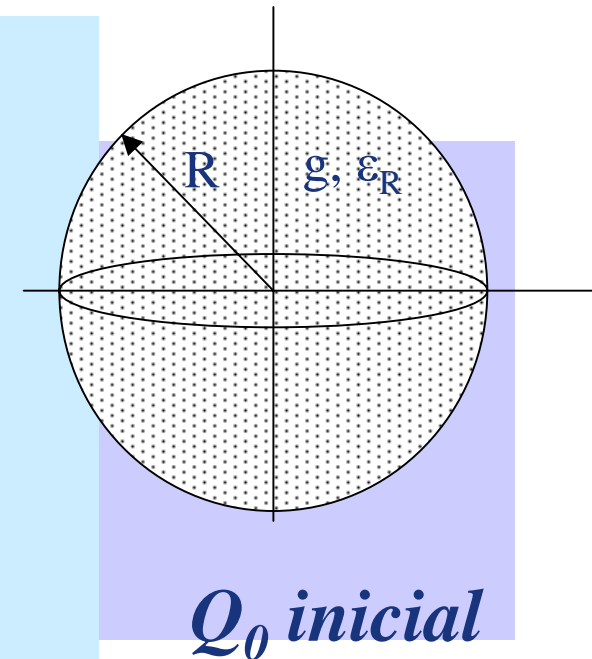
EJEMPLO

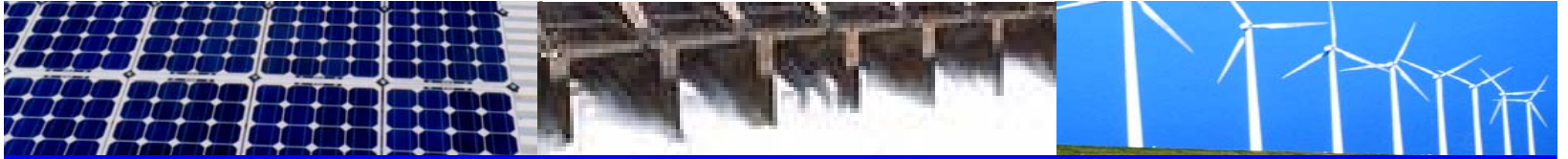
$$\frac{g}{\varepsilon} \rho(t) + \frac{\partial \rho(t)}{\partial t} = 0$$

$$\Rightarrow \iiint_V \left(\frac{g}{\varepsilon} \rho(t) + \frac{\partial \rho(t)}{\partial t} \right) dV = 0$$

$$\Rightarrow \underbrace{\frac{g}{\varepsilon} \iiint_V \rho(t) dV}_{Q(t)} + \frac{\partial}{\partial t} \underbrace{\iiint_V \rho(t) dV}_{Q(t)} = 0$$

$$\Rightarrow \frac{g}{\varepsilon} Q(t) + \frac{\partial Q(t)}{\partial t} = 0 \Rightarrow Q(t) = Q_0 e^{-t/T_R}$$



