

# Discrete Photonics in Waveguide Arrays

Mario I. Molina

Departamento de Física, MSI-Nucleus on Advanced Optics,  
and Center for Optics and Photonics (CEFOP), Facultad de Ciencias,  
Universidad de Chile, Santiago, Chile



<http://fisica.ciencias.uchile.cl/nonopt/NLOG.html>  
<http://www.cefop.cl/>





## Why study physics of discrete systems?

Testbed to test general phenomenology

Richer physics than continuous counterpart

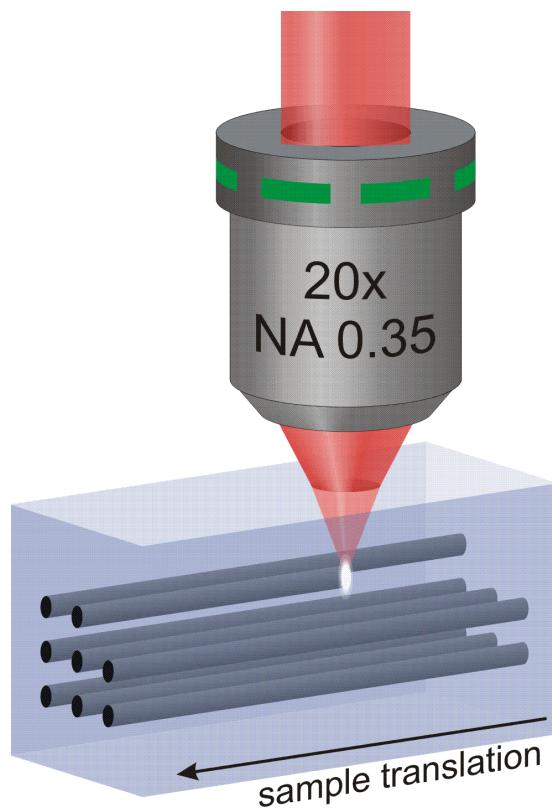
