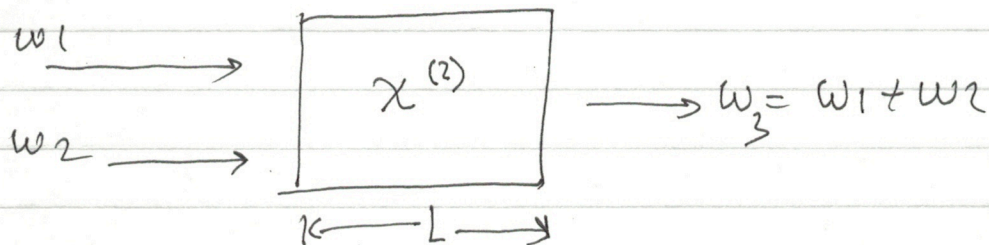


## Ec. de onda acopladas para la generación de suma de frecuencias



En ausencia de no linealidad (bajas intensidades)

$$\vec{E}_3(z,t) = A_3 e^{i(k_3 z - \omega_3 t)} + c.c. \quad (1)$$

$$k_3 = \frac{n_3 \omega_3}{c}, \quad n_3 = \sqrt{\epsilon^{(1)}(\omega_3)} \quad (2)$$

Cuando se aumenta la intensidad,  $A_3 \rightarrow A_3(z)$  función de variable

$$P_{NL}(z,t) = P_3 e^{-i\omega_3 t} + c.c. \quad (3)$$

$$\text{donde } P_3 = 4d E_1 E_2 \quad (= 2\chi^{(2)} E_1 E_2) \quad (4)$$

$$E_i(z,t) = E_i e^{-i\omega_i t} + c.c. \quad (i=1,2) \quad (5)$$

con  $E_i = A_i e^{i k_i z}$

$$\Rightarrow P_3 = 4d A_1 A_2 e^{i(k_1 + k_2)z} \quad (6)$$

se sust. todo esto en la ec. de onda:



se cancelan  
mutuo'!

$$\left( \frac{\partial^2 A_3}{\partial z^2} + 2ik_3 \frac{\partial A_3}{\partial z} - \cancel{k_3^2 A_3} + \frac{\epsilon^{(1)}(\omega_3) \omega_3^2 A_3}{c^2} \right) e^{i(k_3 z - \omega_3 t)} + c.c.$$

$$= -\frac{16\pi d\omega_3^2}{c^2} A_1 A_2 e^{i(k_1 + k_2 - k_3)z} + c.c.$$

$$\rightarrow \frac{\partial^2 A_3}{\partial z^2} + 2ik_3 \frac{\partial A_3}{\partial z} = -\frac{16\pi d\omega_3^2}{c^2} A_1 A_2 e^{i(k_1 + k_2 - k_3)z}$$

Ahora suponemos  $\left| \frac{\partial^2 A_3}{\partial z^2} \right| \ll \left| k_3 \frac{\partial A_3}{\partial z} \right|$   
(Aprox. de amplitud lenta' variable)

$$\frac{dA_3}{dz} = \frac{8\pi i d\omega_3^2}{k_3 c^2} A_1 A_2 e^{i(\Delta K)z}$$

(7)

donde  $\Delta K \equiv k_1 + k_2 - k_3$  (parámetro de mismatch) (8)

Del lado del medio no-lineal, también venimos  $A_1$  y  $A_2$ .  
Repetiendo el procedimiento anterior para  $\omega_1, \omega_2$ .

$$\frac{dA_1}{dz} = \frac{8\pi i d\omega_1^2}{k_1 c^2} A_3 A_2^* e^{-i\Delta K z}$$

(9)

$$\frac{dA_2}{dz} = \frac{8\pi i d\omega_2^2}{k_2 c^2} A_3 A_1^* e^{-i\Delta K z}$$

(10)

Tarea: obtener el set completo (7), (9), (10) a partir de las ec. de onda y la aprox. de amplitud lenta' variable.