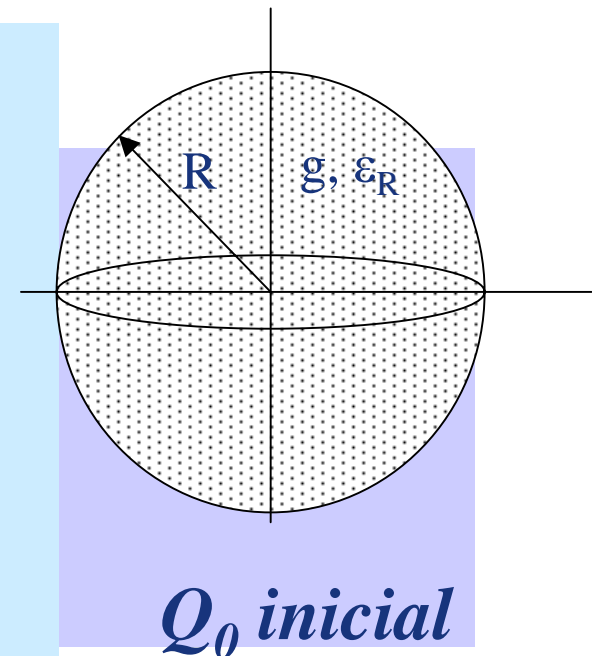


Ecuación de Continuidad en Medios materiales

EJEMPLO

$$Q(t) = Q_0 e^{-t/T_R}$$

| | | |
|-------------------------|------------------------------------|-----------------------------|
| | cobre | Cuarzo fusionado |
| T_R | $1.53 \times 10^{-19} \text{ seg}$ | 51.2 días |





Condiciones de Borde para \vec{J}

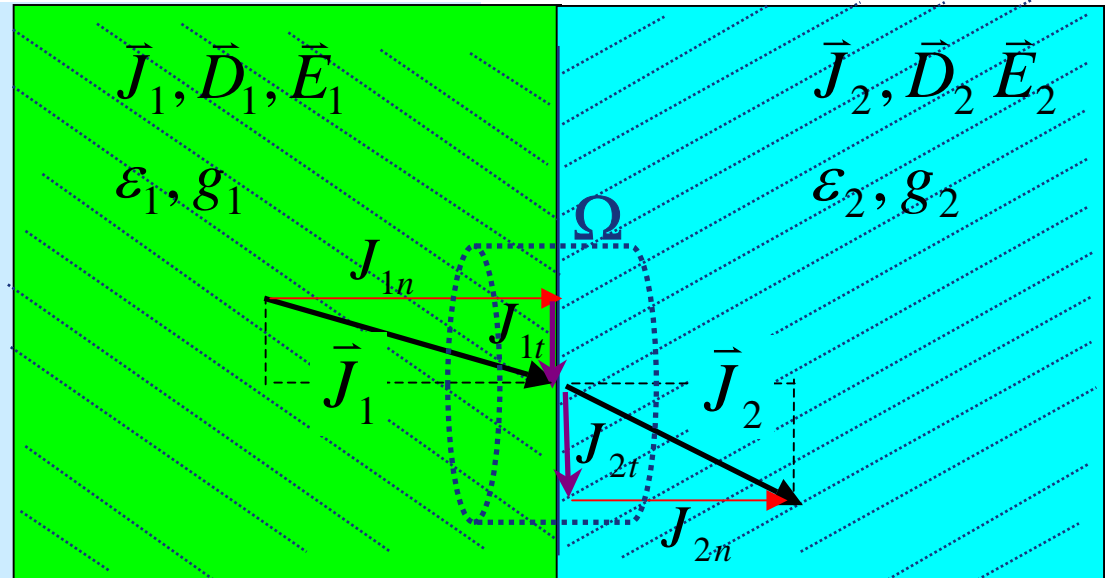
$$\nabla \times \vec{E} = 0$$

$$E_{1t} = E_{2t} \Rightarrow \frac{J_{1t}}{g_1} = \frac{J_{2t}}{g_2}$$

$$\nabla \cdot \vec{D} = \rho$$

$$D_{1N} - D_{2N} = \sigma_{\text{libre}}$$

$$\epsilon_1 E_{1n} - \epsilon_2 E_{2n} = \sigma_l \Rightarrow \epsilon_1 \frac{J_{1n}}{g_1} - \epsilon_2 \frac{J_{2n}}{g_2} = \sigma_l$$