Greater potential for applications



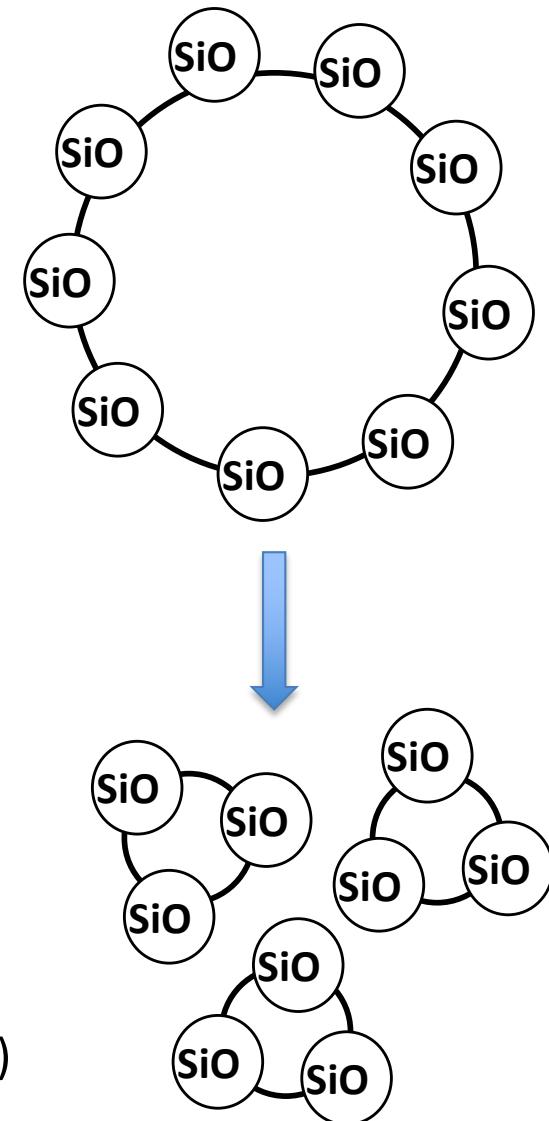
# Waveguides in fused silica



A. Szameit et al, Opt. Express 13,10552 (2005).



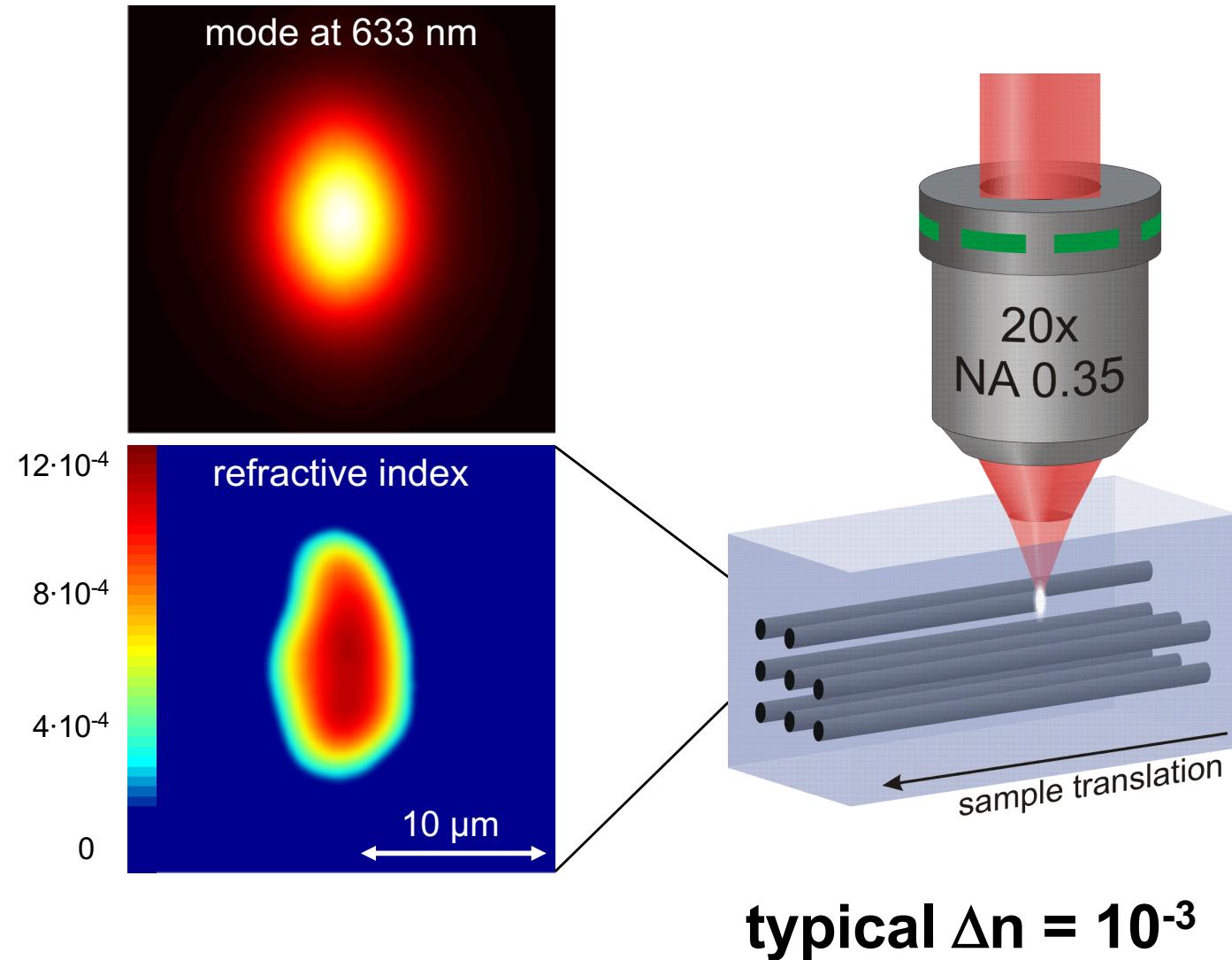
9-ring structure



3-ring structures  
(densification,  
refractive index increase)



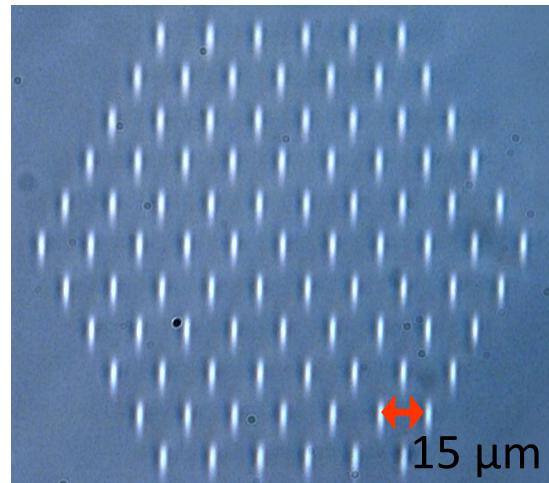
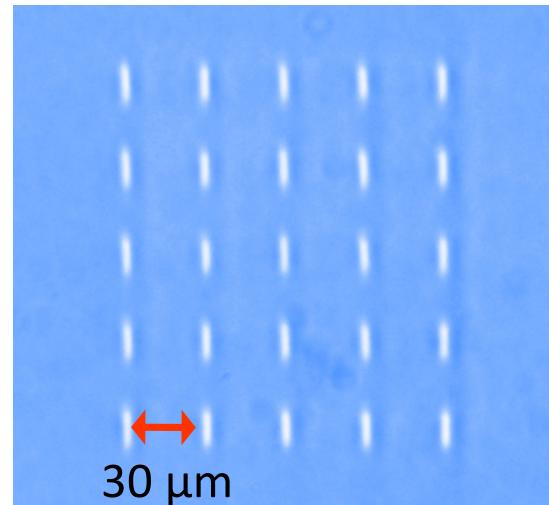
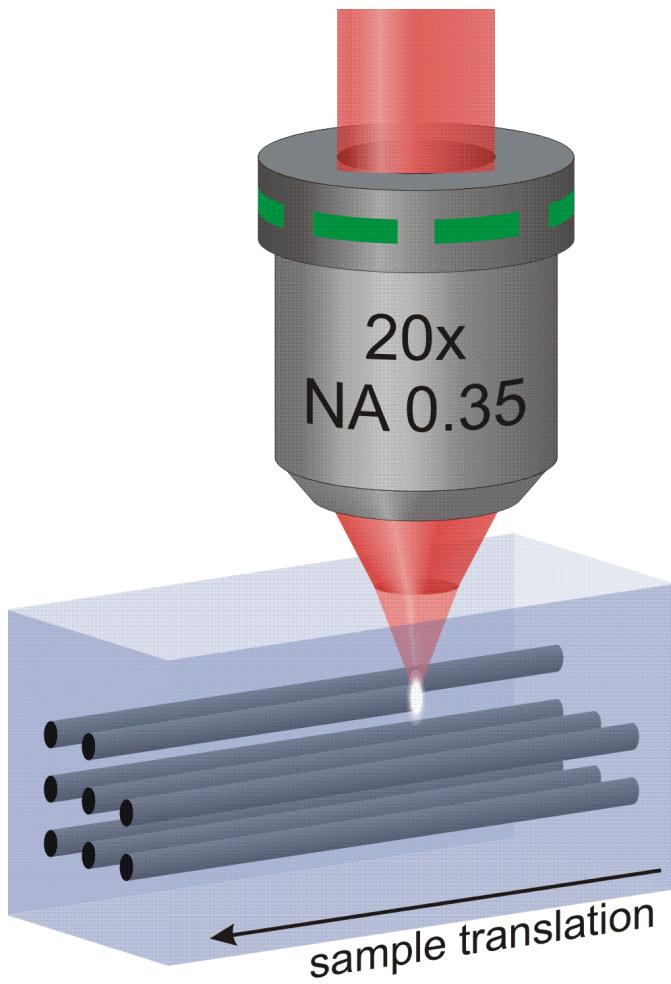
# Waveguides in fused silica