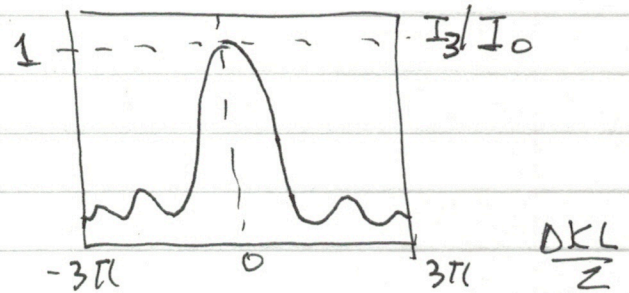


Case I :  $A_1(z)$ ,  $A_2(z) \approx \text{cte.}$

$$A_3(L) = \frac{8\pi i \omega_3^2 d A_1 A_2}{k_3 c^2} \int_0^L e^{i \Delta K z} dz = \frac{8\pi i d \omega_3^2 A_1 A_2}{k_3 c^2} \left( \frac{e^{i \Delta K L} - 1}{i \Delta K} \right)$$

Intensity (a la salida):  $I_3 = \frac{n_3 c}{2\pi} |A_3|^2 = \frac{32\pi^2 d^2 \omega_3^2 |A_1|^2 |A_2|^2 n_3}{k_3^2 c^3} \left| \frac{e^{i \Delta K L} - 1}{\Delta K} \right|^2$

$$I_3 = I_0 \left[ \frac{\sin\left(\frac{\Delta K L}{2}\right)}{\left(\frac{\Delta K L}{2}\right)} \right]^2$$



$\frac{I_3(L)}{I_0}$  es máxima para  $\Delta K = 0$  ("phase-matched")

## Relaciones de Manley-Rowe

$$I_i = \frac{n_i c}{2\pi} A_i A_i^*$$

|                        |
|------------------------|
| $\omega_1 \rightarrow$ |
| $\omega_2 \rightarrow$ |
| $\omega_3 \rightarrow$ |

$$\frac{dI_i}{dz} = \frac{n_i c}{2\pi} \left[ A_i^* \frac{dA_i}{dz} + A_i \frac{dA_i^*}{dz} \right]$$

$$\Rightarrow \frac{dI_1}{dz} = 4d\omega_1 (i A_3 A_1^* A_2^* e^{-i \Delta K z} + \text{c.c.}) = -8d\omega_1 \text{Im} [A_3 A_1^* A_2^* e^{-i \Delta K z}]$$

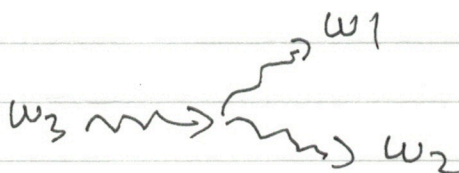
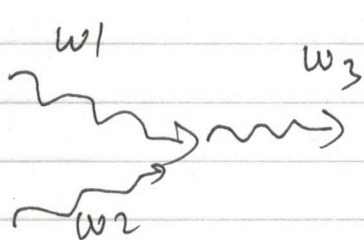
$$\frac{dI_2}{dz} = -8d\omega_2 \text{Im} [A_3 A_1^* A_2^* e^{-i \Delta K z}]$$

$$\Rightarrow \frac{dI_3}{dz} = -8d\omega_3 \operatorname{Im} [A_3^* A_1 A_2 e^{i\Delta K z}] = 8d\omega_3 \operatorname{Im} [A_3 A_1^* A_2^* e^{-i\Delta K z}]$$

$$\frac{d}{dz} \left( \frac{I_1}{\omega_1} \right) = \frac{d}{dz} \left( \frac{I_2}{\omega_2} \right) = - \frac{d}{dz} \left( \frac{I_3}{\omega_3} \right)$$

MC:  $\frac{I_i}{\omega_i} = \frac{\# \text{ fotones}}{v \cdot \text{densidad} \cdot \text{det.}}$

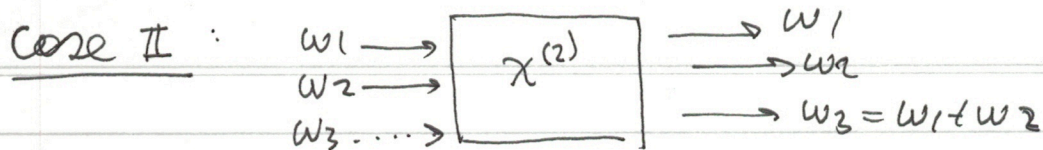
$\Rightarrow$  tener de creación (destrucción) de fotones a frec.  $\omega_1$   
 = tener " " " " " " "  $\omega_2$   
 = tener destrucción (creación) " " "  $\omega_3$