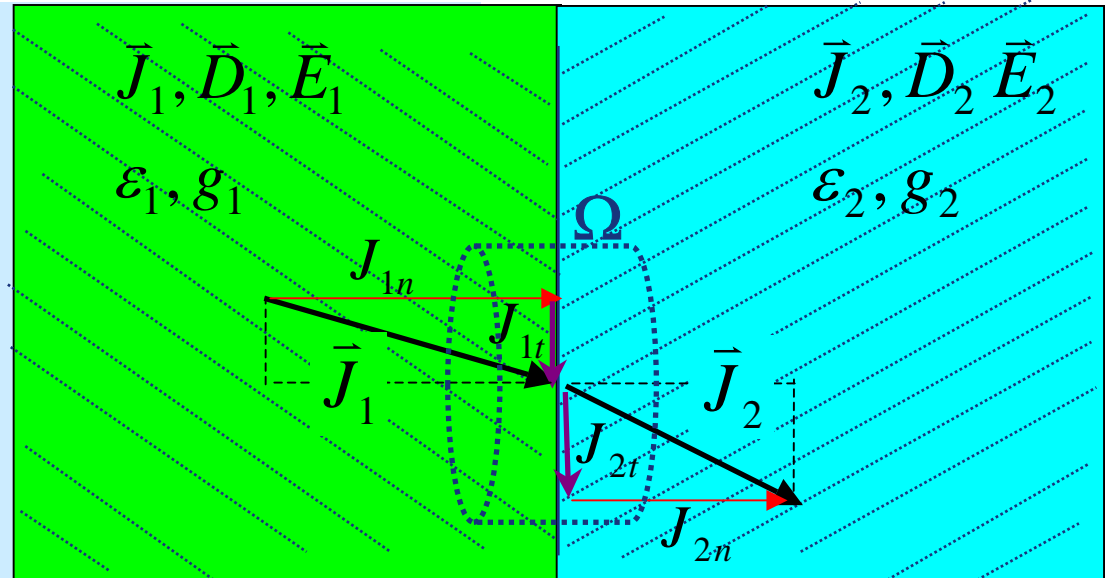




Condiciones de Borde para \vec{J}

$$\frac{J_{1t}}{g_1} = \frac{J_{2t}}{g_2}$$

$$\epsilon_1 \frac{J_{1n}}{g_1} - \epsilon_2 \frac{J_{2n}}{g_2} = \sigma_l$$



I. Situación Estacionaria $\frac{\partial \rho(t)}{\partial t} = 0$

$$\nabla \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0 \Rightarrow \nabla \cdot \vec{J} = 0$$