# Waveguides in fused silica

OFOP



AS et al., Opt. Lett. **33**, 663 (2008).

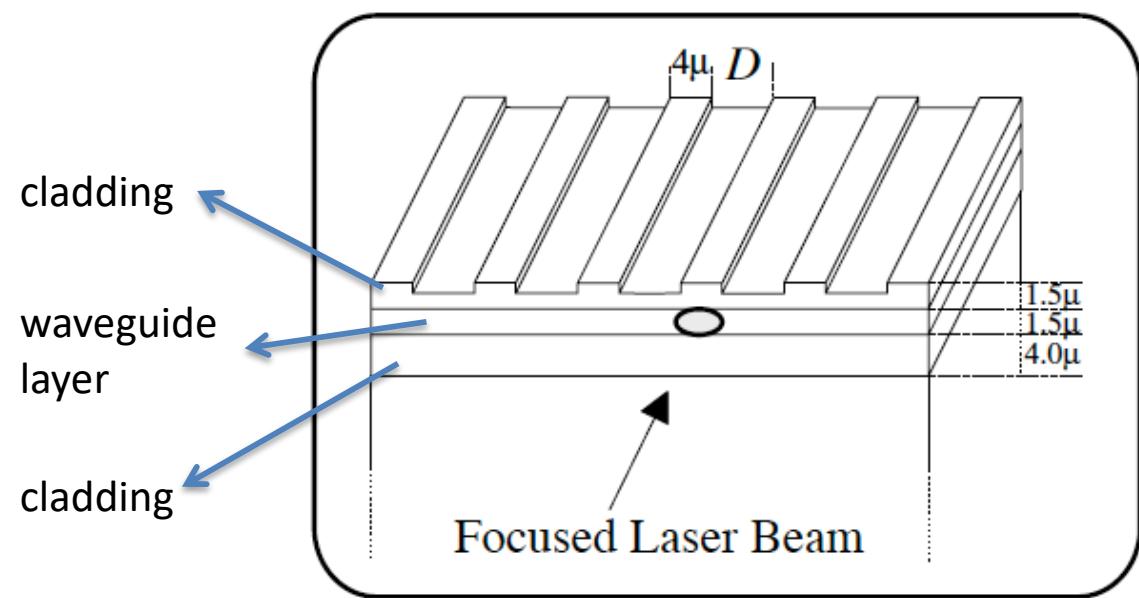
AS et al., Appl. Phys. B **82**, 507 (2006).