Tarea 2: Demuestras que la intensidad total  $I_1 + I_2 + I_3$  es conservada.

$$\frac{dI_1}{dz} + \frac{dI_2}{dz} + \frac{dI_3}{dz} = -8d \operatorname{Im} [A_3 A_1^* A_2^* e^{-i\Delta K z}] \underbrace{\{ \omega_1 + \omega_2 - \omega_3 \}}_0 = 0$$





campo  $E_2$  fuerte,  $A_2(z) \approx \text{cte}$  (PUMP) (1)  
 $E_1$  débil

$$\Rightarrow \frac{dA_1}{dz} = K_1 A_3 e^{-i\Delta K z} \quad (2)$$

$$\frac{dA_3}{dz} = K_3 A_1 e^{i\Delta K z} \quad (3)$$

donde  $K_1 = \frac{8\pi i \omega_1^2 d A_2^*}{K_1 c^2}$  ;  $K_3 = \frac{8\pi i \omega_3^2 d A_2}{K_3 c^2}$  (4)

$\Delta K = K_1 + K_2 - K_3$  parámetro de mismatch. (5)

(a)  $\Delta K = 0$  :  $\frac{dA_1}{dz} = K_1 A_3$  ;  $\frac{dA_3}{dz} = K_3 A_1$

$$\Rightarrow \frac{d^2 A_1}{dz^2} = K_1 \frac{dA_3}{dz} = K_1 K_3 A_1 = -K^2 A_1 \quad (5')$$

donde  $K^2 = \frac{64\pi^2 \omega_1^2 \omega_3^2 d^2 |A_2|^2}{K_1 K_3 c^4}$  (6')

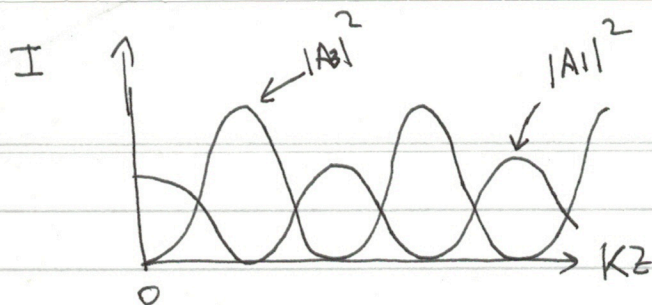
luego  $A_1(z) = B \cos(Kz) + C \sin(Kz)$

y  $A_3(z) = \frac{1}{K_1} \frac{dA_1}{dz} = -\frac{BK}{K_1} \sin(Kz) + \frac{CK}{K_1} \cos(Kz)$

Si  $A_3(0) = 0$  ~  $A_1(0) \neq 0$ ,

$$\Rightarrow \begin{cases} A_1(z) = A_1(0) \cos(Kz) \\ A_3(z) = -A_1(0) \frac{K}{K_1} \sin(Kz) \end{cases} \quad (7)$$





$$(A_2 = \text{cte}, \Delta K = 0)$$

$$(b) \Delta K \neq 0 \quad \left. \begin{aligned} \dot{A}_1 &= K_1 A_3 e^{-i\Delta K z} + i\Delta K A_1 e^{+i\Delta K z} \\ A_3 &= K_3 A_1 e^{-i\Delta K z} \end{aligned} \right\} (8)$$

$$\Rightarrow \dot{A}_1 = K_1 A_3 e^{-i\Delta K z} + K_1 A_3 (-i\Delta K) e^{-i\Delta K z} \\ = K_1 e^{-i\Delta K z} (K_3 A_1) e^{-i\Delta K z} - i K_1 \Delta K A_3 e^{-i\Delta K z}$$

$$= K_1 K_3 A_1 - \cancel{i K_1 \Delta K A_1 e^{-i\Delta K z}} \quad = K_1 K_3 A_1 - i \Delta K A_1$$