$$\oiint_{S(\Omega)} \vec{D} \cdot d\vec{s} = 0 \Rightarrow J_{1n} = J_{2n}$$

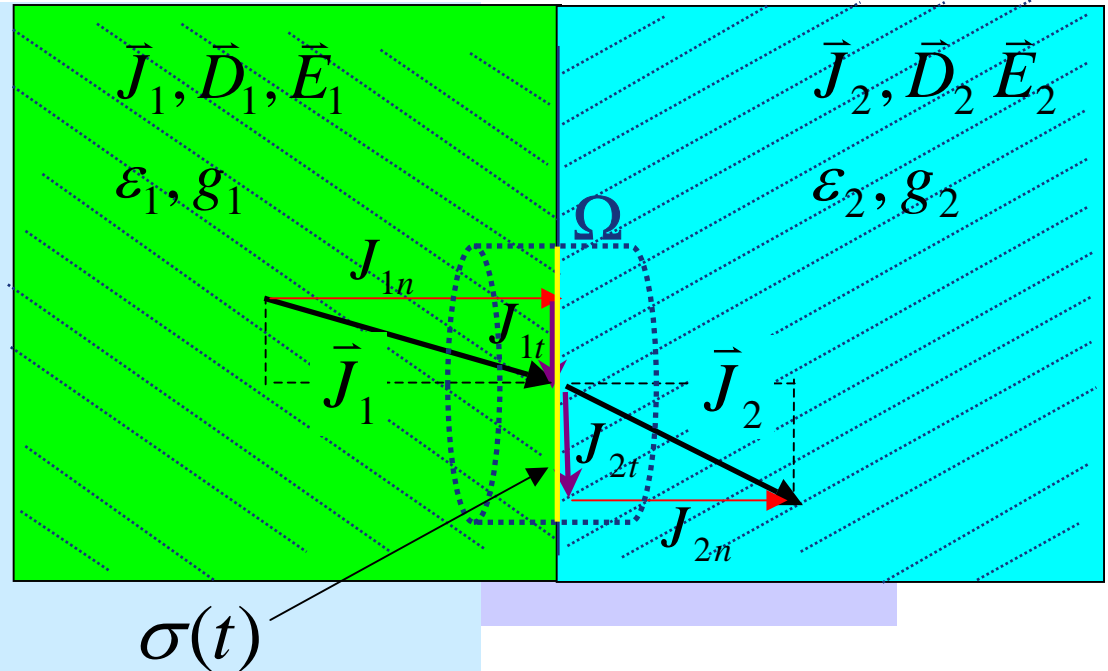


Condiciones de Borde para \vec{J}

II. Situación transitoria

$$\frac{\partial \rho(t)}{\partial t} \neq 0$$

$$\oiint_{S(\Omega)} \vec{J} \cdot d\vec{S} = J_{2n} \Delta S - J_{1n} \Delta S$$



Haciendo tender la altura del cilindro a cero $\frac{\partial Q}{\partial t}$ se acumula sólo en la superficie que limita los medios

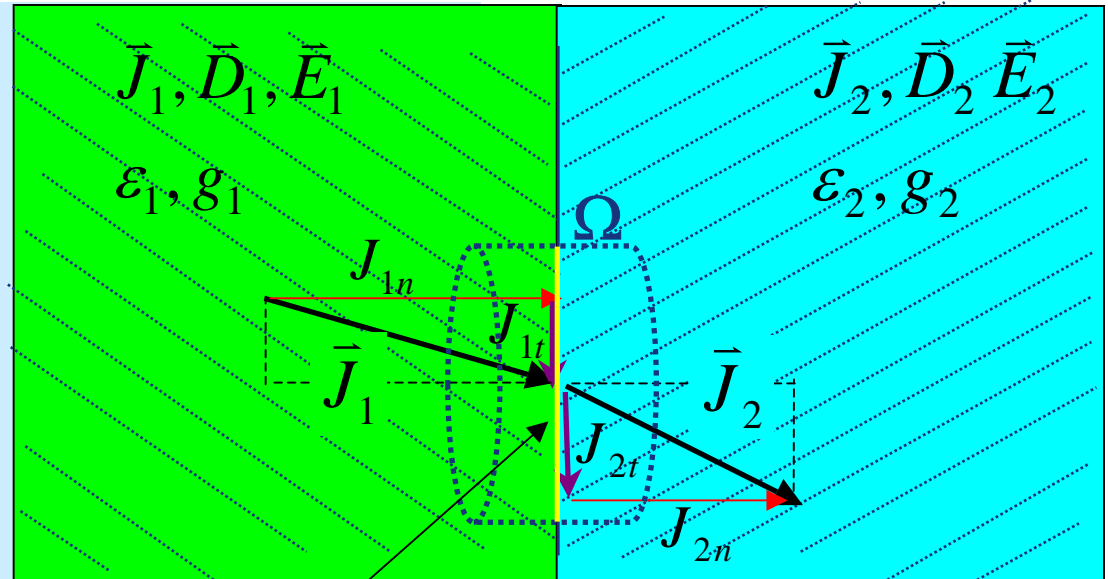


Condiciones de Borde para \vec{J}

$$\Rightarrow \frac{\partial Q}{\partial t} = \frac{\partial}{\partial t} (\sigma \cdot \Delta S)$$

$$\Rightarrow J_{2n} \Delta S - J_{1n} \Delta S + \frac{\partial \sigma}{\partial t} \Delta S = 0$$

$$\Rightarrow J_{2n} - J_{1n} + \frac{\partial \sigma}{\partial t} = 0$$



$\sigma(t)$

$$\therefore J_{2n} - J_{1n} = \frac{\partial \sigma(t)}{\partial t}$$