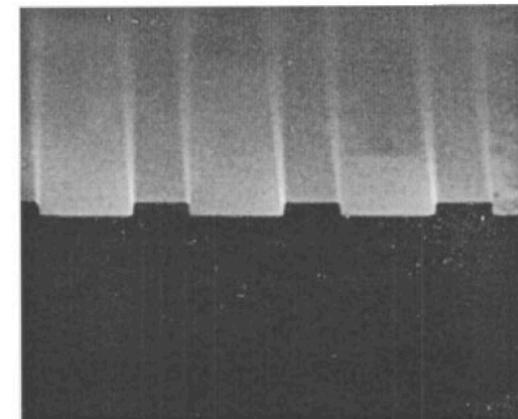
# Semiconductor Waveguides



P. Millar, J.S. Aitchson, J.U. Kang, G.I. Stegeman, J. Opt. Soc. Am. B 14, 3224 (1997).



Substrate: Ga As  
Cladding:  $\text{Al}_{0.24}\text{Ga}_{0.76}\text{As}$   
Waveguide layer:  $\text{Al}_{0.18}\text{Ga}_{0.82}\text{As}$

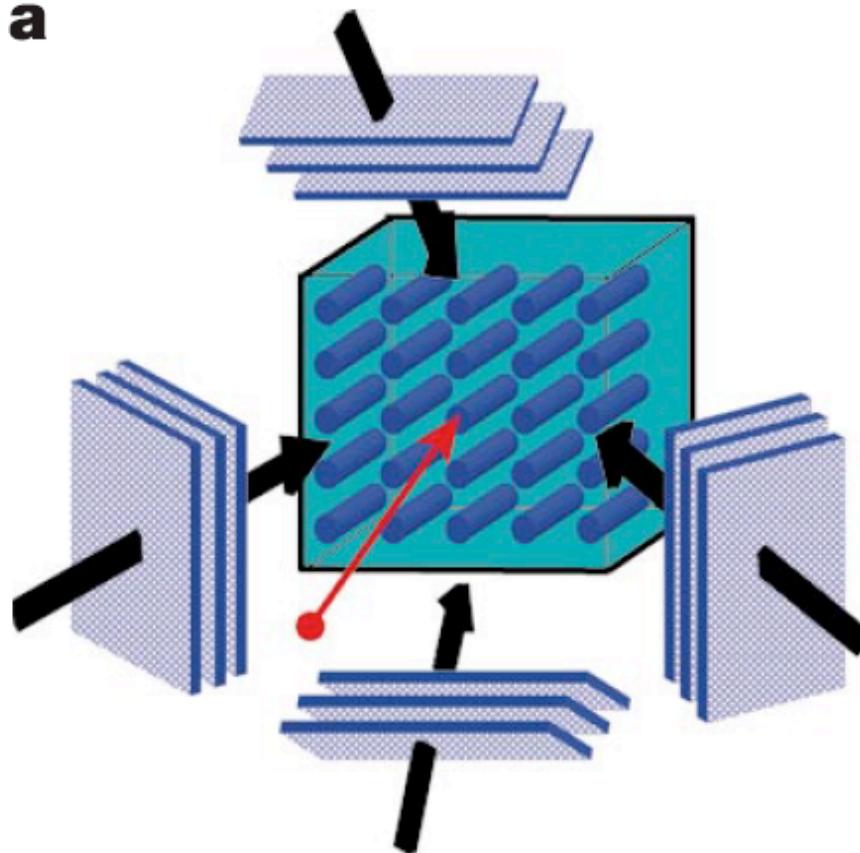




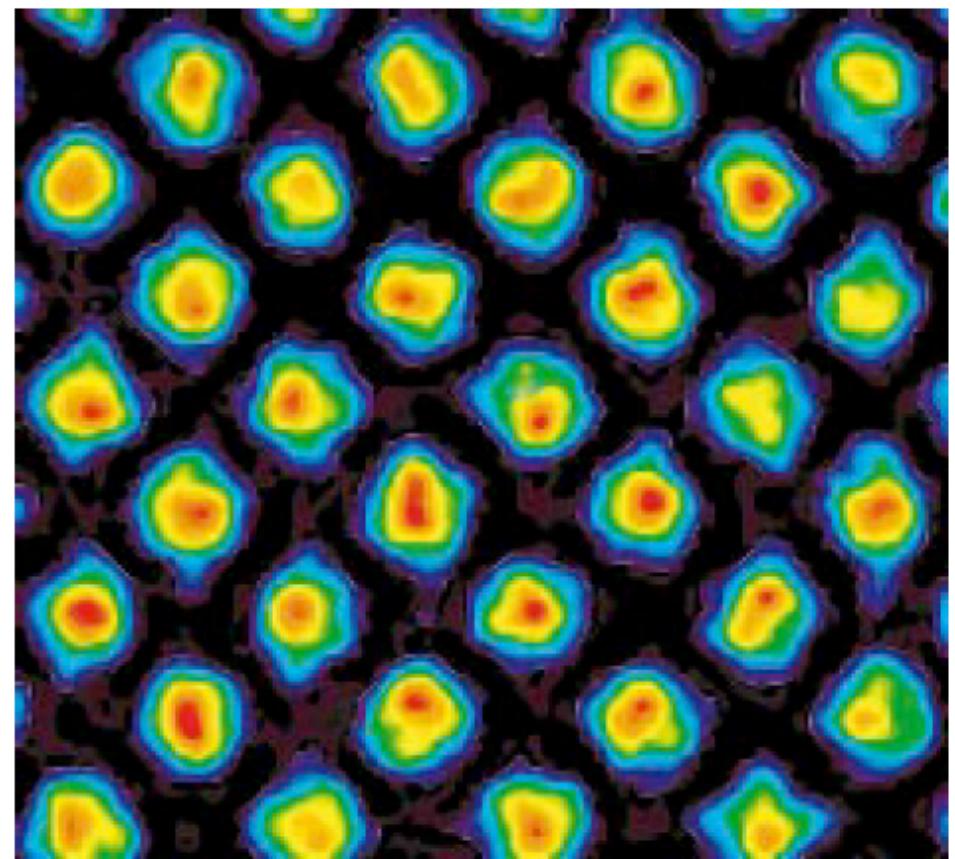
# Photorefractive Waveguides

CEFRON

a



b



Light → releases electrons → drift → local E fields → electro-optic effect → distribution of refractive indices