$$\therefore \frac{d^2 A_1}{dz^2} = K_1 K_3 A_1 - i \Delta K \frac{dA_1}{dz} \quad (9)$$

$$A_1 \sim e^{\lambda z} : \lambda^2 = K_1 K_3 - i \Delta K \lambda$$

$$\lambda^2 + i \Delta K \lambda - K_1 K_3 = 0 \Rightarrow \lambda = \frac{-i \Delta K \pm \sqrt{-\Delta K^2 + 4 K_1 K_3}}{2}$$

$$\therefore \text{see, } \lambda = -i \frac{\Delta K}{2} \pm \frac{1}{2} \sqrt{\Delta K^2 + 4 K_1 K_3}$$

$$\therefore A_1(z) = e^{-i \frac{\Delta K}{2} z} \left\{ \alpha e^{i \sqrt{K^2 + (\Delta K/2)^2} z} + \beta e^{-i \sqrt{K^2 + (\Delta K/2)^2} z} \right\}$$

$$A_1(z) = e^{-i \frac{\Delta K}{2} z} \left\{ \alpha \cos(qz) + \beta \sin(qz) \right\} \quad (10)$$

$$\text{where } q(z) = \sqrt{K^2 + (\Delta K/2)^2}$$

On se pose,  $A_3(z) = \bar{K}_1 e^{-i \Delta K z} A_1$

$$= \bar{K}_1 e^{-i \Delta K z} \left\{ -i \frac{\Delta K}{2} e^{-i \frac{\Delta K z}{2}} [\alpha \cos(qz) + \beta \sin(qz)] + e^{-i \frac{\Delta K z}{2}} [-\alpha q \sin(qz) + \beta q \cos(qz)] \right\}$$

$$= \bar{K}_1 e^{-i \frac{\Delta K z}{2}} \left\{ \left( \alpha + \beta q \right) \cos(qz) - \left( \alpha q - i \frac{\Delta K}{2} \beta \right) \sin(qz) \right\} \quad (11)$$

$$A_1(0) = \alpha$$

$$A_3(0) = \bar{K}_1 \left( -i \frac{\Delta K}{2} \alpha + \beta q \right) = \bar{K}_1 \left( -i \frac{\Delta K}{2} A_1(0) + \beta q \right)$$

$$A_3(0) = -\frac{i \Delta K}{2 \bar{K}_1} A_1(0) + \beta \frac{q}{\bar{K}_1} \Rightarrow \boxed{\beta = \frac{\bar{K}_1}{q} A_3(0) + i \frac{\Delta K}{2 q} A_1(0)} \quad (12)$$

$$\therefore \left. \begin{aligned} A_1(z) &= e^{-i \frac{\Delta K z}{2}} \left[ A_1(0) \cos(qz) + \left( i \frac{\Delta K}{2 q} A_1(0) + \frac{\bar{K}_1}{q} A_3(0) \right) \sin(qz) \right] \\ A_3(z) &= e^{+i \frac{\Delta K z}{2}} \left[ A_3(0) \cos(qz) + \left( -i \frac{\Delta K}{2 q} A_3(0) + \frac{\bar{K}_3}{q} A_1(0) \right) \sin(qz) \right] \end{aligned} \right\} (13)$$

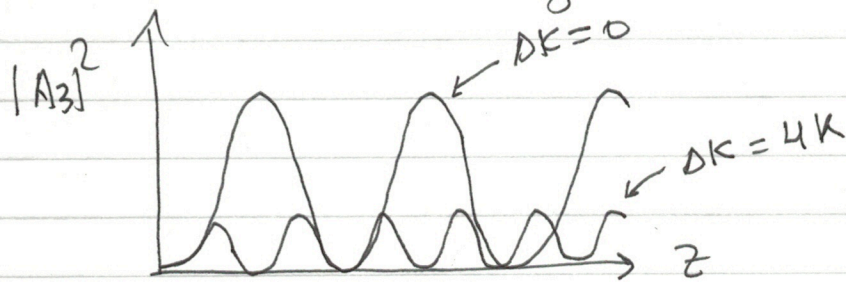
caso special:  $A_3(0) = 0$

$$\Rightarrow \left. \begin{aligned} A_1(z) &= e^{-i \frac{\Delta K z}{2}} \left[ A_1(0) \cos(qz) + i \frac{\Delta K}{2 q} A_1(0) \sin(qz) \right] \\ A_3(z) &= e^{+i \frac{\Delta K z}{2}} \frac{\bar{K}_3}{q} A_1(0) \sin(qz) \end{aligned} \right\} (14)$$