# Coupled-modes theory



Maxwell:

$$\vec{\nabla} \cdot \vec{D} = 0$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{H} = \frac{\partial}{\partial t} \vec{D}$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial}{\partial t} \vec{B}$$

$$\vec{D} = \vec{E} + \vec{P}$$

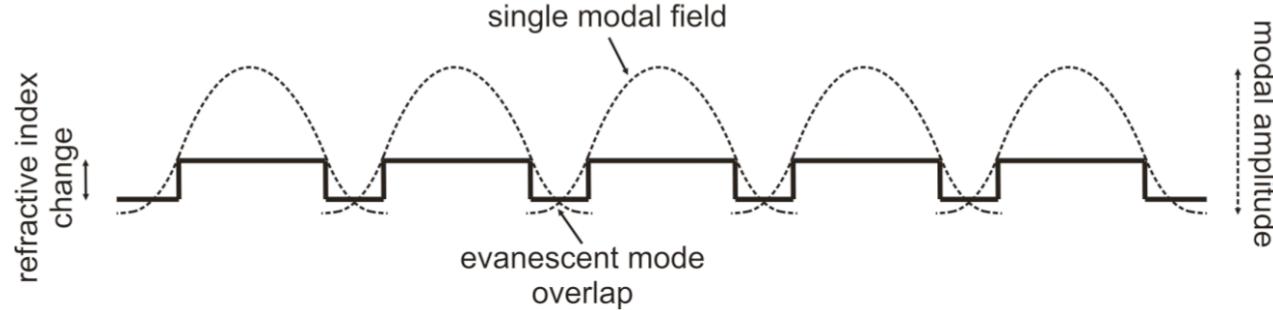
$$\vec{H} = \vec{B} + \vec{M}$$

$$\nabla^2 \vec{E} - \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = \frac{1}{c^2} \frac{\partial^2 \vec{P}}{\partial t^2}$$

$$\vec{P} = \chi^{(1)} \vec{E} + \chi^{(3)} |\vec{E}|^2 \vec{E}$$



# Teoría de modos acoplados



$$E(x, z) = \sum_{n=-\infty}^{\infty} C_n(z) \phi(x - x_n)$$

$$\left| \frac{d^2 C_n}{dz^2} \right| \ll k_0 \left| \frac{d C_n}{dz} \right| \quad n = n_0 + n_2 |E|^2 \quad \text{Kerr}$$

$$i \frac{d C_n}{dz} + V(C_{n+1} + C_{n-1}) + \gamma |C_n|^2 C_n = 0$$

Discrete nonlinear Schrodinger (DNLS) equation



## Coupled-modes theory



$$P = \sum_n |C_n|^2$$

$$H = \sum_n \{ V(C_n C_{n+1}^* + C_n^* C_{n+1}) + (\gamma/2) |C_n|^4 \}$$

} Conserved quantities

$$q_n = C_n; \quad p_n = i C_n^*$$

$$(d/dt)q_n = \partial H / \partial p_n \quad (d/dt)p_n = -\partial H / \partial q_n$$

Hamiltonian system

$$C_n = u_n \exp(i\beta z) \quad \text{Stationary mode}$$

$$-\beta u_n + (u_{n+1} + u_{n-1}) + \chi |u_n|^2 u_n = 0$$

Nonlinear eigenvalue equation



## Coupled-modes theory



Finding the localized nonlinear mode

$$-EC_n + V(C_{n+1} + C_{n-1}) + \chi|C_n|^2C_n = 0$$

$$\lambda \equiv E/V, \quad \phi_n \equiv \sqrt{\chi/V}C_n$$

$$-\lambda\phi_n + (\phi_{n+1} + \phi_{n-1}) + |\phi_n|^2\phi_n = 0$$

$\vec{F}(\vec{\phi}) = 0$  use Newton-Raphson

Need good seed (anticontinuous limit)

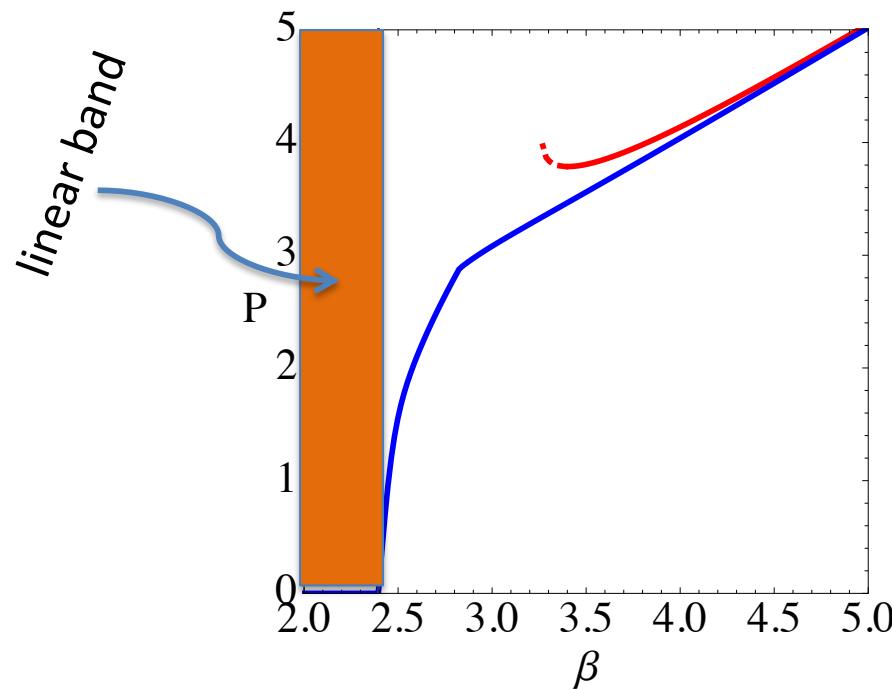
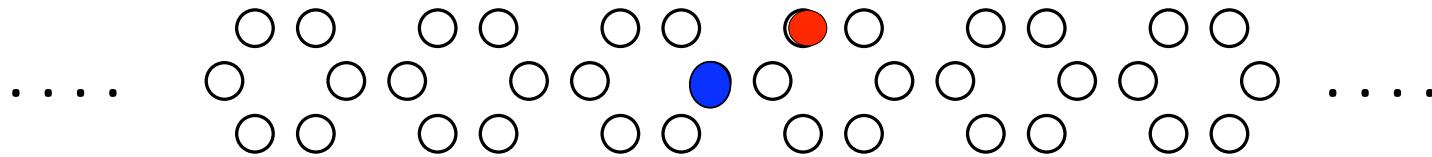
Find many solution families (characterized by power  
vs prop.const. curve)