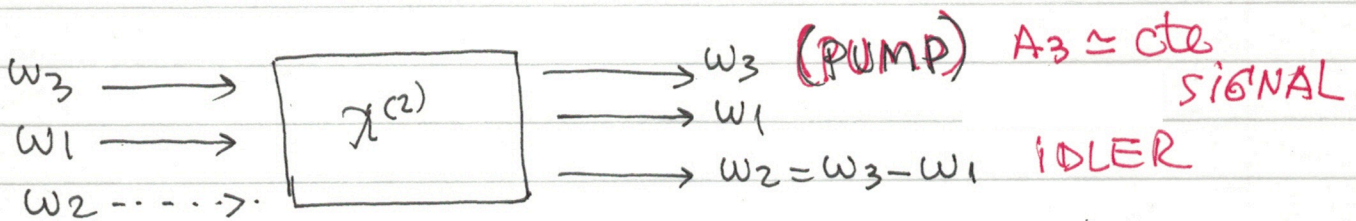


$$\Rightarrow I_3 = |A_3(z)|^2 = |A_1(0)|^2 \frac{|k_3|^2}{g^2} \sin^2(gz)$$

$$(g = \sqrt{k^2 + (\Delta k/2)^2})$$



## Diferencia de frecuencias, amplificación paramétrica



$$\frac{dA_1}{dz} = \frac{8\pi i \omega_1^2 d}{k_1 c^2} A_3 A_2^* e^{i \Delta k z}$$

$$\frac{dA_2}{dz} = \frac{8\pi i \omega_2^2 d}{k_2 c^2} A_3 A_1^* e^{i \Delta k z}$$

(1)

$$(2) \quad \underline{\Delta k = 0} : \frac{d^2 A_2}{dz^2} = \frac{64\pi^2 \omega_1^2 \omega_2^2 d^2}{k_1 k_2 c^4} A_3 A_3^* A_2 \equiv K A_2 \quad (2)$$

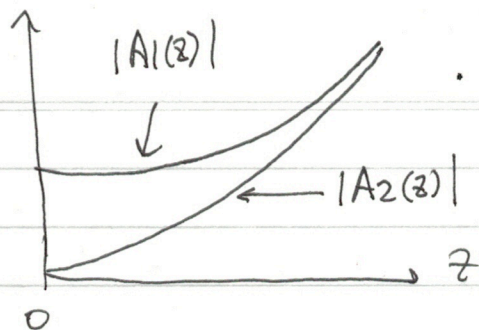
$$\Rightarrow A_2(z) = C \sinh(Kz) + D \cosh(Kz) \quad (3)$$

Si  $A_2(0) = 0$ ,  $A_1(0) = \text{arbitrario}$

$$\Rightarrow A_1(z) = A_1(0) \cosh(Kz)$$

$$A_2(z) = i \frac{8\pi \omega_2^2 d |A_3|^2}{A_3} A_1(0) \sinh(Kz)$$

(4)



(b)  $\Delta K \neq 0$  (Demonstrates Coupled Modes)

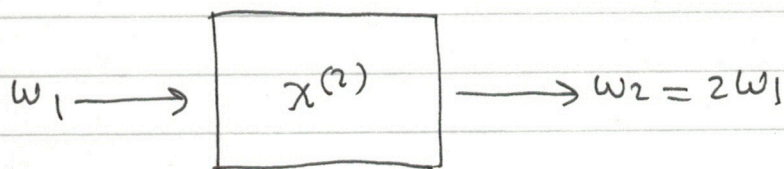
$$A_1(z) = \left[ A_1(0) \left( \cosh(qz) - \frac{i\Delta K}{2q} \sinh(qz) \right) + \frac{K_1}{q} A_2^*(0) \sinh(qz) \right] e^{i\frac{\Delta K}{2}z}$$

$$A_2(z) = \left[ A_2(0) \left( \cosh(qz) - \frac{i\Delta K}{2q} \sinh(qz) \right) + \frac{K_2}{q} A_1^*(0) \sinh(qz) \right] e^{i\frac{\Delta K}{2}z}$$

where  $q \equiv \sqrt{K_1 K_2^* - (\Delta K/2)^2}$ ,  $K_j \equiv \frac{8\pi i \omega_j^2 d A_3}{K_j c^2}$ .