# Coupled-modes theory



Example: Graphene ribbon



$$P = \sum_n |\phi_n|^2$$

Conserved quantity



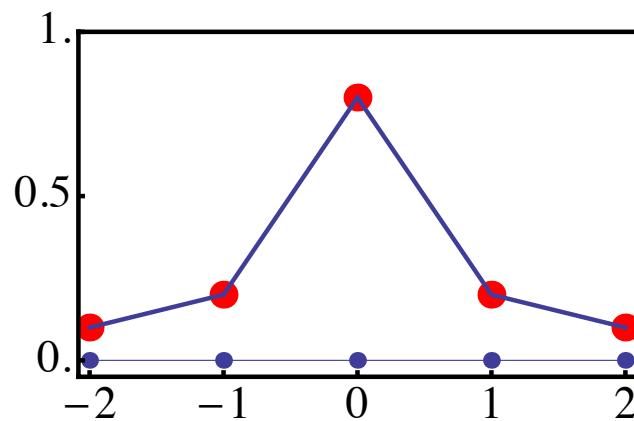
# Coupled-modes theory



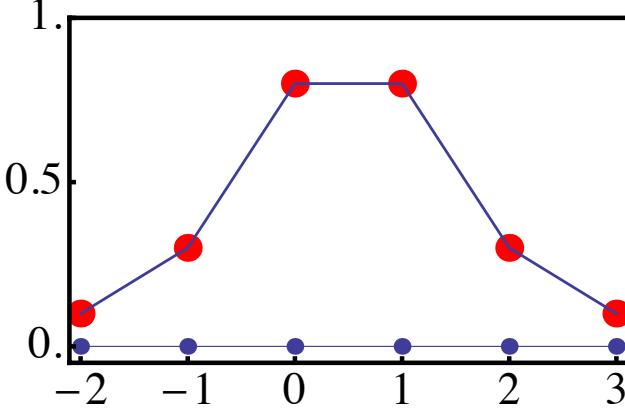
Pure 1D CASE



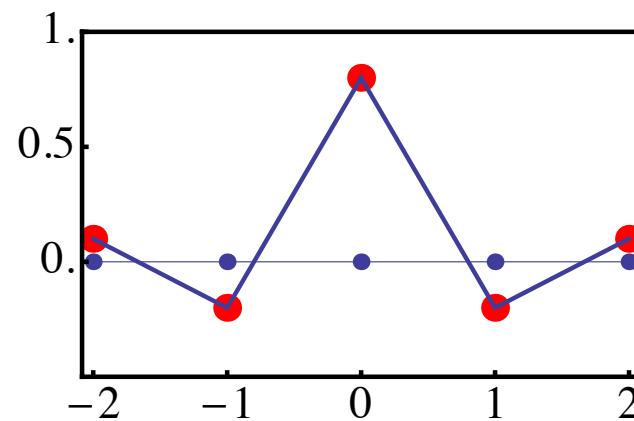
odd  
unstaggered



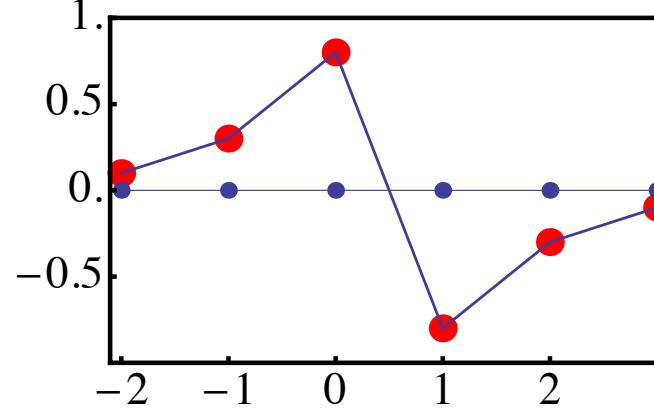
even



odd  
staggered



twisted





## Coupled-modes theory



### Linear stability

$$C_n(z) = \phi_n e^{-i\lambda z} \quad \text{sol. of DNLS}$$

$$C_n(z) \rightarrow (\phi_n + \delta\phi_n) e^{-i\lambda z}, \quad |\delta\phi_n/\phi_n| \ll 1$$

$$\implies \text{Equation for } \delta\phi_n = \delta u_n + i\delta v_n$$

$$\text{define } \delta \vec{u} = (\delta u_1, \delta u_2, \dots, \delta u_N), \quad \delta \vec{v} = (\delta v_1, \delta v_2, \dots, \delta v_N)$$

$$\mathcal{A}_{nm} = \delta_{n,m+1} + \delta_{n,m-1} + (\lambda + \phi_n^2) \delta_{n,m}$$

$$\mathcal{B}_{nm} = \delta_{n,m+1} + \delta_{n,m-1} + (\lambda + 3\phi_n^2) \delta_{n,m}$$

$$\boxed{\ddot{\delta \vec{U}} + \mathcal{B}\mathcal{A} \ \delta \vec{U} = 0 \quad \text{and} \quad \ddot{\delta \vec{V}} + \mathcal{A}\mathcal{B} \ \delta \vec{V} = 0}$$



## Coupled-modes theory



$\{m\}$ =eigenvalues of  $\mathcal{AB}$  = eigenvalues of  $\mathcal{BA}$

instability gain

$$G^* = \text{Max} \left\{ \sqrt{(1/2)(-\text{Re}[m] + \sqrt{\text{Re}[m]^2 + \text{Im}[m]^2})} \right\}$$

$G^* = 0$  stable

$G^* > 0$  unstable

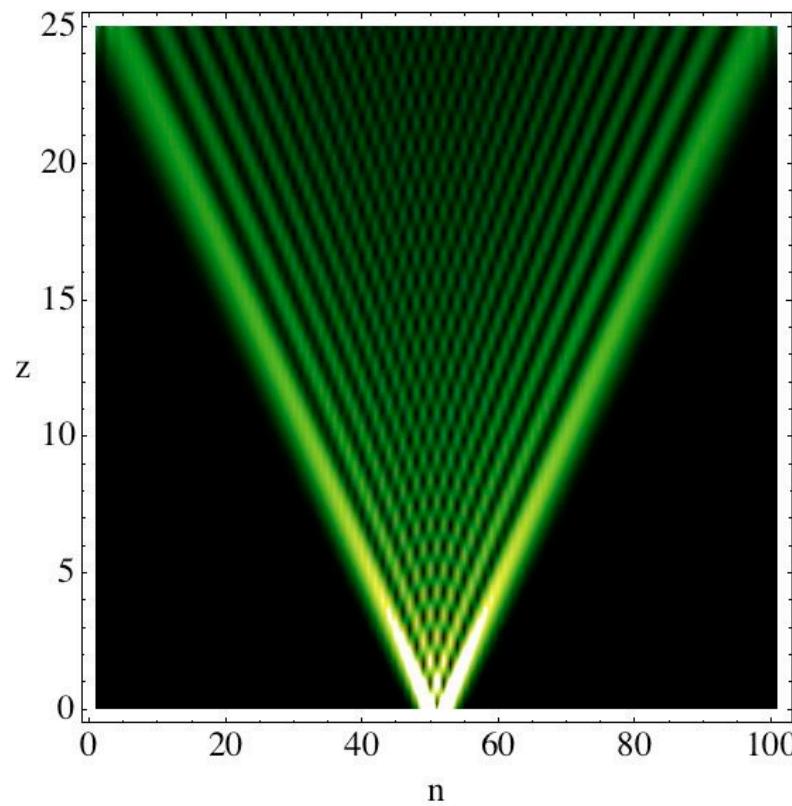


# Coupled-modes theory

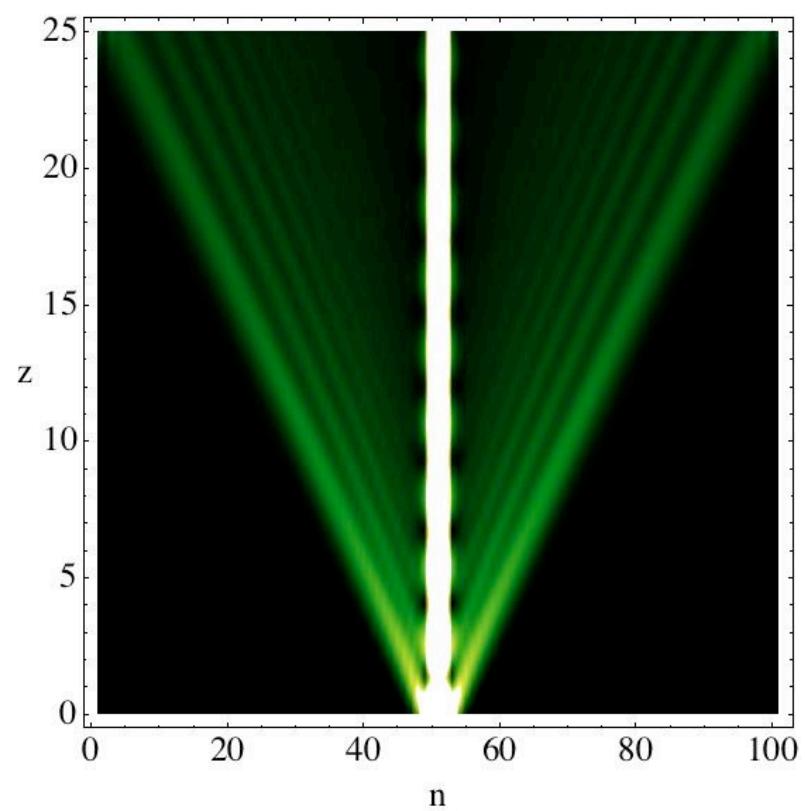


## Numerical propagation

Discrete diffraction

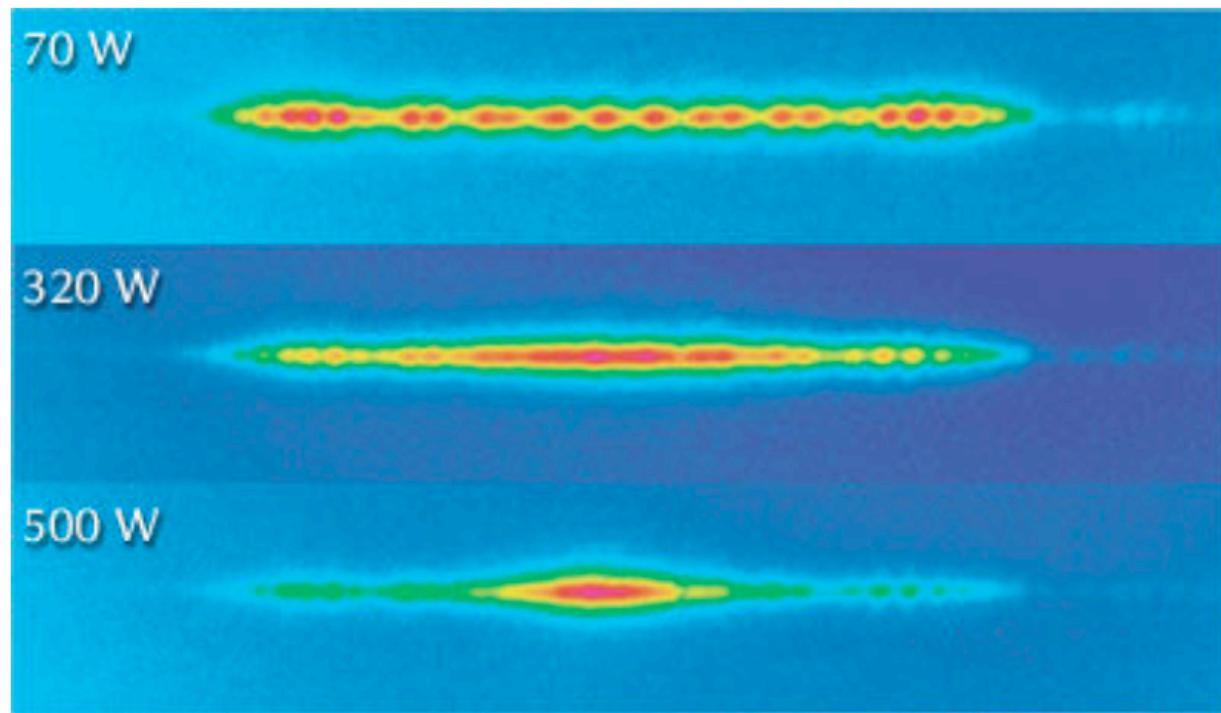


Discrete soliton formation





# First experimental observation of discrete soliton



Difraccion discreta

Soliton discreto

H. Eisenberg et al, PRL 81, 3383 (1998).