## Generation de 2<sup>de</sup> harmonique



$$E(z,t) = E_1(z,t) + E_2(z,t) \quad (1)$$

$$E_j(z,t) = E_j(z) e^{-i\omega_j t} + c.c. \quad (1')$$

$$E_j(z) = A_j(z) e^{i k_j z} \quad (1'')$$

donc  $k_j = n_j \omega_j / c$ ,  $n_j = \sqrt{\epsilon^{(1)}(\omega_j)}$

$$\frac{\partial^2 E_j}{\partial z^2} - \frac{\epsilon^{(1)}(\omega_j)}{c^2} \frac{\partial^2 E_j}{\partial t^2} = \frac{4\pi}{c^2} \frac{\partial^2 P_j}{\partial t^2} \quad (2)$$

$$P_{NL}(z,t) = P_1(z,t) + P_2(z,t), \quad (3)$$

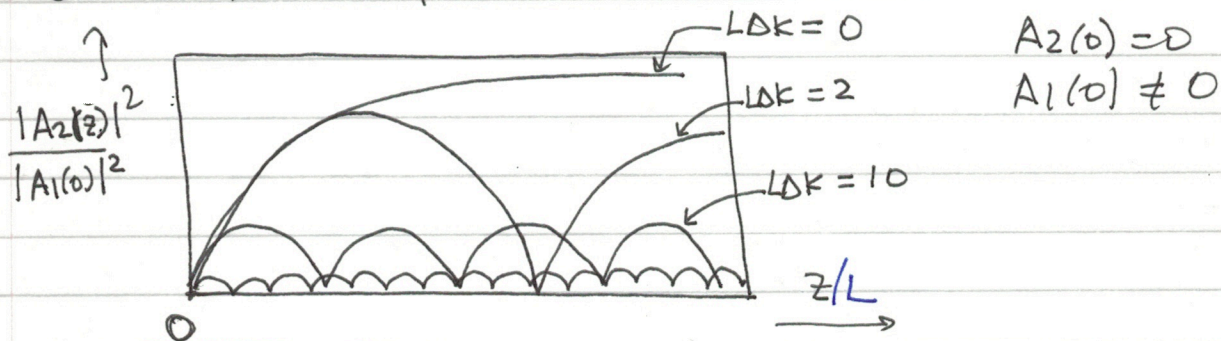
car  $P_j(z,t) = P_j(z) e^{-i\omega_j t} + c.c. \quad (j=1,2) \quad (3')$

donc  $P_1(z) = 2\chi^{(2)} E_2 E_1^* = \overset{4d}{2\chi^{(2)}} A_2 A_1^* e^{i(k_2 - k_1)z}$   
 $P_2(z) = \chi^{(2)} E_1^2 = \overset{2d}{\chi^{(2)}} A_1^2 e^{i2k_1 z}$  } (4)

$$\Rightarrow \left. \begin{aligned} \frac{dA_1}{dz} &= \frac{8\pi i \omega_1^2 d}{k_1 c^2} A_2 A_1^* e^{-i\Delta k z} \\ \frac{dA_2}{dz} &= \frac{4\pi i \omega_2^2 d}{k_2 c^2} A_1^2 e^{i\Delta k z} \end{aligned} \right\} \quad (5)$$

donc  $\Delta k \equiv 2k_1 - k_2$

Sistema (5) puede ser resuelto en forma exacta, por medio de f.m. elípticas de Jacobi.



Tarea: Integrar (5) numéricamente y graficar

$$\frac{I_2(z)}{I} \text{ v/s } \frac{z}{L} \text{ para } A_1(0) = \sqrt{I} \text{ y } A_2(0) = 0$$

para varios valores de  $\Delta S = L\Delta K$

Para  $\Delta K \neq 0$ , la eficiencia de conversión a 2<sup>da</sup> armónica decrece drásticamente.